

CORRELATION BASED PASSIVE IMAGING WITH A WHITE NOISE SOURCE

TAPIO HELIN^a, MATTI LASSAS^a, LAURI OKSANEN^b AND TEEMU SAKSALA^a

^aDepartment of Mathematics and Statistics, University of Helsinki, Finland

^bDepartment of Mathematics, University College London, UK

ABSTRACT. Passive imaging refers to problems where waves generated by unknown sources are recorded and used to image the medium through which they travel. The sources are typically modelled as a random variable and it is assumed that some statistical information is available. In this paper we study the stochastic wave equation $\partial_t^2 u - \Delta_g u = \chi W$, where W is a random variable with the white noise statistics on \mathbb{R}^{1+n} , $n \geq 3$, χ is a smooth function vanishing for negative times and outside a compact set in space, and Δ_g is the Laplace–Beltrami operator associated to a smooth non-trapping Riemannian metric tensor g on \mathbb{R}^n . The metric tensor g models the medium to be imaged, and we assume that it coincides with the Euclidean metric outside a compact set. We consider the empirical correlations on an open set $\mathcal{X} \subset \mathbb{R}^n$,

$$C_T(t_1, x_1, t_2, x_2) = \frac{1}{T} \int_0^T u(t_1 + s, x_1) u(t_2 + s, x_2) ds, \quad t_1, t_2 > 0, \quad x_1, x_2 \in \mathcal{X},$$

for $T > 0$. Supposing that χ is non-zero on \mathcal{X} and constant in time after $t > 1$, we show that in the limit $T \rightarrow \infty$, the data C_T becomes statistically stable, that is, independent of the realization of W . Our main result is that, with probability one, this limit determines the Riemannian manifold (\mathbb{R}^n, g) up to an isometry. To our knowledge, this is the first result showing that a medium can be determined in a passive imaging setting, without assuming a separation of scales.

1. INTRODUCTION

In passive imaging, waves generated by unknown sources are recorded and used to image the medium through which they travel. Passiveness refers to the observer having only little or no control over the source (think earthquakes in seismic imaging). However, some statistical information of the source may be available and it can be useful to model the source as a random variable: while the statistics of the random variable is known, its realization remains unknown.

Passive imaging has had a fundamental impact to seismic and various other imaging modalities. We refer to the recent book by Garnier and Papanicolaou [18] for an

Date: October 18, 2016.

extensive review of the field. The previous mathematical theory is, to a large extent, based on assuming some physical scaling regime. Such an approach has produced a number of important and efficient numerical methods. However, our key finding in the present paper is that exact recovery of an unknown medium is also possible without any scaling assumptions. The proof of this is based on a reduction to a deterministic inverse problem.

In this work we consider the wave equation

$$(1) \quad \begin{aligned} \partial_t^2 u(t, x) - \Delta_g u(t, x) &= \chi(t, x)W(t, x) \quad \text{in } \mathbb{R}_+^{1+n} = (0, \infty) \times \mathbb{R}^n, \\ u|_{t=0} &= \partial_t u|_{t=0} = 0, \end{aligned}$$

where $n \geq 3$ and Δ_g is the Laplace–Beltrami operator corresponding to a smooth time-independent Riemannian metric g on \mathbb{R}^n . In coordinates $(x_j)_{j=1}^n$ this operator has the following representation.

$$\Delta_g = \sum_{j,k=1}^n |g|^{-1/2} \frac{\partial}{\partial x^j} \left(|g|^{1/2} g^{jk} \frac{\partial}{\partial x^k} u \right),$$

where $[g_{jk}]_{j,k=1}^n = g(x)$, $|g| = \det(g_{jk})$ and $[g^{jk}]_{j,k=1}^n = g(x)^{-1}$. We assume that our source W is a realization of a Gaussian white noise random variable on \mathbb{R}^{1+n} . Moreover, χ stands for a smooth function

$$\chi(t, x) = \chi_0(t)\kappa(x),$$

such that $\chi_0 \in C^\infty(\mathbb{R})$ and

$$\chi_0(t) = \begin{cases} 0, & t \leq 0, \\ 1, & t \geq 1, \end{cases}$$

and $\kappa \in C_0^\infty(\mathbb{R}^n)$. We assume that there exists an open and non-empty set $\mathcal{X} \subset \mathbb{R}^n$ where κ is non-vanishing. The source χW can be modelled as a random variable taking values in a local Sobolev space with negative index, and the same is true for the solution u . Contrary to papers such as [10, 37, 40], we do not consider $t \mapsto u(t, \cdot)$ as a random process.

The problem we study is the following: suppose we can record the empirical correlation

$$(2) \quad C_T(t_1, x_1, t_2, x_2) = \frac{1}{T} \int_0^T u(t_1 + s, x_1) u(t_2 + s, x_2) ds,$$

for $t_1, t_2 > 0$, $x_1, x_2 \in \mathcal{X}$ and $T > 0$. What information does this data yield regarding the metric g ? For any finite T , the correlation C_T is random in the sense that it depends on the realization of the source. A fundamental part of our result below is to

show that this data becomes *statistically stable*, i.e. independent of the realization, as T increases. More precisely, we show that the limit

$$\lim_{T \rightarrow \infty} \langle C_T, f \otimes h \rangle_{\mathcal{D}' \times C_0^\infty(\mathbb{R}^{2+2n})}, \quad f, g \in C_0^\infty(\mathbb{R}^{1+n}),$$

is deterministic, see Theorem 3 below. Thereafter, the paper is devoted to showing that this stability enables the recovery of g :

Theorem 1. *Let $n \geq 3$. Suppose that g is non-trapping and that g coincides with the Euclidean metric outside a compact set. Let $u = \mathbb{U}(\omega)$ be the solution of (1) where $W = \mathbb{W}(\omega)$ is a realization of the white noise \mathbb{W} on \mathbb{R}^{1+n} . Then with probability one, the empirical correlations (2) defined in the sense of generalized random variables in $\mathcal{D}'((\mathbb{R} \times \mathcal{X})^2)$ for $T > 0$, determine the Riemannian manifold (\mathbb{R}^n, g) up to an isometry.*

Recall that a metric tensor g on \mathbb{R}^n is non-trapping if for each compact $K \subset \mathbb{R}^n$ there exists $T > 0$ such that for each $(p, \xi) \in T\mathbb{R}^n$, $p \in K$, $\|\xi\|_g = 1$, it holds that $\gamma_{p,\xi}(t) \notin K$ when $t \geq T$. Here we denote by $\gamma_{p,\xi}$ the unique maximal geodesic of metric g that satisfies the following initial conditions

$$\gamma_{p,\xi}(0) = p \text{ and } \dot{\gamma}_{p,\xi}(0) = \xi.$$

Note that the covariance data (2) is determined by the measurement $u|_{(0,\infty) \times \mathcal{X}}$. This implies the following corollary:

Corollary 1. *The measurement $u|_{(0,\infty) \times \mathcal{X}}$, with a single realization of the white noise source, determines the Riemannian manifold (\mathbb{R}^n, g) , up to an isometry, with probability one under the assumptions of Theorem 1.*

The statistical stability of C_T , $T > 0$, allows us to reduce the passive imaging problem to a deterministic inverse problem, that we then solve. As this deterministic problem is of independent interest, we solve it in a more general geometric setting formulated as follows:

Theorem 2. *Let (N, g) be a smooth and complete Riemannian manifold of dimension $n \geq 2$. Let $\mathcal{X} \subset N$ be an open set. Consider the following initial value problem for the wave equation*

$$(3) \quad \begin{aligned} \partial_t^2 w(t, x) - \Delta_g w(t, x) &= f, \quad \text{in } (0, \infty) \times N, \\ w|_{t=0} = \partial_t w|_{t=0} &= 0. \end{aligned}$$

Let $\Lambda_{\mathcal{X}} : C_0^\infty((0, \infty) \times \mathcal{X}) \rightarrow C^\infty((0, \infty) \times \mathcal{X})$ be the local source-to-solution operator defined by

$$\Lambda_{\mathcal{X}} f = w|_{(0,\infty) \times \mathcal{X}}.$$

Then the data $(\mathcal{X}, \Lambda_{\mathcal{X}})$ determines (N, g) up to an isometry. More precisely this means the following:

Let (N_i, g_i) , $i = 1, 2$, be a smooth and complete Riemannian manifold. Let $\mathcal{X}_i \subset N_i$ be an open set, and assume that there exists a diffeomorphism

$$(4) \quad \phi : \mathcal{X}_1 \rightarrow \mathcal{X}_2$$

that satisfies

$$(5) \quad \phi^*(\Lambda_{\mathcal{X}_2} f) = \Lambda_{\mathcal{X}_1}(\phi^* f), \quad \text{for all } f \in C_0^\infty((0, \infty) \times \mathcal{X}_2).$$

Then (N_1, g_1) and (N_2, g_2) are Riemannian isometric.

Above the pullback ϕ^* of ϕ is defined by $\phi^* f = f \circ \tilde{\phi}$, where $\tilde{\phi}$ is the lift of ϕ on $(0, \infty) \times \mathcal{X}_1$, that is, $\tilde{\phi}(t, x) = (t, \phi(x))$ $t > 0$, $x \in \mathcal{X}_1$.

Lastly we will point the connection of Theorem 2 to the following Inverse spectral problem of Laplace-Beltrami operator.

Corollary 2. *Let (N, g) be a smooth and compact Riemannian manifold of dimension $n \geq 2$ with out boundary. Let $\mathcal{X} \subset N$ be an open set. Let $(\varphi_k)_{k=1}^\infty \subset C^\infty(N)$ be the collection of orthonormal eigenfunctions of operator Δ_g in $L^2(N)$. Let $(\lambda_k)_{k=1}^\infty$ be the collection of corresponding eigenvalues of Δ_g . Then the Spectral data*

$$(6) \quad (\mathcal{X}, (\varphi_k|_{\mathcal{X}})_{k=1}^\infty, (\lambda_k)_{k=1}^\infty)$$

determines (N, g) up to isometry.

1.1. Outline the paper. We begin by showing that the empirical correlation C_T is well-defined in Section 2. In Section 3 we show the statistical stability discussed above, and in Section 4 we reduce the proof of Theorem 1 to that of Theorem 2. We prove Theorem 2 in Section 5. For the convenience of the reader, we have collected some well-known results in an appendix.

1.2. Previous literature. For previous mathematical results on passive imaging problems we refer to [12, 17]. The monograph [18] gives a thorough review of the related literature. Passive imaging problems arise in geophysical applications. In seismic imaging ambient noise sources, that appear due to nonlinear interaction of ocean waves with the ocean bottom, can be utilized to image the wave speed in the subsurface of the Earth, see e.g. [42, 43, 52].

We also mention the closely related topic of imaging random media by time reversal techniques [2, 7, 8, 14] as well as inverse scattering from random potential or random boundary conditions [9, 22, 34].

Let us now turn to results on deterministic inverse problems similar to Theorem 2. In such coefficient determination problems, it is typical to use the Dirichlet-to-Neumann map to model the data. Apart from immediate applications, this is reasonable since several other types of data can be reduced to the Dirichlet-to-Neumann case. For instance, in [38] an inverse scattering problem is solved via a reduction to

the inverse conductivity problem in [44], and the latter uses the Dirichlet-to-Neumann map as data. In the present paper, however, we do not perform a reduction to the Dirichlet-to-Neumann case but adapt techniques originally developed in that case to the case of local source-to-solution map Λ_χ .

The approach that we use is a modification of the Boundary Control method. This method was first developed by Belishev to the acoustic wave equation on \mathbb{R}^n with an isotropic wave speed [4]. A geometric version of the method, suitable when the wave speed is given by a Riemannian metric tensor as in the present paper, was introduced by Belishev and Kurylev [5]. We refer to [29] for a thorough review of the related literature. Local reconstruction of the geometry from the local source-to-solution map Λ_χ has been studied as a part of iterative schemes, see e.g. [25, 32]. In the present paper we give a global uniqueness proof that does not rely on an iterative scheme.

We restrict our attention to the unique solvability of the inverse problem but note that several variants of the Boundary Control method have been studied computationally [3, 13, 27, 41] and stability questions have been investigated [1, 31, 36].

This work continues the line of research started by the authors in [20, 21], where similar unique solvability of the geometry was considered for random and pseudo-random boundary sources. A novel feature of this paper is that we consider passive imaging, when the source is not assumed to be known.

2. THE STOCHASTIC DIRECT PROBLEM

In this section we show that the running averages C_T , $T > 0$, are well-defined as random variables. Let us first recall the concept of generalized Gaussian random variable [19]. A cylindrical set in a locally convex vector space V with the dual V' is of the form

$$\{u \in V \mid (\langle \ell_1, u \rangle, \dots, \langle \ell_k, u \rangle) \in B\},$$

where $k \geq 1$, $\ell_1, \dots, \ell_k \in V'$, and B is a Borel subset of \mathbb{R}^k , i.e., $B \in \mathcal{B}(\mathbb{R}^k)$. Above, we write $\langle \cdot, \cdot \rangle = \langle \cdot, \cdot \rangle_{V' \times V}$ for the dual pairing between V' and V . The σ -algebra generated by cylindrical sets in V is denoted by $\mathcal{B}_c(V)$. Notice that the cylindrical σ -algebra is always a subset of the Borel σ -algebra, and the two σ -algebras are known to coincide if V is a separable Fréchet space [6, Thm. A.3.7.].

We denote the rapidly decaying functions on \mathbb{R}^n by $\mathcal{S}(\mathbb{R}^d)$. The topological dual of $\mathcal{S}(\mathbb{R}^d)$ is the space of tempered distributions $\mathcal{S}'(\mathbb{R}^d)$. It is well-known that $\mathcal{S}'(\mathbb{R}^d)$ is a locally convex topological vector space (even nuclear).

Throughout the paper, let $(\Omega, \mathcal{F}, \mathbb{P})$ stand for a complete probability space.

Definition 1. *A generalized random variable is a measurable function*

$$X : (\Omega, \mathcal{F}) \rightarrow (\mathcal{S}'(\mathbb{R}^d), \mathcal{B}_c(\mathcal{S}'(\mathbb{R}^d))).$$

A generalized random variable X is called Gaussian, if for all $\phi_1, \dots, \phi_k \in \mathcal{S}(\mathbb{R}^d)$, $k \in \mathbb{N}$, the mapping

$$\Omega \ni \omega \mapsto (\langle X(\omega), \phi_1 \rangle, \dots, \langle X(\omega), \phi_k \rangle) \in \mathbb{R}^k$$

is a Gaussian random variable.

The probability law of a generalized Gaussian random variable X is determined by the expectation $\mathbb{E}X$ and the covariance operator $C_X : \mathcal{S}(\mathbb{R}^d) \rightarrow \mathcal{S}'(\mathbb{R}^d)$ defined by

$$(7) \quad \langle \psi_1, C_X \psi_2 \rangle = \mathbb{E}(\langle X - \mathbb{E}X, \psi_1 \rangle \langle X - \mathbb{E}X, \psi_2 \rangle).$$

If X is zero-mean and satisfies $C_X = \iota$, where $\iota : \mathcal{S}(\mathbb{R}^d) \rightarrow \mathcal{S}'(\mathbb{R}^d)$ is the identity operator $\iota(\phi) = \phi$, then X is called Gaussian white noise.

Remark 1. *The construction above is identical for generalized random variables obtaining values in the space of generalized functions $\mathcal{D}'(\mathbb{R}^d)$. This was also the original formulation in [19].*

It was proved by Kusuoka in [33] that for any $\epsilon > 0$, white noise satisfies

$$(8) \quad \mathbb{W} \in H^{-d/2-\epsilon}(\mathbb{R}^d; \langle x \rangle^{-d/2-\epsilon}) \quad \text{almost surely,}$$

where the weight function is defined by $\langle x \rangle = (1 + |x|^2)^{1/2}$. Moreover, we have $H^{-d/2-\epsilon}(\mathbb{R}^d; \langle x \rangle^{-d/2-\epsilon}) \in \mathcal{B}_c(\mathcal{S}'(\mathbb{R}^d))$ (see e.g. [16, Prop. 7]) and therefore we can consider \mathbb{W} as a random variable restricted to $H^{-d/2-\epsilon}(\mathbb{R}^d; \langle x \rangle^{-d/2-\epsilon})$ assigned with the cylindrical σ -algebra. Since the weighted Sobolev space is separable (and Fréchet), the cylindrical σ -algebra coincides with the Borel σ -algebra and \mathbb{W} is Borel measurable in $H^{-d/2-\epsilon}(\mathbb{R}^d; \langle x \rangle^{-d/2-\epsilon})$. Finally, since we have a continuous embedding $H^{-d/2-\epsilon}(\mathbb{R}^d; \langle x \rangle^{-d/2-\epsilon}) \subset H_{loc}^{-d/2-\epsilon}(\mathbb{R}^d)$, we can identify \mathbb{W} as a random variable

$$\mathbb{W} : (\Omega, \mathcal{F}) \rightarrow (H_{loc}^{-d/2-\epsilon}(\mathbb{R}^d), \mathcal{B}(H_{loc}^{-d/2-\epsilon}(\mathbb{R}^d))).$$

We denote by \square_χ^{-1} the solution operator of (1), that is, $\square_\chi^{-1}(W) = u$ where u solves (1) and u is defined to be zero for negative times. Then

$$\square_\chi^{-1} : H_{loc}^\sigma(\mathbb{R}^{1+n}) \rightarrow H_{loc}^{\sigma+1}(\mathbb{R}^{1+n}), \quad \sigma \in \mathbb{R},$$

is continuous, see e.g. [24, Thm. 23.2.4]. We denote by τ^s the translation by $s \in \mathbb{R}$ in time, that is,

$$\tau^s \phi(t) = \phi(t + s), \quad \phi \in C_0^\infty(\mathbb{R}),$$

and extend this definition to $\mathcal{D}'(\mathbb{R})$ by

$$\langle \tau^s w, \phi \rangle_{\mathcal{D}' \times C_0^\infty(\mathbb{R})} = \langle w, \tau^{-s} \phi \rangle_{\mathcal{D}' \times C_0^\infty(\mathbb{R})}.$$

The function

$$\Phi : (0, T) \times \mathbb{R} \rightarrow \mathbb{R}, \quad \Phi(s, t) = \tau^s \phi(t)$$

is smooth, and moreover $\Phi = 0$ when $t \notin (0, T + R)$ where $R > 0$ is such that $\text{supp}(\phi) \subset (0, R)$. Hence function

$$(9) \quad s \mapsto \langle w, \Phi(s, \cdot) \rangle_{\mathcal{D}' \times C_0^\infty(\mathbb{R})} = \langle \tau^s w, \phi \rangle_{\mathcal{D}' \times C_0^\infty(\mathbb{R})}$$

is smooth for all $w \in \mathcal{D}'(\mathbb{R})$ and $\phi \in C_0^\infty(\mathbb{R})$, see [23, Thm. 2.1.3]. An analogous argument shows that

$$s \mapsto \langle \tau^s w \otimes \tau^s w, \phi \rangle_{\mathcal{D}' \times C_0^\infty(\mathbb{R}^{2+2n})}$$

is smooth for all $w \in \mathcal{D}'(\mathbb{R}^{1+n})$ and $\phi \in C_0^\infty(\mathbb{R}^{2+2n})$. Here \otimes denotes the tensor product of distributions, see e.g. [23, Thm. 5.1.1] for the definition.

For a fixed $T > 0$, we define the map

$$A_T(w) = \frac{1}{T} \int_0^T \tau^s w \otimes \tau^s w ds, \quad w \in H_{loc}^\sigma(\mathbb{R}^{1+n}),$$

in the sense of the Pettis integral, that is,

$$\langle A_T(w), \phi \rangle_{\mathcal{D}' \times C_0^\infty(\mathbb{R}^{2+2n})} = \frac{1}{T} \int_0^T \langle \tau^s w \otimes \tau^s w, \phi \rangle_{\mathcal{D}' \times C_0^\infty(\mathbb{R}^{2+2n})} ds.$$

The integral above defines $A_T(w)$ as a generalized function in $\mathcal{D}'(\mathbb{R}^{2+2n})$ and, moreover, yields a continuous map in the following sense:

Lemma 1. *The map $A_T : H_{loc}^{-\sigma}(\mathbb{R}^{1+n}) \rightarrow H_{loc}^{-\sigma}(\mathbb{R}^{2+2n})$, $\sigma \in \mathbb{R}$, is continuous.*

Proof. We recall that the topology of $H_{loc}^{-\sigma}(\mathbb{R}^{1+n})$ is induced by the semi-norms

$$w \mapsto \|\psi w\|_{H^{-\sigma}(\mathbb{R}^{1+n})}, \quad \psi \in C_0^\infty(\mathbb{R}^{1+n}).$$

Let $w_0 \in H_{loc}^{-\sigma}(\mathbb{R}^{1+n})$, $\psi \in C_0^\infty(\mathbb{R}^{2+2n})$ and $\epsilon > 0$. In order to show that A_T is continuous, it is enough to show [47, p. 64] that there are $\tilde{\psi} \in C_0^\infty(\mathbb{R}^{1+n})$ and $\delta > 0$ such that

$$\left\| \tilde{\psi}(w - w_0) \right\|_{H^{-\sigma}(\mathbb{R}^{1+n})} < \delta \quad \text{implies} \quad \|\psi(A_T(w) - A_T(w_0))\|_{H^{-\sigma}(\mathbb{R}^{2+2n})} < \epsilon.$$

We choose $\tilde{\psi} \in C_0^\infty(\mathbb{R}^{1+n})$ so that $(\tilde{\psi} \otimes \tilde{\psi})\tau_1^{-s}\tau_2^{-s}\psi = \tau_1^{-s}\tau_2^{-s}\psi$ for all $s \in (0, T)$. Here τ_j^{-s} , $j = 1, 2$, act in the different time variables. Let $\phi \in H^\sigma(\mathbb{R}^{2+2n})$. It follows that

$$\begin{aligned} & |\langle \psi(A_T(w) - A_T(w_0)), \phi \rangle_{H^{-\sigma} \times H^\sigma(\mathbb{R}^{2+2n})}| \\ & \leq \frac{1}{T} \int_0^T \left\| (\tilde{\psi} \otimes \tilde{\psi})(w \otimes w - w_0 \otimes w_0) \right\|_{H^{-\sigma}(\mathbb{R}^{2+2n})} \left\| \tau_1^{-s}\tau_2^{-s}\psi\phi \right\|_{H^\sigma(\mathbb{R}^{2+2n})} ds \\ & \leq C \left\| \tilde{\psi}(w - w_0) \otimes \tilde{\psi}w + \tilde{\psi}w_0 \otimes \tilde{\psi}(w - w_0) \right\|_{H^{-\sigma}(\mathbb{R}^{2+2n})} \|\phi\|_{H^\sigma(\mathbb{R}^{2+2n})}. \end{aligned}$$

Finally, for small $\delta > 0$

$$\begin{aligned} & \left\| \tilde{\psi}(w - w_0) \otimes \tilde{\psi}w + \tilde{\psi}w_0 \otimes \tilde{\psi}(w - w_0) \right\|_{H^{-\sigma}(\mathbb{R}^{2+2n})} \\ & \leq \delta \left\| \tilde{\psi}w \right\|_{H^{-\sigma}(\mathbb{R}^{1+n})} + \left\| \tilde{\psi}w_0 \right\|_{H^{-\sigma}(\mathbb{R}^{1+n})} \delta \leq C\delta. \end{aligned}$$

□

By combining the continuity results above, we define $C_T(\omega) = A_T(\square_\chi^{-1}(\mathbb{W}(\omega)))$, $T > 0$, and see that

$$C_T : \Omega \rightarrow (H_{loc}^\sigma(\mathbb{R}^{2+2n}), \mathcal{B}(H_{loc}^\sigma(\mathbb{R}^{2+2n}))), \quad \sigma < -\frac{1+n}{2} + 1,$$

is a random variable.

Remark 2. *Since the weighted Sobolev space $H^{-d/2-\epsilon}(\mathbb{R}^d; \langle x \rangle^{-d/2-\epsilon})$ is separable, the random variable \mathbb{W} in (8) has the Radon property [6]. Notice carefully that the Radon property is transferred through any continuous mappings and therefore also C_T is Radon.*

3. THE STOCHASTIC INVERSE PROBLEM AND STATISTICAL STABILITY

For any function $f \in C_0^\infty(\mathbb{R}^{1+n})$, let us define $v^f = v$ as the solution of a time reversed wave equation

$$(10) \quad \begin{aligned} \partial_t^2 v - \Delta_g v &= f \quad \text{in } (-\infty, S) \times \mathbb{R}^n, \\ v|_{t=S} = \partial_t v|_{t=S} &= 0, \end{aligned}$$

where $S \in \mathbb{R}$ is large enough so that $f \in C_0^\infty((-\infty, S) \times \mathbb{R}^n)$. In this section we show the following theorem.

Theorem 3. *Suppose that $n \geq 3$, (\mathbb{R}^n, g) is non-trapping and that g coincides with the Euclidean metric outside a compact set. Let $\mathbb{D} \subset C_0^\infty((0, \infty) \times \mathcal{X})$ be a countable set. There exists $\Omega_0 \subset \Omega$ such that $\mathbb{P}(\Omega_0) = 0$ and for all $\omega \in \Omega \setminus \Omega_0$ and all $f, h \in \mathbb{D}$, it holds that*

$$\lim_{T \rightarrow \infty} \langle C_T(\omega), f \otimes h \rangle_{\mathcal{D}' \times C_0^\infty(\mathbb{R}^{2+2n})} = \langle \kappa v^f, \kappa v^h \rangle_{L^2(\mathbb{R}^{1+n})}.$$

In what follows, we write $\langle \cdot, \cdot \rangle = \langle \cdot, \cdot \rangle_{\mathcal{D}' \times C_0^\infty(\mathbb{R}^{2+2n})}$.

Lemma 2. *Let $W \in \mathcal{D}'(\mathbb{R}^{1+n})$ and $f \in C_0^\infty(\mathbb{R}^{1+n})$ be arbitrary sources in problems (1) and (10), respectively. Moreover, let u and v^f be the corresponding solutions. Then we have the identity*

$$\langle u, f \rangle = \langle W, \chi v^f \rangle.$$

Proof. Suppose that $W \in C_0^\infty(\mathbb{R}^{1+n})$. The general case follows since test functions are dense in distributions. Next, let v and S be as in (10). Using the shorthand notation $\square_g = \partial_t^2 - \Delta_g$, we have that

$$\langle u, f \rangle = \langle u, \square_g v \rangle_{L^2((0,S) \times \mathbb{R}^n)} = \langle \square_g u, v \rangle_{L^2((0,S) \times \mathbb{R}^n)} = \langle W, \chi v \rangle.$$

This proves the claim. \square

Let us recall the following result regarding the local energy decay which is due to Vainberg [48, 49], see [50] for the formulation as below.

Theorem 4. *Let $u \in C^\infty((0, \infty) \times \mathbb{R}^n)$ solve the problem*

$$\begin{aligned} \partial_t^2 u - \Delta_g u &= 0, \quad \text{in } (0, \infty) \times \mathbb{R}^n, \\ u|_{t=0} &= u_0, \quad \partial_t u|_{t=0} = u_1. \end{aligned}$$

Suppose that u_0 and u_1 are compactly supported. Suppose that (\mathbb{R}^n, g) is non-trapping and that g coincides with the Euclidean metric outside a compact set. Then there is $t_0 > 0$ such that u satisfies local energy decay

$$\int_{\mathbb{R}^n} (|\partial_t u(t, x)|^2 + |\nabla u(t, x)|^2) \chi(x) dx \leq C\eta(t)E_0, \quad t > t_0,$$

for any compactly supported function $\chi \in C_0^\infty(\mathbb{R}^n)$. Here we have

$$E_0 = \int_{\mathbb{R}^n} |\nabla u_0(x)|^2 + |u_1(x)|^2 dx, \quad \eta(t) = \begin{cases} e^{-bt}, & n \geq 3 \text{ odd}, \\ t^{-2n}, & n \geq 2 \text{ even}, \end{cases}$$

and the constants $C, b > 0$ depend on g, χ and the supports of u_0 and u_1 .

We need a decay estimate for the norm $\|u(t, \cdot)\|_{L^2(K)}$ where $K \subset \mathbb{R}^n$ is compact.

Lemma 3. *Let (\mathbb{R}^n, g) be as in Theorem 3 and let u be as in Theorem 4. Let $K \subset \mathbb{R}^n$ be compact. Then there is $t_0 > 0$ such that u satisfies*

$$\|u(t, \cdot)\|_{L^2(K)} \leq C\mu(t)E_0, \quad t > t_0,$$

where

$$(11) \quad \mu(t) = \begin{cases} e^{-bt}, & n \geq 3 \text{ odd}, \\ t^{-2n+1}, & n \geq 4 \text{ even}, \end{cases}$$

Proof. To simplify the notation, we assume without loss of generality that $E_0 = 1$. Let $B(r) = \{\|x\| < r\}$ be the Euclidean ball of radius r and write

$$u_r(t) = \frac{1}{|B(r)|} \int_{B(r)} u(t, x) dx,$$

where $|B(r)|$ is the volume of $B(r)$. Theorem 4 implies $|\partial_t u_r(t)| \leq C\eta(t)$ where the constant $C > 0$ depends on $r > 0$ and g . Thus for $t_0 < t < s$,

$$(12) \quad |u_r(t) - u_r(s)| \leq C \int_t^s \eta(\tau) d\tau = C(\mu(t) - \mu(s)).$$

We see that $\lim_{t \rightarrow \infty} u_r(t)$ exists, and denote the limit by $\bar{u}(r)$.

The Poincaré-Wirtinger inequality

$$\|u(t, \cdot) - u_r(t)\|_{L^2(B(r))} \leq C \|\nabla u(t, \cdot)\|_{L^2(B(r))},$$

together with Theorem 4 and (12), implies that

$$(13) \quad \|u(t, \cdot) - \bar{u}(r)\|_{L^2(B(r))} \leq C\eta(t) + |u_r(t) - \bar{u}(r)| \|1\|_{L^2(B(r))} \leq C\mu(t).$$

In particular, for $0 < r_1 < r_2$, $u(t, \cdot) \rightarrow \bar{u}(r_j)$, $j = 1, 2$, in $L^2(B(r_1))$. Thus $\bar{u}(r)$ does not depend on $r > 0$ and we denote it by \bar{u} .

It remains to show that $\bar{u} = 0$. As $u(t)$ is compactly supported, by the finite speed of propagation, the Gagliardo-Nirenberg-Sobolev inequality implies that

$$\|u(t, \cdot)\|_{L^{p^*}(\mathbb{R}^n)} \leq C \|\nabla u(t, \cdot)\|_{L^2(\mathbb{R}^n)},$$

where p^* is the Sobolev conjugate of 2, that is, $1/p^* = 1/2 - 1/n$. Note that $p^* > 2$. We apply Hölder's inequality with $p = p^*/2$ and $1/p + 1/q = 1$,

$$\int_{B(r)} u^2(t, \cdot) dx \leq \|u^2(t, \cdot)\|_{L^p(B(r))} \|1\|_{L^q(B(r))}.$$

The conservation of energy implies that $\|\nabla u(t, \cdot)\|_{L^2(\mathbb{R}^n)}$, $t > 0$, is bounded. Thus $\|u(t, \cdot)\|_{L^2(B(r))}^2 \leq Cr^{n/q}$ with a constant $C > 0$ independent of r .

To get a contradiction, suppose now that $\bar{u} \neq 0$. Then there is $\epsilon > 0$ such that

$$\|\bar{u}\|_{L^2(B(r))}^2 = \bar{u}^2 \|1\|_{L^2(B(r))}^2 = 2\epsilon r^n.$$

By the convergence (13), for all $r > 0$ there is t_r such that $\|u(t_r, \cdot)\|_{L^2(B(r))}^2 \geq \epsilon r^n$. Thus $r^{n-n/q} \leq C$, $r > 0$, which is a contradiction since $q > 1$. \square

Lemma 4. *Let (\mathbb{R}^n, g) be as in Theorem 3. Suppose that $K \subset \mathbb{R}^n$ is compact and $f \in C_0^\infty(\mathbb{R}^n)$. Let $u \in C^\infty((0, \infty) \times \mathbb{R}^n)$ solve the problem*

$$\begin{aligned} \partial_t^2 u - \Delta_g u &= f, \quad \text{in } (0, \infty) \times \mathbb{R}^n, \\ u|_{t=0} &= \partial_t u|_{t=0} = 0. \end{aligned}$$

Then there exists $t_0 > 0$ such that for all $t > t_0$

$$\|u(t, \cdot)\|_{L^2(K)} \leq C\mu(t) \|f\|_{L^2(\mathbb{R}^{1+n})},$$

where $\mu(t)$ is defined in (11). Here the constants C and t_0 depend on g , K and the support of f .

Proof. Let $t_1 > 0$ be such that $\text{supp}(f) \subset [0, t_1] \times \mathbb{R}^n$. By the finite speed of wave propagation, it holds that $\text{supp}(u|_{t=t_1})$ and $\text{supp}(\partial_t u|_{t=t_1})$ are compact in \mathbb{R}^n . Consider the solution v of the initial value problem

$$\begin{aligned} \partial_t^2 v - \Delta_g v &= 0, \quad \text{in } (t_1, \infty) \times \mathbb{R}^n, \\ v|_{t=t_1} &= u(t_1), \quad \partial_t v|_{t=t_1} = \partial_t u|_{t=t_1}. \end{aligned}$$

By the uniqueness, it must hold that $v = u$. By Lemma 3 there exists $t_0 > t_1$ and constant C independent of $t > t_0$ such that

$$\|u(t, \cdot)\|_{L^2(K)} \leq C\mu(t)E_0, \quad t > t_0,$$

Where $E_0 = \int_{\mathbb{R}^n} |\nabla u(t_1, \cdot)|^2 + |\partial_t u(t_1, \cdot)|^2 dx$. As u is an energy class solution of a wave equation with zero initial values, by the standard energy estimates for the wave equation it holds that

$$E_0 \leq C\|f\|_{L^2(\mathbb{R}^{1+n})}.$$

This proves the claim. \square

Lemma 5. *Let (\mathbb{R}^n, g) be as in Theorem 3. Let $S > 0$ and $f, h \in C_0^\infty((0, S) \times \mathbb{R}^n)$. It follows that*

$$\lim_{T \rightarrow \infty} \mathbb{E}\langle C_T, f \otimes h \rangle = \langle \kappa v^f, \kappa v^h \rangle_{L^2((-\infty, S) \times \mathbb{R}^n)}.$$

Proof. Here we will use notation $f^s(t, x) = f(t + s, x)$ for a time shift $s \in \mathbb{R}$. By the Lemma 2 and standard energy estimates, we have

$$\mathbb{E}\langle u^s, f \rangle^2 = \mathbb{E}\langle \mathbb{W}^s, \chi^s v^f \rangle^2 = \langle \chi^s v^f, \chi^s v^f \rangle \leq C\|f\|_{L^2(\mathbb{R}^{1+n})}^2, \quad s < T,$$

where the constant C depends on T . Therefore, we see that the mapping

$$(\omega, s) \rightarrow \langle u^s(\omega), f \rangle \langle u^s(\omega), h \rangle$$

is integrable on $\Omega \times (0, T)$ with respect to $\mathbb{P} \times dt$. In consequence, together with (7) the Fubini theorem yields

$$\mathbb{E}\langle C_T, f \otimes h \rangle = \frac{1}{T} \int_0^T \mathbb{E}\langle \chi^s \mathbb{W}^s, v^f \rangle \langle \chi^s \mathbb{W}^s, v^h \rangle ds = \frac{1}{T} \int_0^T \langle \chi^s v^f, \chi^s v^h \rangle ds.$$

For the time-shifted characteristic function we have

$$\chi^s(t, x) = \chi_0^s(t) \kappa(x) = \kappa(x) - (1 - \chi_0^s(t)) \kappa(x)$$

and $\text{supp}(1 - \chi_0^s) \subset (-\infty, 1 - s)$. By the local energy decay in Lemma 4, there is a constant $C > 0$ depending on g and the supports of κ and f such that

$$\|v^f\|_{L^2((-\infty, 1-s) \times \text{supp}(\kappa))} \leq C \left(\int_{-\infty}^{1-s} |t|^{-4n+2} dt \right)^{1/2} \|f\|_{L^2(\mathbb{R}^{1+n})} \leq C s^{-2n+\frac{3}{2}} \|f\|_{L^2(\mathbb{R}^{1+n})},$$

for large s . Hence we obtain

$$\mathbb{E}\langle C_T, f \otimes h \rangle = \langle \kappa v^f, \kappa v^h \rangle + \frac{1}{T} \int_0^T R(s) ds,$$

where

$$|R(s)| \leq C s^{-2n+\frac{3}{2}} \|f\|_{L^2(\mathbb{R}^{1+n})} \|h\|_{L^2(\mathbb{R}^{1+n})}$$

To conclude, one has

$$\frac{1}{T} \int_0^T |R(s)| ds \leq C T^{-2n+\frac{3}{2}}$$

and the claim follows. \square

In order to show the statistical stability of the data, we need the following result from ergodic theory (see e.g. [11, p. 94]):

Theorem 5. *Let \tilde{Z}_t , $t \geq 0$, be a real-valued random variables such that $\mathbb{E}\tilde{Z}_t = 0$ and the covariance function $(t, s) \mapsto \mathbb{E}(\tilde{Z}_t \tilde{Z}_s)$, $t, s \geq 0$, is continuous. Assume that for some constants $c, \epsilon > 0$ the condition*

$$|\mathbb{E}(\tilde{Z}_t \tilde{Z}_{t+r})| \leq c(1+r)^{-\epsilon}$$

holds for all $t \geq 0$ and $r \geq 0$. Then,

$$\lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T \tilde{Z}_t dt = 0 \quad \text{almost surely.}$$

Lemma 6. *Let (\mathbb{R}^n, g) be as in Theorem 3. Let $f, h \in C_0^\infty((0, S) \times \mathbb{R}^n)$ and use notation*

$$Z_r = \langle u^r, f \rangle \langle u^r, h \rangle$$

Then there is $C > 0$ depending on n, g , and the supports of κ, f and h such that

$$|\mathbb{E}(Z_r - \mathbb{E}Z_r)(Z_{r+s} - \mathbb{E}Z_{r+s})| \leq C(1+s)^{-n} \|f\|_{L^2(\mathbb{R}^{1+n})}^2 \|h\|_{L^2(\mathbb{R}^{1+n})}^2.$$

Proof. For convenience, let us write $X^r = \langle u^r, f \rangle$ and $Y^r = \langle u^r, h \rangle$. By the Isserlis formula [26] for Gaussian random variables we have

$$\mathbb{E}Z_r Z_{r+s} = \mathbb{E}X^r Y^r \mathbb{E}X^{r+s} Y^{r+s} + \mathbb{E}X^r X^{r+s} \mathbb{E}Y^r Y^{r+s} + \mathbb{E}X^r Y^{r+s} \mathbb{E}Y^r X^{r+s}$$

and, consequently,

$$(14) \quad \mathbb{E}(Z_r - \mathbb{E}Z_r)(Z_{r+s} - \mathbb{E}Z_{r+s}) = \mathbb{E}X^r X^{r+s} \mathbb{E}Y^r Y^{r+s} + \mathbb{E}X^r Y^{r+s} \mathbb{E}Y^r X^{r+s}.$$

We write $v_r^f(t, \cdot) = v^f(t - r, \cdot)$. The local energy decay, Lemma 4, implies

$$\begin{aligned}
|\mathbb{E}X^r Y^{r+s}| &= |\mathbb{E}\langle \chi^r \mathbb{W}^r, v^f \rangle \langle \chi^{r+s} \mathbb{W}^{r+s}, v^h \rangle| \\
&= |\mathbb{E}\langle \chi \mathbb{W}, v_r^f \rangle \langle \chi \mathbb{W}, v_{r+s}^h \rangle| \\
&= |\langle \chi v_r^f, \chi v_{r+s}^h \rangle| \\
(15) \qquad &\leq C(1+s)^{-2n+\frac{3}{2}} \|f\|_{L^2(\mathbb{R}^{1+n})} \|h\|_{L^2(\mathbb{R}^{1+n})},
\end{aligned}$$

since v_r^f is small in $\text{supp}(v_{r+s}^h)$ for $s \gg 0$. \square

Proof of Theorem 3. For a fixed pair of sources (f, h) we set $\tilde{Z}_t = Z_t - \mathbb{E}Z_t$, where $Z_t = \langle u^r, f \rangle \langle u^r, h \rangle$. Continuity of the covariance function of Z_t follows by considering equality (14). Note that the correlations between X^r, X^{r+s}, Y^r and Y^{r+s} on the right hand side of (14) are all represented by inner products between smooth functions in the spirit of (15). Since these inner products are smooth functions with respect to r and s , it follows that the covariance function in (14) is continuous. Next, we combine Lemma 5 and Lemma 6 to validate Theorem 5. As a countable set of source pairs (countable union of zero measurable sets is zero measurable, the claim follows for all $(f, h) \in \mathbb{D}$. \square

We conclude this section with the following simple lemma to quantify the convergence of the data. Notice that Lemma 7 is not needed for the previous proof.

Lemma 7. *Let $f, h \in C_0^\infty((0, S) \times \mathbb{R}^n)$. Then there is $C > 0$ depending on n, g , and the supports of κ, f and h such that*

$$\text{Var}\langle C_T, f \otimes h \rangle \leq CT^{-2} \|f\|_{L^2(\mathbb{R}^{1+n})}^2 \|h\|_{L^2(\mathbb{R}^{1+n})}^2$$

Proof. In the proof of Lemma 5 we showed that the Gaussian random variables X^r and Y^r have a bounded variance independent of r . Since any moment of a Gaussian random variable is bounded by a constant depending on the variance, we see that the mapping

$$(\omega, r, s) \rightarrow X^r Y^r X^s Y^s$$

is integrable over $\Omega \times (0, T) \times (0, T)$ for any fixed $T > 0$ with respect to $\mathbb{P} \times dr \times ds$.

Now the Fubini theorem yields that

$$\begin{aligned}
\mathbb{E}\langle C_T, f \otimes h \rangle^2 &= \frac{1}{T^2} \int_0^T \int_0^T \mathbb{E}X^s Y^s X^r Y^r ds dr \quad \text{and} \\
(\mathbb{E}\langle C_T, f \otimes h \rangle)^2 &= \frac{1}{T^2} \int_0^T \int_0^T \mathbb{E}X^s Y^s \mathbb{E}X^r Y^r ds dr.
\end{aligned}$$

It follows by equation (14) and estimate (15) that

$$\text{Var}(\langle C_T, f \otimes h \rangle) \leq C \|f\|_{L^2(\mathbb{R}^{1+n})}^2 \|h\|_{L^2(\mathbb{R}^{1+n})}^2 \frac{1}{T^2} \int_0^T \int_0^T (1 + |r - s|)^{-4n+3} ds dr$$

and the claim follows by estimating the the double integral in time by

$$\begin{aligned}
\int_0^T \int_0^T (1 + |r - s|)^{-4n+3} ds dr &= \frac{1}{2} \int_0^{2T} \int_{2T-s'}^{s'} (1 + r')^{-4n+4} dr' ds' \\
&= \frac{1}{2(1-n)} \int_0^{2T} ((1 + s')^{1-n} - (1 + 2T - s')^{1-n}) ds' \\
&\leq C(1 + T^{5-4n}) \leq C.
\end{aligned}$$

for any $n \geq 3$. □

4. REDUCTION TO THE DETERMINISTIC INVERSE PROBLEM

In this section we will show the following theorem.

Theorem 6. *Let $\mathbb{D} \subset C_0^\infty((0, \infty) \times \mathcal{X})$ be a dense set, and consider the data*

$$(16) \quad \langle \kappa v^f, \kappa v^h \rangle_{L^2(\mathbb{R}^{1+n})}, \quad f, h \in \mathbb{D},$$

where v^f and v^h solve (10) with the sources f and h , respectively. Then the data (16) determine the local source-to-solution map $\Lambda_{\mathcal{X}}$ as defined in Theorem 2.

It follows from the assumptions in Theorem 1 that the Riemannian manifold (\mathbb{R}^n, g) is complete. Indeed, the metric tensor g coincides with the Euclidean metric e outside a compact set, and therefore there exist uniform constants $c, C > 0$ such that $c\|\cdot\|_e \leq \|\cdot\|_g \leq C\|\cdot\|_e$, where $\|\cdot\|_e$ stands for the Euclidean and $\|\cdot\|_g$ for the Riemannian norm. Thus Theorems 3, 6 and 2 imply Theorem 1.

We will prove two auxiliary lemmas before presenting a proof for Theorem 6. Let d_0 be the Euclidean distance in \mathbb{R}^n , and denote by d_g the Riemannian distance in (\mathbb{R}^n, g) . For $p \in \mathbb{R}^n$ and $r > 0$, we denote the respective open balls by $B_0(p, r)$ and $B_g(p, r)$. We will use the shorthand notation $\square_g = \partial_t^2 - \Delta_g$.

Definition 2. *For $\mathcal{B} \subset (0, \infty) \times \mathbb{R}^n$, we say that $f \in C_0^\infty(\mathcal{B})$ is non-radiating, if $\text{supp}(w^f) \subset \overline{\mathcal{B}}$ for the solution $w = w^f$ of*

$$(17) \quad \begin{aligned} \partial_t^2 w - \Delta_g w &= f \quad \text{in } (0, \infty) \times \mathbb{R}^n, \\ w|_{t=0} &= \partial_t w|_{t=0} = 0. \end{aligned}$$

Furthermore, we define $\mathcal{N}(\mathcal{B}) = \{f \in C_0^\infty(\mathcal{B}) \mid f \text{ is non-radiating}\}$.

Definition 3. *We define the future of a set $\mathcal{B} \subset \mathbb{R}^{1+n}$ by*

$$\begin{aligned}
\mathcal{I}^+(\mathcal{B}) &= \{(t, x) \in \mathbb{R}^{1+n} \mid \text{there exists } (s, y) \in \mathcal{B} \text{ such that } t > s \\
&\quad \text{and } d_g(x, y) < t - s\}.
\end{aligned}$$

Lemma 8. *Let $(t_0, x_0) \in \mathbb{R} \times \mathcal{X}$, $\epsilon > 0$, and define $\mathcal{B} = (t_0 - \epsilon, t_0) \times B_0(x_0, \epsilon)$, and $\mathcal{Q} = (t_0, t_0 + 1) \times \mathcal{X}$. Let $f \in C_0^\infty(\mathcal{B})$. For small $\epsilon > 0$, $f \in \mathcal{N}(\mathcal{B})$ if and only if*

$$(18) \quad \langle \kappa w^f, \kappa w^h \rangle_{L^2(\mathbb{R} \times N)} = 0, \quad h \in C_0^\infty(\mathcal{Q}).$$

Proof. Let $\epsilon > 0$ be small enough so that

$$(19) \quad \mathcal{I}^+(\mathcal{B}) \cap (\{t_0\} \times \mathbb{R}^n) \subset \{t_0\} \times \mathcal{X}.$$

Clearly, $f \in \mathcal{N}(\mathcal{B})$ implies (18). Suppose now that (18) holds. Let $\phi \in C_0^\infty(\mathcal{Q})$. By choosing $h = \square_g(\kappa^{-2}\phi)$, we have $w^h = \kappa^{-2}\phi$ and further

$$\langle w^f, \phi \rangle_{L^2(\mathbb{R}^{1+n})} = \langle \kappa w^f, \kappa w^h \rangle_{L^2(\mathbb{R}^{1+n})} = 0.$$

Thus $w^f = 0$ in \mathcal{Q} . By (19) and the finite speed of wave propagation, it holds that $w^f = 0$ in $(t_0, \infty) \times \mathbb{R}^n$. Using the finite speed of wave propagation once more, we see that $w^f(t, x) = 0$ when $t \in \mathbb{R}$ and $d_g(x, B_0(x_0, \epsilon)) \geq \epsilon$. The exterior domain $E := \mathbb{R}^n \setminus B_0(x_0, \epsilon)$ is connected and $\partial_t^2 w^f - \Delta_g w^f = 0$ in $\mathbb{R} \times E$. Thus $w^f = 0$ in $\mathbb{R} \times E$ by unique continuation (Theorem 10 in the appendix). \square

Lemma 9. *Let $x_1, x_2 \in \mathcal{X}$, and let $\epsilon > 0$ be so small that $B_0(x_j, \epsilon) \subset \mathcal{X}$, $j = 1, 2$. Let $t_0 > 0$ and define $\mathcal{C} = (0, \infty) \times B_0(x_1, \epsilon)$ and $\mathcal{B} = (t_0 - \epsilon, t_0) \times B_0(x_2, \epsilon)$. Then*

$$(20) \quad \mathcal{I}^+(\mathcal{C}) \cap \mathcal{B} = \emptyset,$$

if and only if

$$(21) \quad \langle \kappa w^f, \kappa w^h \rangle_{L^2(\mathbb{R} \times N)} = 0, \quad h \in C_0^\infty(\mathcal{C}), \quad f \in \mathcal{N}(\mathcal{B}).$$

Proof. As f is non-radiating, the finite speed of wave propagation guarantees that (20) implies (21). Suppose now that (20) does not hold. The set $\mathcal{A} := \mathcal{I}^+(\mathcal{C}) \cap \mathcal{B}$ is open and non-empty. Let $\phi \in C_0^\infty(\mathcal{A})$ be non-zero and $\phi \geq 0$. Choose $(s, x) \in \mathcal{A}$ such that $\phi(s, x) > 0$. By approximate controllability (Theorem 11 in the appendix), there exists a source $h \in C_0^\infty(\mathcal{C})$ such that

$$\langle \phi(s), w^h(s) \rangle_{L^2(\mathbb{R}^n)} > 0.$$

Since w^h and ϕ are continuous, there is $\chi \in C_0^\infty(\mathbb{R})$ such that $\langle \chi \phi, w^h \rangle_{L^2(\mathbb{R}^{1+n})} > 0$. We define the function $f = \square_g(\kappa^{-2}\chi \phi) \in C_0^\infty(\mathcal{B})$. Then $f \in \mathcal{N}(\mathcal{B})$ and

$$\langle \kappa w^f, \kappa w^h \rangle_{L^2(\mathbb{R} \times N)} = \langle \chi \phi, w^h \rangle_{L^2(\mathbb{R}^{1+n})} > 0.$$

Therefore (21) is not valid either. \square

Now we are ready to present the proof of Theorem 6.

Proof of Theorem 6. The inner products (16) determine the same inner products for all $f, h \in C_0^\infty((0, \infty) \times \mathcal{X})$ by density. By reversing the time, these again determine the inner products

$$(22) \quad \langle \kappa w^f, \kappa w^h \rangle_{L^2(\mathbb{R}^{1+n})}, \quad f, h \in C_0^\infty((0, \infty) \times \mathcal{X}).$$

Let $x_1, x_2 \in \mathcal{X}$ and $\epsilon, t_0 > 0$ be as in Lemma 9. Observe that

$$(23) \quad d_g(B_0(x_1, \epsilon), B_0(x_2, \epsilon)) = \sup\{t_0 > 0 \mid (21) \text{ is valid}\}$$

and $d_g(x_1, x_2) = \lim_{\epsilon \rightarrow 0} d_g(B_0(x_1, \epsilon), B_0(x_2, \epsilon))$. For \mathcal{B} be as in Lemma 8, we can determine the set $\mathcal{N}(\mathcal{B})$, since the validity of (18) can be tested given the inner products (22). Thus also the validity of (21) can be tested given (22), and the distance function d_g can be determined on $\mathcal{X} \times \mathcal{X}$. These distances determine (\mathcal{X}, g) up to an isometry (see e.g. the proof of Proposition 5 below).

Let $h \in C_0^\infty((0, \infty) \times \mathcal{X})$ and let us show that $w^h|_{(0, \infty) \times \mathcal{X}}$ can be determined from the inner products (22). Let \mathcal{B} be as in Lemma 8. As $(0, \infty) \times \mathcal{X}$ can be covered with a countable number of sets of the form \mathcal{B} , it is enough to show that $w^h|_{\mathcal{B}}$ can be determined. We have already shown that $\mathcal{N}(\mathcal{B})$ can be determined given (22). Let $f \in \mathcal{N}(\mathcal{B})$. Then w^f is a solution of the following initial boundary value problem

$$(24) \quad \begin{aligned} \partial_t^2 w - \Delta_g w &= f \quad \text{in } (0, \infty) \times \mathcal{X}, \\ w|_{\mathbb{R} \times \partial \mathcal{X}} &= 0, \\ w|_{t=0} = \partial_t w|_{t=0} &= 0. \end{aligned}$$

As (\mathcal{X}, g) is known, we can solve the above equation. Thus for every $f \in \mathcal{N}(\mathcal{B})$ we are able to find w^f . In particular, in the inner products

$$(25) \quad \langle w^f, \kappa^2 w^h \rangle_{L^2((0, \infty) \times \mathcal{X})}, \quad f \in \mathcal{N}(\mathcal{B}),$$

the left factor w^f is known. Observe that for any $\phi \in C_0^\infty(\mathcal{B})$ we have $w^f = \phi$ where $f = \square_g \phi \in \mathcal{N}(\mathcal{B})$, and therefore the inclusion

$$(26) \quad \{w^f \mid f \in \mathcal{N}(\mathcal{B})\} \subset L^2(\mathcal{B})$$

is dense. Hence we find $\kappa^2 w^h|_{\mathcal{B}}$ from the inner products (25).

Let us conclude the proof by showing that function $\kappa|_{\mathcal{X}}$ can be determined. We have already shown that, when $f \in \mathcal{N}(\mathcal{B})$, both w^f and $\kappa^2 w^h|_{\mathcal{B}}$ are determined by (22). Thus $\kappa|_{\mathcal{X}}$ can be determined by the density (26). \square

5. THE DETERMINISTIC INVERSE PROBLEM

In this section we prove Theorem 2 in two steps: we show first that local the source-to-solution map $\Lambda_{\mathcal{X}}$ determines a certain family of distance functions, and then that this family determines the geometry g . We work first under the assumption that $d_g|_{\mathcal{X} \times \mathcal{X}}$ is known, and postpone the proof that $\Lambda_{\mathcal{X}}$ determines $d_g|_{\mathcal{X} \times \mathcal{X}}$ in the end of

the section. Recall that in the previous section we already determined $d_g|_{\mathcal{X} \times \mathcal{X}}$, so the step from $\Lambda_{\mathcal{X}}$ to $d_g|_{\mathcal{X} \times \mathcal{X}}$ is needed only in the proof of Theorem 2.

5.1. Reconstruction of a family of distance functions from the local source-to-solution mapping $\Lambda_{\mathcal{X}}$. Consider the following data:

$$(27) \quad (\mathcal{X}, g|_{\mathcal{X}}, d_g|_{\mathcal{X} \times \mathcal{X}}, \Lambda_{\mathcal{X}})$$

Here \mathcal{X} and $g|_{\mathcal{X}}$ stand for the assumption that the Riemannian structure of the open manifold \mathcal{X} is known. We show the following theorem.

Theorem 7. *Let (N, g) be a complete Riemannian manifold. Then the local source-to-solution data (27) determines the following family of distance functions*

$$(28) \quad R_{\mathcal{X}}(N) := \{d_g(x, \cdot)|_{\mathcal{X}} : x \in N\} \subset C(\mathcal{X}).$$

This is to be proved in several steps. Let $T, \epsilon > 0$. For each $r > \epsilon$ and $x \in N$ we define a set

$$S_{\epsilon}(x, r) := (T - (r - \epsilon), T) \times B(x, \epsilon)$$

We denote for any measurable $A \subset N$ the function space

$$L^2(A) := \{u \in L^2(N) : \text{supp}(U) \subset \bar{A}\}.$$

Recall that for any $T > 0$, $f \in C_0^{\infty}(\mathbb{R}_+ \times N)$ the solution $w^f(T, \cdot) \in L^2(N)$.

Lemma 10. *Let $x, y, z \in N$, $\epsilon > 0$ and $\ell_x, \ell_y, \ell_z > \epsilon$. Then the following are equivalent:*

$$(29) \quad B(x, \ell_x) \subset \overline{B(y, \ell_y) \cup B(z, \ell_z)}$$

$$(30) \quad \begin{aligned} & \text{For all } f \in C_0^{\infty}(S_{\epsilon}(x, \ell_x)) \text{ there exists } (f_j)_{j=1}^{\infty} \subset C_0^{\infty}(S_{\epsilon}(y, \ell_y) \cup S_{\epsilon}(z, \ell_z)) \\ & \text{such that } \|w^f(T) - w^{f_j}(T)\|_{L^2(N)} \xrightarrow{j \rightarrow \infty} 0. \end{aligned}$$

Denote by $\ell_{\epsilon} = \max((\ell_z - \epsilon), (\ell_y - \epsilon))$. Here we consider that for $f_j \in C_0^{\infty}(S_{\epsilon}(y, \ell_y) \cup S_{\epsilon}(z, \ell_z))$ function w^{f_j} is the solution of the following wave equation:

$$(31) \quad \begin{cases} \square_g w = f_j, \text{ in } N \times (T - \ell_{\epsilon}, T) \\ w|_{t=T-\ell_{\epsilon}} = \partial_t w|_{t=T-\ell_{\epsilon}} = 0. \end{cases}$$

Proof. Suppose that (29) is valid. Let $f \in C_0^{\infty}(S_{\epsilon}(x, \ell_x))$, then by the finite speed of wave propagation it holds that

$$\text{supp } w^f(T) \subset B(x, \ell_x) \subset \overline{B(y, \ell_y) \cup B(z, \ell_z)}$$

We denote $w^f(T) =: w^f$. Let χ be the characteristic function of the ball $B(y, \ell_y)$ and set $w_y^f = \chi w^f$ and $w_z^f = w^f - w_y^f$. Since the boundary of a geodesic ball is a set of measure zero (see [39]), it holds that $w_y^f \in L^2(B(y, \epsilon))$ and $w_z^f \in L^2(B(z, \epsilon))$. By approximate controllability there exist sequences $(f_y^j)_{j=1}^{\infty} \subset C_0^{\infty}(S_{\epsilon}(y, \ell_y))$ and

$(f_z^j)_{j=1}^\infty \subset C_0^\infty(S_\epsilon(z, \ell_z))$ such that sequences $(w_y^{f_y^j}(T))_{j=1}^\infty$ and $(w_z^{f_z^j}(T))_{j=1}^\infty$ converge to $w_y^f(T)$ and $w_z^f(T)$, respectively, in $L^2(N)$. Therefore sequence

$$f_j = f_y^j + f_z^j \in C_0^\infty(S_\epsilon(y, \ell_y) \cup S_\epsilon(z, \ell_z)), \quad j = 1, 2, \dots$$

satisfies (30).

Suppose that (29) is not valid. Then the open set

$$U := B(x, \ell_x) \setminus \overline{(B(y, \ell_y) \cup B(z, \ell_z))}$$

is not empty. By approximate controllability, we can choose $f \in C_0^\infty(S_\epsilon(x, \ell_x))$ such that $\|w^f(T)\|_{L^2(U)} > 0$. Then by finite speed of wave propagation it holds that

$$\inf\{\|w^f(T) - w^h(T)\|_{L^2(N)} : h \in C_0^\infty(S_\epsilon(y, \ell_y) \cup S_\epsilon(z, \ell_z))\} > 0.$$

Therefore (30) is not true. \square

For any point $(p, \xi) \in TM$, $\|\xi\|_g = 1$ we will denote the cut distance function

$$\tau(p, \xi) = \sup\{t > 0 : d_g(p, \gamma_{p, \xi}(t)) = t\}.$$

Let $\alpha, \beta : (0, 1) \rightarrow N$ be curves such that $\alpha(1) = \beta(0)$. Then we denote by $\alpha\beta$ the concatenated curve.

Lemma 11. *Let (N, g) be a complete Riemannian manifold. Let $x, y \in N$ and let $\gamma_{y, \xi}$ be a distance minimizing geodesic from y to x . Let $s := d_g(x, y)$. Let $r > 0$. If $\tau(y, \xi) < s + r$, then*

$$(32) \quad \text{there exists } \epsilon > 0 \text{ such that } B(x, r + \epsilon) \subset \overline{B(y, s + r)}.$$

Also if (32) is valid then $\tau(y, \xi) \leq s + r$.

Indeed

$$\tau(y, \xi) = \inf\{s + r > 0 : r, s > 0, \gamma_{y, \xi}([0, s]) \subset \mathcal{X}, (32) \text{ holds}\}.$$

Proof. Let $r > 0$ and denote $p = \gamma_{y, \xi}(s + r)$.

Suppose that (32) is valid. Let $\delta \in (0, \epsilon)$ and consider a point

$$z = \gamma_{y, \xi}(s + r + \delta) \in B(x, r + \epsilon).$$

By (32) $d_g(z, y) \leq s + r$. Thus $\tau(y, \xi) < s + r + \delta$. Since δ was arbitrary we have $\tau(y, \xi) \leq s + r$.

Suppose that $\tau(y, \xi) < s + r$. We show first that

$$(33) \quad \overline{B(x, y)} \subset B(y, s + r).$$

By triangle inequality it suffices to show that $\partial B(x, y) \subset B(y, s + r)$. Let $z \in \partial B(x, y)$. By triangle inequality $d_g(z, y) \leq s + r$. Let α be a minimizing geodesic from x to z . Suppose first that α is not the geodesic continuation of segment $\gamma_{y, \xi}([0, s])$. Since a curve $\gamma_{y, \xi}\alpha$ has a length $s + r$ and it is not smooth at x , it must hold

that $d_g(z, y) < s + r$. If α is the geodesic continuation of segment $\gamma_{y,\xi}((0, s))$, then $z = \gamma_{y,\xi}(s + r) = p$. Since $\tau(y, \xi) < s + r$, it holds that $d_g(y, p) < s + r$. Thus (33) follows. Therefore $\text{dist}_g(\partial B(x, r), \partial B(y, s + r)) > 0$ and (32) is valid. \square

Next we provide a method to find the cut distance function τ .

Proposition 1. *For any $y \in \mathcal{X}$ and $\xi \in S_y N$ we can find $\tau(y, \xi)$ from the local source-to-solution data (27).*

Proof. Let $y \in \mathcal{X}$ and $\xi \in S_y N$. Given the data (27) we can find the geodesic segment $\gamma_{y,\xi}([0, s])$ for small values $s > 0$.

Let $s > 0$ be so small that $\gamma_{y,\xi}([0, s]) \subset \mathcal{X}$. We denote $x = \gamma_{y,\xi}(s)$. Let $r > 0$. Consider the relation (32). By Lemma 11, relation (32) determines $\tau(y, \xi)$.

Choose $\epsilon > 0$ so small that

$$B(y, \epsilon) \cup B(x, \epsilon) \subset \mathcal{X}.$$

By taking $z = y$, $l_y = r + s = l_z$, $l_x = r + \epsilon$ as in Lemma 10 we see that (32) is equivalent with relation (30). Using the Blagovestchenskii identity (see (45) in the appendix) we see that the source-to-solution data (27) determines (30). \square

Lemma 12. *It holds that*

$$\{\gamma_{y,\xi}(t) \in N : y \in \mathcal{X}, \xi \in S_y N, t < \tau(y, \xi)\} = N.$$

Proof. Let $p \in N$ and choose any $y \in \mathcal{X}$. Let $\gamma_{y,\xi}$ be a distance minimizing geodesic from y to p . We denote by $r = d_g(y, p)$. Then it holds that $r \leq \tau(y, \xi)$. Choose $s \in (0, r)$ such that $y_1 := \gamma_{y,\xi}(s)$, $\gamma_{y,\xi}([0, s]) \subset \mathcal{X}$. Let $\xi_1 := \dot{\gamma}_{y,\xi}(s)$. We will show that $r - s < \tau(y_1, \xi_1)$ and this proves the claim of this lemma.

Suppose that $\tau(y_1, \xi_1) \leq r - s$. By the symmetry of cut points, it holds that $\tau(p, \eta) \leq r - s$, where $\eta := -\dot{\gamma}_{y,\xi}(r)$. Thus there exists $t \in (0, s)$ such that for a point $z := \gamma_{y,\xi}(t)$ it holds $d_g(p, z) < r - t$. Then it also holds that

$$d_g(y, p) \leq d_g(y, z) + d_g(z, p) < t + r - t = r.$$

This is a contradiction and therefore $r - s < \tau(y_1, \xi_1)$. \square

Notice that the assumption \mathcal{X} is open is crucial in Lemma 12. For instance consider the cylinder

$$\{e^{i\pi t} \in \mathbb{C} : t \in [-1, 1]\} \times (-1, 1),$$

and let $\mathcal{X} = \{1\} \times (-1, 1)$ and $p = (-1, 0)$. Then it holds that every point in \mathcal{X} is a cut point of p .

Proposition 2. *Let $z, y \in \mathcal{X}$, $\xi \in T_y \mathcal{X}$, $\|\eta\| = 1$ and $\tilde{r} < \tau(y, \eta)$. Then the local source-to-solution data (27) determines $d_g(p, z)$, where $p = \gamma_{y,\xi}(\tilde{r})$.*

Proof. Let $s \in (0, \tilde{r})$ be such that $\gamma_{y,\xi}([0, s]) \subset \mathcal{X}$. We denote by $x = \gamma_{y,\xi}(s)$. Let $r := \tilde{r} - s$.

Let $R > 0$. By Lemma 10 the inclusion

$$(34) \quad B(x, r + \epsilon) \subset \overline{B(y, r + s) \cup B(z, R)}$$

is valid for all $\epsilon > 0$ small enough if and only if the equation (30) is valid with $\ell_x = r + \epsilon$, $\ell_y = r + s$ and $\ell_z = R$. Using the Blagovestchenskii identity the local source-to-solution data (27) determines (30). We will show that

$$d_g(p, z) = R^* := \inf\{R > 0 : \text{Formula (34) is valid for } R \text{ and some } \epsilon > 0\}.$$

Suppose that (34) is valid. Since we assumed that $r + s < \tau(y, \xi)$, it holds that $p \in \overline{B(z, R)}$. Thus $d_g(p, z) \leq R^*$.

Suppose that $R \in (d_g(p, z), R^*)$. Then for any $\epsilon > 0$ (34) is not valid. Choose for every $k \in \mathbb{N}$ a point

$$p_k \in B(x, r + 1/k) \setminus \overline{B(y, r + s) \cup B(z, R)}.$$

By compactness of $\overline{B(x, r + 1)}$ we may assume that $p_k \rightarrow \tilde{p} \in \partial B(x, r)$ as $k \rightarrow \infty$. By similar argument as in the proof of Lemma 11 we deduce that $\tilde{p} = p$. Since $p \in B(z, R)$ we get a contradiction with the choice of sequence $(p_k)_{k=1}^\infty$. Therefore interval $(d_g(p, z), R^*) = \emptyset$ and $R^* = d_g(p, z)$. \square

Let $p \in N$ and $z \in \mathcal{X}$. By Lemma 12 it holds that there exists $y \in \mathcal{X}$ and an unit vector $\xi \in S_y N$ such that $p = \gamma_{y,\xi}(\tilde{r})$, for some $\tilde{r} < \tau(y, \xi)$. By Propositions 1 and 2 we have reconstructed $R_{\mathcal{X}}(N)$. Therefore Theorem 7 is proved.

5.2. Reconstruction of the Riemannian manifold from the distance functions. So far we have been able to find the following *distance data*

$$(35) \quad (\mathcal{X}, g|_{\mathcal{X}}, R_{\mathcal{X}}(N)),$$

where $R_{\mathcal{X}}(N)$ is defined by (28). In this section we will show, how one can reconstruct the topological, smooth and Riemannian structures from the distance data (35). The rest of the paper is devoted to showing the following theorem:

Theorem 8. *Let (N, g) be a complete smooth Riemannian manifold without a boundary. Let $U \subset N$ be open, bounded and have a smooth boundary. Suppose that the topological and smooth structure of U are known, and $g|_U$ is also known. Then*

$$R(N) := \{d_g(\cdot, x)|_{\overline{U}} : x \in N\} \subset C(\overline{U})$$

determines, topological, smooth and Riemannian structure of N up to isometry.

Since \bar{U} is compact, $C(\bar{U})$ is a Banach space when equipped with L^∞ -norm. We define the mapping

$$R : N \rightarrow C(\bar{U}), R(x) = r_x = d_g(x, \cdot)|_{\bar{U}}.$$

Our aim is to construct such a Riemannian structure in $R(N) \subset C(\bar{U})$ that $R : N \rightarrow R(N)$ is a Riemannian isometry.

Lemma 13. *Mapping R is continuous and one-to-one.*

Proof. Let $x, y \in N$. Then by the triangle inequality

$$\|R(x) - R(y)\|_{L^\infty(\bar{U})} = \sup_{z \in \bar{U}} |r_x(z) - r_y(z)| \leq d_g(x, y).$$

Thus R is continuous.

Suppose that $x, y \in N$ satisfy $r_x = r_y$. If $x \in \bar{U}$ then $r_y(x) = 0$ and thus $x = y$. Therefore we can assume that $x, y \in N \setminus \bar{U}$. Since \bar{U} is compact there exists a closest point $z \in \bar{U}$ to x . Then $z \in \partial U$ and it is also a closest point of \bar{U} to y . Since ∂U is smooth $n - 1$ dimensional submanifold of N , the distance minimizing unit speed geodesic γ from z to x is orthogonal to ∂U . Since both x and y are points of the exterior of U , it holds by the uniqueness of geodesics that

$$x = \gamma(r_x(z)) = \gamma(r_y(z)) = y.$$

This completes the proof. \square

Next we will recall two topological results that allow us to prove that mapping $R : N \rightarrow R(N)$ is a homeomorphism.

Definition 4. *Let X be a topological space. We say that a sequence $(x_j)_{j=1}^\infty$ in X escapes to infinity, if for every compact $K \subset X$, $x_j \in K$ for at most finitely many $j \in \mathbb{N}$.*

For the proofs of the following two lemmas see for instance [51].

Lemma 14. *Let (X, d_X) and (Y, d_Y) be metric spaces. Let $f : X \rightarrow Y$ be continuous. Then f is proper if and only if for every sequence $(x_j)_{j=1}^\infty \subset X$ that escapes to infinity the image sequence $(f(x_j))_{j=1}^\infty \subset Y$ escapes to infinity.*

Lemma 15. *Let (X, d_X) and (Y, d_Y) be metric spaces. Let $f : X \rightarrow Y$ be one-to-one, continuous and proper. Then mapping the f is closed.*

Proposition 3. *Mapping $R : N \rightarrow R(N)$ is a homeomorphism.*

Proof. If N is bounded, then N is compact, and the claim follows from basic topology. Suppose that N is not bounded. Let $(x_j)_{j=1}^\infty \subset N$ be a sequence that escapes to

infinity. Let $x_0 \in \bar{U}$. We define $X_j := \overline{B(x_0, j)}$ for every $j \in \mathbb{N}$ and $Y_j = R(X_j)$. Then $\cup_{j=1}^{\infty} X_j = N$ and thus

$$\lim_{j \rightarrow \infty} d_g(x_0, x_j) = \infty.$$

We write $R(x_0) =: r_0$ and $R(x_j) =: r_j$. Then

$$d_{\infty}(r_0, r_j) \geq |d_g(x_0, x_0) - d_g(x_0, x_j)| = d_g(x_0, x_j).$$

Thus $d_{\infty}(r_0, r_j) \rightarrow \infty$ as $j \rightarrow \infty$. Since a compact set of a metric space is always bounded, it holds that sequence $(r_j)_{j=1}^{\infty}$ escapes to infinity. Therefore R is a proper mapping and by Lemma 15 it is closed. \square

By Proposition 3, the topological structure of N has been found. Next we will show, how to construct such a smooth atlases on N and $R(N)$ that the mapping R is a diffeomorphism.

Let $z \in \bar{U}$ and $x \in N$. Denote by $\omega(x)$ the cut locus of x . Recall that $r_x := d_g(x, \cdot)|_U$ is smooth at z if and only if $z \neq x$ or $z \notin \omega(x)$ (see Lemma 2.1.11 and Theorem 2.1.14 of [30]). Using also the fact that $z \in \omega(x)$ if and only if $x \in \omega(z)$ we can find the cut locus $\omega(z)$ from data (35). We denote by

$$\omega^T(z) = \exp_z^{-1} \omega(z)$$

the tangential cut locus of z . Notice that $\omega(z)$ and $\omega^T(z)$ are closed.

We define a mapping Φ_z by

$$\Phi_z(r) := -r(z)\nabla_g r|_z \in T_z(N), \quad r \in R(N) \setminus R(\omega(z)).$$

By the following lemma it holds

$$(36) \quad \Phi_z \circ R|_{N \setminus \omega(z)} = \exp_z^{-1},$$

where \exp_z is the exponential mapping of z .

Lemma 16. *Let $x \in N$. We denote $B := T_z N \setminus \omega^T(z)$. Then the following are equivalent:*

$$(37) \quad \eta \in B \text{ and } \exp_z(\eta) = x$$

$$(38) \quad \nabla_g d_g(x, \cdot)|_z \in T_z N \text{ exists and } \eta = -d_g(x, z)\nabla_g d_g(x, \cdot)|_z.$$

Proof. Suppose that formula (37) is valid. Since exponential mapping \exp_z is one-to-one in the set B , the point z is not in the cut locus of x and therefore the function $d_g(x, \cdot)$ is smooth at z . Thus $\nabla_g d_g(x, \cdot)|_z \in T_z N$ exists and $\eta = -d_g(x, z)\nabla_g d_g(x, \cdot)|_z$. Therefore (38) is also valid.

Suppose that formula (38) is valid. Then it holds that $d_g(x, \cdot)$ is smooth at z . Thus x is not in the cut locus of z and therefore $\xi := -\nabla_g d_g(x, \cdot)|_z$ is the initial

velocity of the unique distance minimizing geodesic from z to x . We have

$$\exp_z(\eta) = \gamma_{z,\xi}(d_g(x, z)) = x \in \exp_z(B).$$

□

We define the smooth structure on $R(N)$ by using mappings $\Phi_z, z \in \bar{U}$. By (36) each mapping Φ_z is a topological coordinate mapping. Let $z, w \in \bar{U}$. Then the composition

$$\Phi_z \circ \Phi_w^{-1} = (\Phi_z \circ R) \circ (\Phi_w \circ R)^{-1} = \exp_z^{-1} \circ \exp_w$$

is smooth whenever defined. Moreover, R is clearly smooth when the smooth structure of $R(N)$ is defined in this way. Therefore we have proved the following proposition.

Proposition 4. *The mapping $R : N \rightarrow R(N)$ is a diffeomorphism.*

We define a metric tensor $\tilde{g} := (R^{-1})^*g$ on $R(N)$, that is, \tilde{g} is the push forward of g . Then $(R(N), \tilde{g})$ and (N, g) are Riemannian isometric. In the next proposition, we provide a method to construct representation of \tilde{g} in local coordinates of $R(N)$.

Proposition 5. *Let $\tilde{g} := (R^{-1})^*g$. We can construct the metric tensor \tilde{g} on $R(N)$ from the distance data (35).*

Proof. Let $r_0 \in R(N)$. We write $x_0 := R^{-1}(r_0)$. By Lemma 12 it holds that there exists a point $z \in U$ that is not in the cut locus of x_0 . Let $U' \subset U$ be an open neighborhood of z such that $d_g(\cdot, y)$ is smooth at x_0 for any $y \in U'$.

It holds that

$$\nabla_g d_g(\cdot, y)|_{x_0} = -\dot{\gamma}_{y, x_0}(d_g(y, x_0)) \in S_{x_0}N,$$

where γ_{y, x_0} is the unique unit speed distance minimizing geodesic from y to x_0 . Since U' is open and \exp_{x_0} is continuous the set $\exp_{x_0}^{-1}U' \subset T_{x_0}N$ is open. Therefore the set

$$\mathcal{V} := \{\nabla_g d_g(\cdot, y)|_{x_0} \in S_{x_0}N : y \in U'\}$$

is open in $S_{x_0}N$. Let $(x, \xi) \in TN$. We will use the notation

$$\xi^b := \langle \xi, \cdot \rangle_g \in T_x^*N.$$

Since R is a diffeomorphism it holds that

$$\mathcal{W}^* := R_*\mathcal{V}^* = \{(\nabla d_g(R^{-1}(\cdot), y)|_{r_0})^b \in S_{r_0}^*R(N) : y \in U'\}$$

is open. For any point $y \in U'$ we define an evaluation function $E_y : R(N) \rightarrow \mathbb{R}$ with the formula $E_y(r) = r(y)$. Notice that

$$dE_y|_{r_0} = (\nabla d_g(R^{-1}(\cdot), y)|_{r_0})^b,$$

and therefore

$$\mathcal{W}^* = \{dE_y|_{r_0} \in S_{r_0}^*R(N) : y \in U'\}.$$

As we know the smooth structure of $R(N)$ we can find the set \mathcal{W}^* . The last step is to show that set \mathcal{W}^* determines $\tilde{g}(r_0)$.

Let

$$\mathbb{R}_+\mathcal{W}^* := \{sv \in T_{r_0}^*R(N) : v \in \mathcal{W}^*, s > 0\}$$

be the open cone generated by \mathcal{W}^* . Let $\{E_j\}_{j=1}^n$ be a local coordinate system at r_0 . For any $s > 0$ and $v \in \mathcal{W}^*$ it holds in coordinates $\{E_j\}_{j=1}^n$ that

$$F(sv) := s^2 \tilde{g}^{ij}(r_0) v_i v_j = s^2.$$

We know the function $F : \mathbb{R}_+\mathcal{W}^* \rightarrow \mathbb{R}$, and $\mathbb{R}_+\mathcal{W}^*$ is open, we get

$$\tilde{g}^{ij}(r_0) = \frac{\partial}{\partial E_i} \frac{\partial}{\partial E_j} F.$$

□

By Propositions 3, 4 and 5 we can reconstruct $(R(N), \tilde{g})$, more over (N, g) and $(R(N), \tilde{g})$ are isometric as Riemannian manifolds. Thus we have proved Theorem 8.

In order to prove Theorem 2 we still need the next small lemma.

Lemma 17. *Let (N, g) and \mathcal{X} be as in the formulation of Theorem 2. Then data $\{\mathcal{X}, \Lambda_{\mathcal{X}}\}$ determines the distance function d_g on $\mathcal{X} \times \mathcal{X}$.*

Proof. Let $x, y \in \mathcal{X}$. Since \mathcal{X} is a smooth manifold, we may choose an auxiliary metric d_0 on \mathcal{X} that gives the same topology as g . Let $\epsilon > 0$ and consider the metric ball $B_{d_0}(x, \epsilon)$. We write $\mathcal{B}_\epsilon := (0, \infty) \times B_{d_0}(x, \epsilon)$ and

$$t_\epsilon = \inf\{t > 0 : \text{there is } f \in C_0^\infty(\mathcal{B}_\epsilon) \text{ such that } \text{supp}(\Lambda_{\mathcal{X}} f)(t, \cdot) \cap B_{d_0}(y, \epsilon) \neq \emptyset\}.$$

By the finite speed of wave propagation and the approximate controllability the equality

$$t_\epsilon = \text{dist}_g(B_{d_0}(x, \epsilon), B_{d_0}(y, \epsilon))$$

holds. Thus the following limit is valid

$$d_g(x, y) = \lim_{\epsilon \rightarrow 0} t_\epsilon.$$

□

Now we are finally ready to give a proof for Theorem 2.

Proof of Theorem 2. By making \mathcal{X}_i smaller, if needed, we may assume without lost of generality that \mathcal{X}_i is precompact with smooth boundary and that $\phi : \overline{\mathcal{X}}_1 \rightarrow \overline{\mathcal{X}}_2$ (see (4)) is a diffeomorphism. Denote $R(N_i) = \{d_i(x, \cdot)|_{\overline{\mathcal{X}}_i} : x \in N_i\}$ and consider a mapping

$$R_i : N_i \rightarrow R(N_i), \quad i = 1, 2, \quad R_i(x) = d_i(x, \cdot)|_{\overline{\mathcal{X}}_i}.$$

By Lemma 17, Proposition 5 and equation (5) it holds that

$$(39) \quad d_1(\cdot, \cdot)|_{\mathcal{X}_1 \times \mathcal{X}_1} = d_2(\phi(\cdot), \phi(\cdot))|_{\mathcal{X}_1 \times \mathcal{X}_1} \text{ and } g_1|_{\mathcal{X}_1} = \phi^* g_2|_{\mathcal{X}_2}.$$

Therefore we may assume that $\phi : \overline{\mathcal{X}}_1 \rightarrow \overline{\mathcal{X}}_2$ is a Riemannian isometry. By Proposition 11 the following relation

$$\tau_1(y, \xi) = \tau_2(\phi(y), \phi(\xi)), \quad y \in \mathcal{X}_1, \quad \xi \in S_y N_1$$

is valid. Therefore by Proposition 2 it holds that

$$R_2(N_2) = \Phi(R_1(N_1)),$$

where

$$\Phi : C(\overline{\mathcal{X}}_1) \rightarrow C(\overline{\mathcal{X}}_2), \quad \Phi(f) = f \circ \phi^{-1}.$$

Moreover by Theorems 7 and 8 the mappings $R_i : N_i \rightarrow R_i(N_i)$ are Riemannian isometries.

With out lost of generality we assume that $\overline{\mathcal{X}}_1 \subset V$, where (V, α) is a coordinate chart for N_1 . Write $\alpha \circ \phi =: \tilde{\alpha}$, $W = \alpha(V)$ and define Riemannian isometries

$$\alpha^* : R(N_1) \rightarrow \alpha^*(R_1(N_1)) \subset C(W), \quad \alpha^*(r)(x) = r(\alpha^{-1}(x))$$

and

$$\tilde{\alpha}^* : R(N_2) \rightarrow \tilde{\alpha}^*(R_2(N_2)) \subset C(W), \quad \tilde{\alpha}^*(r)(y) = r(\tilde{\alpha}^{-1}(y)).$$

Thus we have proved that mapping

$$N_1 \xrightarrow{R_1} R_1(N_1) \xrightarrow{\alpha^*} \alpha^*(R_1(N_1)) \xrightarrow{id} \tilde{\alpha}^*(R_2(N_2)) \xrightarrow{(\tilde{\alpha}^*)^{-1}} R_2(N_2) \xrightarrow{R_2^{-1}} N_2,$$

is a Riemannian isometry. This ends the proof. \square

Lastly we will give a proof for Corollary 2.

Proof of Corollary 2. Since N is a compact manifold without a boundary we have

$$0 = \lambda_1 < \lambda_2 \leq \lambda_3 \dots$$

Let $f \in C_0^\infty((0, \infty) \times \mathcal{X})$ and $w = w^f$ be the solution of the initial value problem (3). For each $j \in \mathbb{N}$ we define the j^{th} Fourier coefficient

$$I_j(t) = \langle w(t, \cdot), \varphi_j \rangle_{L^2(N)}.$$

Since w is smooth, also I_j is smooth. By Greens formula and the initial conditions of (3) it holds that

$$(40) \quad \begin{cases} \frac{d^2}{dt^2} I_j(t) - \lambda_j I_j(t) = \int_{\mathcal{X}} f(t, x) \varphi_j(x) dV_g(x) \\ I_j(0) = \frac{d}{dt} I_j(0) = 0. \end{cases}$$

Solve the ordinary differential equation (40) to get

$$I_j(t) = \int_0^t \int_{\mathcal{X}} s_j(t-s) f(s, x) \varphi_j(x) dV_g(x) ds, \quad j \geq 1$$

where

$$s_1(t) = t \text{ and } s_j(t) = \frac{\sin(\sqrt{\lambda_j}(t))}{\sqrt{\lambda_j}}, \text{ for } j > 1.$$

Notice that apriori the volume form $dV_g|_{\mathcal{X}}$ is not given. However without a loss of generality we may assume that \mathcal{X} is contained in a coordinate patch of N . Thus we can assume, that we are given some volume form ω on \mathcal{X} . Therefore there exists a unique smooth function $\eta : \mathcal{X} \rightarrow (0, \infty)$ such that

$$\eta dV_g|_{\mathcal{X}} = \omega.$$

We write

$$\tilde{I}_j(t) = \int_0^t \int_{\mathcal{X}} s_j(t-s) f(t, x) \varphi_j(x) \omega(x) ds.$$

By direct computations and initial values of (40) we have

$$(41) \quad \sum_{j=1}^{\infty} \tilde{I}_j(t) \varphi_j(x) = w^{\eta f}(t, x).$$

Thus for every $f \in C_0^\infty((0, \infty) \times \mathcal{X})$ the Fourier coefficients $\tilde{I}_j(t)$ can be recovered from the Spectral data (6). We conclude that we have recovered the mapping

$$\sum_{j=1}^{\infty} \tilde{I}_j(t) \varphi_j(x)|_{x \in \mathcal{X}} = w^{\eta f}(t, x)|_{x \in \mathcal{X}} = (\Lambda_{\mathcal{X}} M_{\eta}) f.$$

Here M_{η} is the multiplier operator $M_{\eta} f(t, x) = \eta(x) f(t, x)$. Let $R(h(t, x)) = h(-t, x)$. Then

$$(\Lambda_{\mathcal{X}} M_{\eta})^* = M_{\eta}^* \Lambda_{\mathcal{X}}^* = M_{\eta} R \Lambda_{\mathcal{X}} R = R(M_{\eta} \Lambda_{\mathcal{X}}) R$$

(see Lemma 18 in the appendix), so that we have recovered the operator $M_{\eta} \Lambda_{\mathcal{X}}$. Therefore the claim follows from Theorem 2. \square

6. APPENDIX

We will use the assumptions and notations of Theorem 2. Let $T > 0$, $p \in N$ and $a > 1$. Let $C_{p,T}$ be the cone

$$C_{p,T} := \{(t, q) \in \mathbb{R} \times N : 0 \leq t \leq T, d_N(p, q) \leq T - t\}.$$

Theorem 9 (Finite speed of propagation). *Let $f \in L^2(\mathbb{R} \times N)$. Suppose that u solves*

$$\begin{cases} (\partial_t^2 - \Delta_g)u = f, & \text{in } (0, \infty) \times N \\ f|_{C_{p,T}} = 0 \\ u|_{B(p,T) \times \{t=0\}} = \partial_t u|_{B(p,T) \times \{t=0\}} = 0, \end{cases}$$

Then

$$u|_{C_{p,T}} = 0.$$

Proof. See [46]. □

Consider an open double cone created by a cylindrical set $(0, 2T) \times \mathcal{X}$

$$C(T, \mathcal{X}) = \left\{ (x, t) \in (0, 2T) \times N : \begin{array}{l} \text{dist}_g(x, \mathcal{X}) < t, t \leq T \\ \text{dist}_g(x, \mathcal{X}) < 2T - t, t \geq T \end{array} \right\}$$

We write

$$M(T, \mathcal{X}) = \{x \in N : \text{dist}_g(x, \mathcal{X}) \leq T\},$$

for the domain of influence of set \mathcal{X} .

Theorem 10 (Tataru's unique continuation). *Let $\mathcal{X} \subset N$ be open and bounded. Let $u \in C_0^\infty(\mathbb{R} \times N)$. Suppose that $(\partial_t^2 - \Delta_g)u = 0$ in $(0, 2T) \times M(T, \mathcal{X})$ and $u|_{(0,2T) \times \mathcal{X}} \equiv 0$. Then $u|_{C(T, \mathcal{X})} \equiv 0$.*

Proof. See [28] for a local result and [45] for the global result. □

We use a short hand notation

$$\mathcal{F}_{\mathcal{X}, T} := \{f \in C_0^\infty(\mathbb{R} \times N) : \text{supp } f \subset (0, T) \times \mathcal{X}\}.$$

Theorem 11 (Approximate controllability). *Let $\mathcal{X} \subset N$ be open and bounded. For any $T > 0$ set*

$$\mathcal{W}_T := \{w^f(T) : f \in \mathcal{F}_{\mathcal{X}, T}\}$$

is dense in Hilbert space $L^2(M(T, \mathcal{X}))$.

Proof. By the finite speed of wave propagation $\mathcal{W}_T \subset L^2(M(T, \mathcal{X}))$. Since $L^2(M(T, \mathcal{X}))$ is a Hilbert space, it suffices to prove that $\mathcal{W}_T^\perp = \{0\}$. Suppose that $\phi \in L^2(M(T, \mathcal{X}))$ is such that $(w^f(T), \phi)_{L^2(N)} = 0$ for all $f \in \mathcal{F}_{\mathcal{X}, T}$. Let $u \in C^\infty(\mathbb{R} \times N)$ solve

$$(42) \quad \begin{cases} (\partial_t^2 - \Delta_g)u = 0, & \text{in } (0, T) \times N \\ u|_{t=T} = 0, \quad \partial_t u|_{t=T} = \phi. \end{cases}$$

Let $f \in \mathcal{F}_{B, T}$. By the finite speed of wave propagation, there exists a compact set of N that contains the $\text{supp } w^f(t)$ for each $t \in (0, T)$. We use the Green identities to see that

$$\langle f, u \rangle_{L^2((0, T) \times N)} = \langle \square_g w^f, u \rangle_{L^2((0, T) \times N)} - \langle w^f, \square_g u \rangle_{L^2((0, T) \times N)} = 0.$$

Since $\mathcal{F}_{\mathcal{X},T}$ is dense in $L^2((0,T) \times \mathcal{X})$, it holds that $u \equiv 0$ in $(0,T] \times \mathcal{X}$.

Let U solve

$$(43) \quad \begin{cases} (\partial_t^2 - \Delta_g)U = 0, & \text{in } (0, 2T) \times N \\ U|_{t=0} = u(0), \quad \partial_t U|_{t=0} = \partial_t u|_{t=0}. \end{cases}$$

By equations (42) and (43) it holds $U|_{[0,T] \times N} = u$. More over the function $\tilde{u}(t, x) = -u(2T - t, x)$ solves the wave equation

$$(44) \quad \begin{cases} (\partial_t^2 - \Delta_g)\tilde{u} = 0, & \text{in } (T, 2T) \times N \\ \tilde{u}|_{t=T} = 0, \quad \partial_t \tilde{u}|_{t=T} = \phi, \end{cases}$$

since $\tilde{u}(T, x) = -u(2T - T, x) = 0$ and $\partial_t \tilde{u}|_{t=T} = \partial_t u(2T - T) = \phi$. Therefore in particular $U|_{(0,2T) \times \mathcal{X}} \equiv 0$.

By unique continuation (Theorem 10), it holds that $U|_{C(T,\mathcal{X})} \equiv 0$. Since $M(T, \mathcal{X}) \times \{T\} \subset C(T, \mathcal{X})$ we have

$$\phi|_{M(T,\mathcal{X})} = \partial_t U|_{t=T}|_{M(T,\mathcal{X})} = 0.$$

□

Next our aim is to prove the Blagovestchenskii identity on a complete Riemannian manifold (N, g) .

Theorem 12. *Let (N, g) be a complete Riemannian manifold. Let $T > 0$, $\mathcal{X} \subset N$ be open and bounded. Let $f, h \in \mathcal{F}_{B,2T}$, then*

$$(45) \quad \langle w^f(T), w^h(T) \rangle_{L^2(N)} = \langle f(T), (J\Lambda_{\mathcal{X}} - \Lambda_{\mathcal{X}}^* J)h(T) \rangle_{L^2((0,T) \times N)}$$

where the operator $J : L^2(0, 2T) \rightarrow L^2(0, T)$ is defined as

$$J\phi(t) = \frac{1}{2} \int_t^{2T-t} \phi(s) ds.$$

Proof. Let $f, h \in \mathcal{F}_{B,2T}$ and consider the mapping $W : [0, 2T] \times [0, 2T] \rightarrow \mathbb{R}$,

$$W(t, s) = \langle w^f(t), w^h(s) \rangle_{L^2(N)}.$$

Then using Greens formula

$$\begin{aligned} (\partial_t^2 - \partial_s^2)W(t, s) &= (\partial_t^2 - \partial_s^2)\langle w^f(t), w^h(s) \rangle_{L^2(N)} \\ &= \langle f(t), \Lambda_{B,2T}h(s) \rangle_{L^2(N)} - \langle \Lambda_{B,2T}f(t), h(s) \rangle_{L^2(N)} := F(t, s). \end{aligned}$$

Notice that there is no boundary terms due finite speed of wave propagation. The function $(t, s) \mapsto F(t, s)$ can be computed, if the local source-to-solution mapping $\Lambda_{\mathcal{X}}$ is given. By (3) it holds that

$$W(0, s) = 0 = \partial_t W(t, s)|_{t=0}.$$

Thus w is the solution of the following $(1 + 1)$ -dimensional initial value problem:

$$(46) \quad \begin{cases} (\partial_t^2 - \partial_s^2)W = F, & \text{in } (0, 2T) \times \mathbb{R} \\ W|_{t=0} = \partial_t W|_{t=0} = 0. \end{cases}$$

Recall that the following formula

$$(47) \quad W(t, s) = \frac{1}{2} \int_0^t \int_{s-\tau}^{s+\tau} F(t - \tau, y) \, dy d\tau, \quad s \in \mathbb{R}, t \in [0, 2T],$$

solves (46) (see e.q. [15]). By the change of variables $T - s = \tau$, we conclude

$$\begin{aligned} W(T, T) &= \frac{1}{2} \int_0^T \int_\tau^{2T-\tau} F(\tau, y) \, dy d\tau. \\ &= \langle f, J\Lambda_{\mathcal{X}}h \rangle_{L^2(\mathcal{X} \times (0, T))} - \langle \Lambda_{\mathcal{X}}f, Jh \rangle_{L^2(\mathcal{X} \times (0, T))}. \end{aligned}$$

□

Lemma 18. *The adjoint mapping of $\Lambda_{\mathcal{X}}$ in $L^2((0, T) \times \mathcal{X})$ is $R\Lambda_{\mathcal{X}}R$, where*

$$Rh(t, x) = h(T - t, x).$$

Proof. Let $f, h \in \mathcal{F}_{\mathcal{X}, T}$ and consider the wave equations

$$(48) \quad \begin{cases} (\partial_t^2 - \Delta_g)w = f, & \text{in } (0, T) \times N \\ w|_{t=0} = \partial_t w|_{t=0} = 0 \end{cases} \quad \text{and} \quad \begin{cases} (\partial_t^2 - \Delta_g)u = h, & \text{in } (0, T) \times N \\ u|_{t=T} = \partial_t u|_{t=T} = 0. \end{cases}$$

We start with observing that

$$\langle f, u \rangle_{L^2((0, T) \times N)} - \langle w, h \rangle_{L^2((0, T) \times N)} = 0.$$

This holds due the computations we have done in the proof of Theorem 11. Therefore

$$\langle f, u \rangle_{L^2((0, T) \times \mathcal{X})} - \langle \Lambda_{\mathcal{X}}f, h \rangle_{L^2((0, T) \times \mathcal{X})} = 0 \quad \text{and} \quad \Lambda_{\mathcal{X}}^*h = u|_{(0, T) \times \mathcal{X}}.$$

Replace $f = Rh$. Then

$$\square Ru = \square u(T - \cdot, \cdot) = h(T - \cdot, \cdot) = Rh \quad \text{and} \quad Ru(0, \cdot) = \partial_t Ru(0, \cdot) = 0.$$

By (48) it holds that

$$Ru|_{(0, T) \times \mathcal{X}} = w|_{(0, T) \times \mathcal{X}} = \Lambda_{\mathcal{X}}f = \Lambda_{\mathcal{X}}Rh.$$

Since $R \circ R = id_{L^2((0, T) \times \mathcal{X})}$ we get

$$u|_{(0, T) \times \mathcal{X}} = R\Lambda_{\mathcal{X}}Rh.$$

□

REFERENCES

- [1] M. Anderson, A. Katsuda, Y. Kurylev, M. Lassas, and M. Taylor. Boundary regularity for the ricci equation, geometric convergence, and gelfands inverse boundary problem. *Inventiones mathematicae*, 158(2):261–321, 2004.
- [2] G. Bal and L. Ryzhik. Time reversal and refocusing in random media. *SIAM Journal on Applied Mathematics*, 63(5):1475–1498, 2003.
- [3] M. Belishev and V. Y. Gotlib. Dynamical variant of the bc-method: theory and numerical testing. *Journal of Inverse and Ill-Posed Problems*, 7(3):221–240, 1999.
- [4] M. I. Belishev. An approach to multidimensional inverse problems for the wave equation. In *Dokl. Akad. Nauk SSSR*, volume 297, pages 524–527, 1987.
- [5] M. I. Belishev and Y. V. Kuryiev. To the reconstruction of a riemannian manifold via its spectral data (bc–method). *Communications in partial differential equations*, 17(5–6):767–804, 1992.
- [6] V. I. Bogachev and V. I. Bogachev. *Gaussian measures*, volume 62. American Mathematical Society Providence, 1998.
- [7] L. Borcea, G. Papanicolaou, and C. Tsogka. Theory and applications of time reversal and interferometric imaging. *Inverse Problems*, 19(6):S139, 2003.
- [8] L. Borcea, G. Papanicolaou, C. Tsogka, and J. Berryman. Imaging and time reversal in random media. *Inverse Problems*, 18(5):1247, 2002.
- [9] P. Caro, T. Helin, and M. Lassas. Inverse scattering for a random potential. *arXiv preprint arXiv:1605.08710*, 2016.
- [10] D. Conus, R. C. Dalang, et al. The non-linear stochastic wave equation in high dimensions. *Electron. J. Probab*, 13(22):629–670, 2008.
- [11] H. Cramér and M. R. Leadbetter. *Stationary and related stochastic processes. Sample function properties and their applications*. John Wiley & Sons, Inc., New York-London-Sydney, 1967.
- [12] M. V. de Hoop, J. Garnier, S. F. Holman, and K. Sølna. Retrieval of a green’s function with reflections from partly coherent waves generated by a wave packet using cross correlations. *SIAM Journal on Applied Mathematics*, 73(1):493–522, 2013.
- [13] M. V. de Hoop, P. Kepley, and L. Oksanen. On the construction of virtual interior point source travel time distances from the hyperbolic Neumann-to-Dirichlet map. *SIAM J. Appl. Math.*, 76(2):805–825, 2016.
- [14] M. V. De Hoop and K. Solna. Estimating a green’s function from field-field correlations in a random medium. *SIAM Journal on Applied Mathematics*, 69(4):909–932, 2009.
- [15] L. C. Evans. *Partial differential equations*, volume 19 of *Graduate Studies in Mathematics*. American Mathematical Society, Providence, RI, second edition, 2010.
- [16] J. Fageot, A. Fallah, and M. Unser. Multidimensional Lévy white noises in weighted Besov spaces. *arXiv preprint arXiv:1603.01502*, 2016.
- [17] J. Garnier and G. Papanicolaou. Passive sensor imaging using cross correlations of noisy signals in a scattering medium. *SIAM Journal on Imaging Sciences*, 2(2):396–437, 2009.
- [18] J. Garnier and G. Papanicolaou. *Passive Imaging with Ambient Noise*. Cambridge University Press, 2016.
- [19] I. M. Gel’fand and N. Y. Vilenkin. *Generalized functions. Vol. 4*. Academic Press [Harcourt Brace Jovanovich, Publishers], New York-London, 1964 [1977]. Applications of harmonic analysis, Translated from the Russian by Amiel Feinstein.

- [20] T. Helin, M. Lassas, and L. Oksanen. An inverse problem for the wave equation with one measurement and the pseudorandom source. *Analysis & PDE*, 5(5):887–912, 2012.
- [21] T. Helin, M. Lassas, and L. Oksanen. Inverse problem for the wave equation with a white noise source. *Communications in Mathematical Physics*, 332(3):933–953, 2014.
- [22] T. Helin, M. Lassas, and L. Päivärinta. Inverse acoustic scattering problem in half-space with anisotropic random impedance. *arXiv preprint arXiv:1407.2481*, 2014.
- [23] L. Hörmander. *The analysis of linear partial differential operators. I*. Springer Study Edition. Springer-Verlag, Berlin, second edition, 1990. Distribution theory and Fourier analysis.
- [24] L. Hörmander. *The analysis of linear partial differential operators III: Pseudo-differential operators*, volume 274. Springer Science & Business Media, 2007.
- [25] H. Isozaki, Y. Kurylev, and M. Lassas. Forward and inverse scattering on manifolds with asymptotically cylindrical ends. *J. Funct. Anal.*, 258(6):2060–2118, 2010.
- [26] L. Isserlis. On a formula for the product-moment coefficient of any order of a normal frequency distribution in any number of variables. *Biometrika*, 12(1/2):134–139, 1918.
- [27] S. I. Kabanikhin, A. D. Satybaev, and M. A. Shishlenin. *Direct methods of solving multidimensional inverse hyperbolic problems*, volume 48. Walter de Gruyter, 2004.
- [28] A. Kachalov, Y. Kurylev, and M. Lassas. *Inverse boundary spectral problems*. CRC Press, 2001.
- [29] A. Kachalov, Y. Kurylev, and M. Lassas. *Inverse boundary spectral problems*, volume 123 of *Monographs and Surveys in Pure and Applied Mathematics*. Chapman & Hall/CRC, Boca Raton, FL, 2001.
- [30] W. P. Klingenberg. *Riemannian geometry*, volume 1. Walter de Gruyter, 1995.
- [31] J. Korpela, M. Lassas, and L. Oksanen. Regularization strategy for an inverse problem for a $1+1$ dimensional wave equation. *Inverse Problems*, 32(6):065001, 2016.
- [32] Y. Kurylev, L. Oksanen, and G. P. Paternain. Inverse problems for the connection Laplacian. *Submitted. Preprint arXiv:1509.02645*, 2015.
- [33] S. Kusuoka. The support property of a Gaussian white noise and its applications. *J. Fac. Sci. Univ. Tokyo Sect. IA Math.*, 29(2):387–400, 1982.
- [34] M. Lassas, L. Päivärinta, and E. Saksman. Inverse scattering problem for a two dimensional random potential. *Communications in Mathematical Physics*, 279(3):669–703, 2008.
- [35] M. Lassas and T. Saksala. Determination of a riemannian manifold from the distance difference functions. *arXiv preprint arXiv:1510.06157*, 2015.
- [36] S. Liu and L. Oksanen. A Lipschitz stable reconstruction formula for the inverse problem for the wave equation. *Trans. Amer. Math. Soc.*, 368(1):319–335, 2016.
- [37] A. Millet and M. Sanz-Solé. A stochastic wave equation in two space dimension: smoothness of the law. *Annals of Probability*, pages 803–844, 1999.
- [38] A. I. Nachman. Reconstructions from boundary measurements. *Ann. of Math. (2)*, 128(3):531–576, 1988.
- [39] L. Oksanen. Solving an inverse problem for the wave equation by using a minimization algorithm and time-reversed measurements. *Inverse Probl. Imaging*, 5(3):731–744, 2011.
- [40] M. Ondreját. Existence of global mild and strong solutions to stochastic hyperbolic evolution equations driven by a spatially homogeneous Wiener process. *J. Evol. Equ.*, 4(2):169–191, 2004.
- [41] L. Pestov, V. Bolgova, and O. Kazarina. Numerical recovering of a density by the bc-method. *Inverse Probl. Imaging*, 4(4):701–712, 2010.
- [42] N. M. Shapiro, M. Campillo, L. Stehly, and M. H. Ritzwoller. High-resolution surface-wave tomography from ambient seismic noise. *Science*, 307(5715):1615–1618, 2005.

- [43] L. Stehly, M. Campillo, and N. Shapiro. A study of the seismic noise from its long-range correlation properties. *Journal of Geophysical Research: Solid Earth*, 111(B10), 2006.
- [44] J. Sylvester and G. Uhlmann. A global uniqueness theorem for an inverse boundary value problem. *Ann. of Math. (2)*, 125(1):153–169, 1987.
- [45] D. Tataru. Unique continuation for solutions to pde’s; between hormander’s theorem and holmgren’s theorem. *Communications in partial differential equations*, 20(5-6):855–884, 1995.
- [46] M. E. Taylor. *Partial differential equations. I*, volume 115 of *Applied Mathematical Sciences*. Springer-Verlag, New York, 1996. Basic theory.
- [47] F. Trèves. *Topological vector spaces, distributions and kernels*. Academic Press, New York-London, 1967.
- [48] B. Vainberg. *Asymptotic methods in equations of mathematical physics*. CRC Press, 1989.
- [49] B. R. Vainberg. On the short wave asymptotic behaviour of solutions of stationary problems and the asymptotic behaviour as $t \rightarrow \infty$ of solutions of non-stationary problems. *Russian Mathematical Surveys*, 30(2):1, 1975.
- [50] G. Vodev. Local energy decay of solutions to the wave equation for nontrapping metrics. *Arkiv för Matematik*, 42(2):379–397, 2004.
- [51] S. Willard. *General topology*. Dover Publications, Inc., Mineola, NY, 2004.
- [52] H. Yao, R. D. van Der Hilst, and V. Maarten. Surface-wave array tomography in se tibet from ambient seismic noise and two-station analysisi. phase velocity maps. *Geophysical Journal International*, 166(2):732–744, 2006.

E-mail address: teemu.saksala@helsinki.fi