

ONE DIMENSIONAL INVERSE SPECTRAL BOUNDARY PROBLEM

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STUDENTS' SEMINAR

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Contents

- 1 What is an Inverse problem?
- 2 Properties of 2nd order differential operators
- 3 Formulation of the main problem
- 4 The sketch of proofs

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An Inverse problem

5	1	3	9
4	2	2	8
3	5	7	15
12	8	12	

Direct problem

			2
			7
			12
10	5	6	

Inverse problem

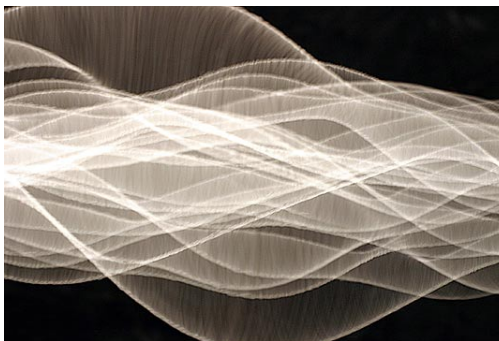
- Does this problem have a solution?
- If there is a solution, is it unique?
- Do we have some prior information about the numbers?

Vibrating string 1

Let $a, b \in \mathbb{R}$, $a \neq 0$. Recall that the 2nd order equation

$$\begin{cases} a^2 \frac{d^2}{dx^2} v(x) = 0, & x \in (0, 1) \\ u(0) = 0, & \frac{d}{dx} u(0) = b, \end{cases}$$

describes the the motion of a vibrating string. Here a is related to the material parameters of the string.



Vibrating string 2

Direct Problem: If numbers a, b are given then

$$\begin{cases} a^2 \frac{d^2}{dx^2} v(x) = 0, & x \in (0, 1) \\ u(0) = 0, & \frac{d}{dx} u(0) = b, \end{cases}$$

has a unique solution.

Inverse Problem: Find a , if some information about the operator $a^2 \frac{d^2}{dx^2}$ is given.

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Definitions and domain

In this talk we will consider 2nd order differential operators that have a general form

$$A = a(x) \frac{d^2}{dx^2} + b(x) \frac{d}{dx} + c(x)$$

where $x \in [0, 1]$ and $a, b, c \in C^\infty([0, 1])$, $a(x) > 0$ and $a(0) = 1$.

We define the domain of A

$$D(A) := H_0^1(0, 1) \cap H^2(0, 1) \sim \{f \in C^2(0, 1) : f(0) = f(1) = 0\}$$

Then $A : D(A) \rightarrow L^2(0, 1)$.

Spectrum of a differential operator

Recall that a function $\varphi \in D(A)$ is an eigen function of A , if $\varphi \neq 0$ and there exists $\lambda \in \mathbb{R}$ such that

$$A\varphi = \lambda\varphi.$$

Actually one can show that every eigen function of A is smooth.

Theorem

There exists a L^2 -orthonormal sequence $(\varphi_j)_{j=1}^{\infty} \subset D(A)$ of eigen functions of differential operator A such that

- $0 < \lambda_1 \leq \lambda_2 \leq \lambda_3 \leq \lambda_4 \dots \rightarrow \infty$
- $\varphi_1(x) \neq 0, x \in (0, 1)$
- For any $f \in L^2(0, 1)$, $f(x) = \sum_{j=1}^{\infty} (f|\varphi_j)_2 \varphi_j(x)$

Proof: Take a course PDE 2 (Spring 2017) or Spectral theorem (Fall 2016)

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Spectral boundary data

We say that the spectral boundary data (SBD) of differential operator A is the collection

$$\{(\lambda_j)_{j=1}^{\infty}, (\dot{\varphi}_j(0))_{j=1}^{\infty}\}.$$

Problem (Inverse spectral boundary problem)

Let

$$A = a(x) \frac{d^2}{dx^2} + b(x) \frac{d}{dx} + c(x)$$

where $x \in [0, 1]$ and $a, b, c \in C^\infty([0, 1])$, $a(x) > 0$, $a(0) = 1$.
Suppose that the spectral boundary data

$$\{(\lambda_j)_{j=1}^{\infty}, (\dot{\varphi}_j(0))_{j=1}^{\infty}\}$$

is given. Can you find functions a, b, c ?

Gauge transformations

Let $\kappa \in C^\infty([0, 1])$ such that

$$\kappa(0) = 1, \text{ and } \kappa(x) > 0, x \in [0, 1].$$

We define a Gauge transformation A_κ of differential operator A by formula

$$A_\kappa f = \kappa A \begin{pmatrix} f \\ \kappa \end{pmatrix}, f \in D(A).$$

Let $\varphi \in D(A)$ be an eigen function of A w.r.t. eigen value λ . Then function $\varphi_\kappa := \kappa\varphi$ satisfies

$$A_\kappa \varphi_\kappa = \lambda \varphi_\kappa \text{ and } \frac{d}{dx} \varphi_\kappa(0) = \dot{\kappa}(0)\varphi(0) + \kappa(0)\dot{\varphi}(0) = \dot{\varphi}(0).$$

Therefore any Gauge transform of operator A preserves the SBD.

Changes of Coordinates 1

Let $\ell > 0$, $X : [0, \ell] \rightarrow [0, 1]$ be a smooth function such that

$$\dot{X}(y) > 0, X(0) = 0, \dot{X}(0) = 1 \text{ and } X(\ell) = 1,$$

Any such a function is called a change of coordinates. Recall that in these new coordinates we have

$$\frac{d}{dx} = \left(\frac{dX}{dy} \right)^{-1} \frac{d}{dy}$$

and

$$\frac{d^2}{dx^2} = \left(\frac{dX}{dy} \right)^{-2} \left[\frac{d^2}{dy^2} - \frac{d^2X}{dy^2} \left(\frac{dX}{dy} \right)^{-1} \frac{d}{dy} \right].$$

Changes of Coordinates 2

Thus operator A transforms to operator A^X defined as

$$A^X f(y) := a_X(y) \frac{d^2}{dy^2} f(y) + b_X(y) \frac{d}{dy} f(y) + c_X(y) f(y),$$

where

$$a_X(y) = a(X(y)) \dot{X}(y)^{-2}$$

$$b_X(y) = -a(X(y)) \dot{X}(y)^{-3} \ddot{X}(y) + \dot{X}(y)^{-1} b(X(y))$$

$$c_X(y) = c(X(y)).$$

Let $\varphi \in D(A)$ be an eigen function of A w.r.t. eigen value λ .

Define $\varphi_X := \varphi \circ X$. Then

$$A_X \varphi_X = \lambda \varphi_X \text{ and } \dot{\varphi}_X(0) = \dot{\varphi}(X(0)) \dot{X}(0) = \dot{\varphi}(0).$$

Thus a change of coordinates preserves SBD

The invariance of the spectral boundary data

Theorem

Let A and B be two second order differential operators as before. Then SBD of A coincides with SBD of B if and only if there exists a change of coordinates X and a Gauge transform κ such that

$$B = (A^X)_{\kappa}.$$

Main theorem

We consider 2nd order differential operators with special forms

$$A := -\frac{d^2}{dx^2} + q(x) \text{ and } B := -a(x)^2 \frac{d^2}{dx^2}$$

where $q, a \in C^\infty([0, 1])$, $a(x) > 0$, $x \in [0, 1]$ and $a(0) = 1$.

Theorem (Inverse spectral boundary problem)

Suppose that the boundary spectral data

$$\{(\lambda_j)_{j=1}^\infty, (\dot{\psi}_j(0))_{j=1}^\infty\}$$

of operator A (respectively B) is given. Then we can reconstruct the potential q (respectively the wave speed a).

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We will provide a proof for the case

$$A := \frac{d^2}{dx^2} + q(x).$$

All we need to do is to recover the first eigen function φ_1 since then

$$q(x) = \frac{\frac{d^2}{dx^2}\varphi_1(x) + \lambda_1\varphi_1(x)}{\varphi_1(x)}.$$

Recall that we know that $\varphi_1(x) \neq 0$.

Initial/Boundary value problem of Wave equation

To solve the Inverse spectral boundary problem we will employ one dimensional wave equation

$$(*) \begin{cases} \left(\frac{d^2}{dt^2} - \frac{d^2}{dx^2} + q(x) \right) u(t, x) = 0, & (t, x) \in (0, 1) \times (0, 1) \\ u(t, 0) = f(t), \quad u(t, 1) = 0 \\ u(0, x) = \frac{\partial}{\partial t} u(0, x) = 0, \end{cases}$$

where $f \in C_0^\infty(0, 1)$ is called a boundary source.

Theorem

Let $f \in C_0^\infty(0, 1)$. Then there exists a unique $u^f(t, x) \in C^\infty((0, 1) \times (0, 1))$ that solves $()$.*

Proof: Take course PDE 1 next fall!

A series representation of waves

Recall that $(\varphi_j)_{j=1}^{\infty} \subset C^{\infty}((0, 1))$ is an ON basis of $L^2(0, 1)$.

Therefore for every boundary source $f \in C_0^{\infty}((0, 1))$ we can write

$$u^f(t, x) = \sum_{j=1}^{\infty} u_j^f(t) \varphi_j(x),$$

where the Fourier coefficients are given by

$$u_j^f(t) := (u^f(t, \cdot) | \varphi_j)_{L^2(0,1)} = \int_0^1 u^f(t, x) \varphi_j(x) dx.$$

Theorem (Fourier coefficients of waves)

For any $f \in C^{\infty}(0, 1)$ we can find the Fourier coefficients $u_j^f(t)$ from SBD.

Finding the Fourier coefficients from SBD 1

Since u^f is smooth we can differentiate under the integral to get

$$\begin{aligned}
 \frac{d^2}{dt^2} u_j^f(t) &= \int_0^1 \frac{\partial^2}{\partial t^2} u^f(t, x) \varphi_j(x) dx \\
 &= \int_0^1 \left[\frac{\partial^2}{\partial x^2} u^f(t, x) - q(x) u^f(t, x) \right] \varphi_j(x) dx \\
 &= \int_0^1 u^f(t, x) \underbrace{\left[\frac{\partial^2}{\partial x^2} \varphi_j(x) - q(x) \varphi_j(x) \right]}_{=-\lambda_j \varphi_j(x)} dx \\
 &\quad + \frac{\partial}{\partial x} u^f(1, t) \underbrace{\varphi_j(1)}_{=0} - \underbrace{u^f(1, t)}_{=0} \frac{\partial}{\partial x} \varphi_j(1) \\
 &\quad - \frac{\partial}{\partial x} u^f(0, t) \underbrace{\varphi_j(0)}_{=0} + \underbrace{u^f(0, t)}_{=f(t)} \frac{\partial}{\partial x} \varphi_j(0)
 \end{aligned}$$

Finding the Fourier coefficients from SBD 2

Thus we obtain the following initial value problem:

$$\begin{cases} \frac{d^2}{dt^2} u_j^f(t) = -\lambda_j u_j^f(t) + \dot{\varphi}_j(0) f(t) \\ u_j^f(0) = \frac{d}{dt} u_j^f(0) = 0. \end{cases}$$

Solution: Take courses ODE 1 and ODE 2 (Spring 2017).

Thus we conclude that for all $f, h \in C_0^\infty(0, 1)$ we have recovered the Fourier coefficients

$$(u_j^f(t))_{j=1}^\infty, \text{ of the wave } u^f(t, x).$$

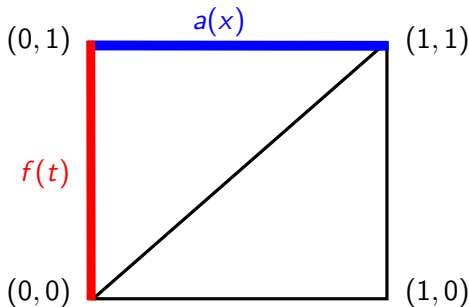
and the inner products

$$(u^f(t, \cdot) | u^h(t, \cdot))_{L^2(0,1)} = \sum_{j=1}^{\infty} u_j^f(t) u_j^h(t)$$

This is the Parseval identity (Funktionaali analyysin peruskurssi Spring 2017).

Controllability

Next we ask can we control the end state of a wave. I.e.



Theorem (Controllability)

Let $a \in C^\infty(0,1)$. There exists a unique $f \in C^\infty(0,1)$ such that

$$u^f(1, x) = a(x).$$

Projectors

Let $t \in [0, 1]$ then we define a projection

$$P_t : L^2(0, 1) \rightarrow L^2(0, 1), P_t(f) = \chi_{[0,t]} f.$$

Define a function $M_{jk} : [0, 1] \rightarrow \mathbb{R}$ by formula

$$M_{jk}(t) = (P_t \varphi_j | \varphi_k)_{L^2(0,1)} = \int_0^t \varphi_j(x) \varphi_k(x) dx.$$

Suppose that function M_{11} is known then

$$\frac{d}{dt} M_{11}(t) = \varphi_1(t)^2 \Rightarrow \text{eigen function } \varphi_1 \text{ is recovered.}$$

Recovery of matrix valued mapping $t \mapsto M_{jk}(t)$

Let $t_0 \in [0, 1]$.

- Choose any smooth orthogonal basis $(g_k)_{k=1}^{\infty}$ of $L^2(0, t_0)$. By controllability theorem

$$\text{span}(u^{g_k}(t_0, \cdot))_{k=1}^{\infty} \subset L^2(0, t_0) \text{ is dense.}$$

- Use Gram-Schmidt to orthonormalise $u^{g_k}(t_0, \cdot)$ to orthonormal basis $(v_k)_{k=1}^{\infty}$ of $L^2(0, t_0)$.
- Since solution mapping $f \mapsto u^f$ is linear it holds that

$$v_k(x) = u^{f_k}(x, t_0), \quad f_k(t) := \sum_{j=1}^k d_{jk} g_j(t), \quad d_{jk} \in \mathbb{R}.$$

- Since $(v_k)_{k=1}^{\infty}$ of $L^2(0, t_0)$ is ON-basis it holds that

$$P_{t_0} \varphi_j = \sum_{\ell=1}^{\infty} (\varphi_j | v_{\ell})_{L^2(0, t_0)} v_{\ell}$$

Thus

$$M_{jk}(t_0) = (P_{t_0} \varphi_j | \varphi_k)_{L^2(0, 1)} = \sum_{\ell=1}^{\infty} (\varphi_j | v_{\ell})_{L^2(0, 1)} (\varphi_k | v_{\ell})_{L^2(0, 1)}$$

- Notice that $(\varphi_j | v_{\ell})_{L^2(0, 1)}$ is a Fourier coefficient of v_{ℓ} w.r.t basis $(\varphi_j)_{j=1}^{\infty}$ i.e

$$(\varphi_j | v_{\ell})_{L^2(0, 1)} = u_j^{f_{\ell}}(t_0).$$

By the Theorem for the Fourier coefficients of waves, we can recover these from SBD.

Thank you for your attention!