

PROBABILISTIC LIMIT SHAPES AND HARMONIC FUNCTIONS

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EASTERN FINLAND

joint work with

Rick Kenyon (Yale)

©Avid seminar - October 2021

TWO SIDES

Gradient variational problems in \mathbb{R}^2
[arXiv:2006.01219](#), 2020

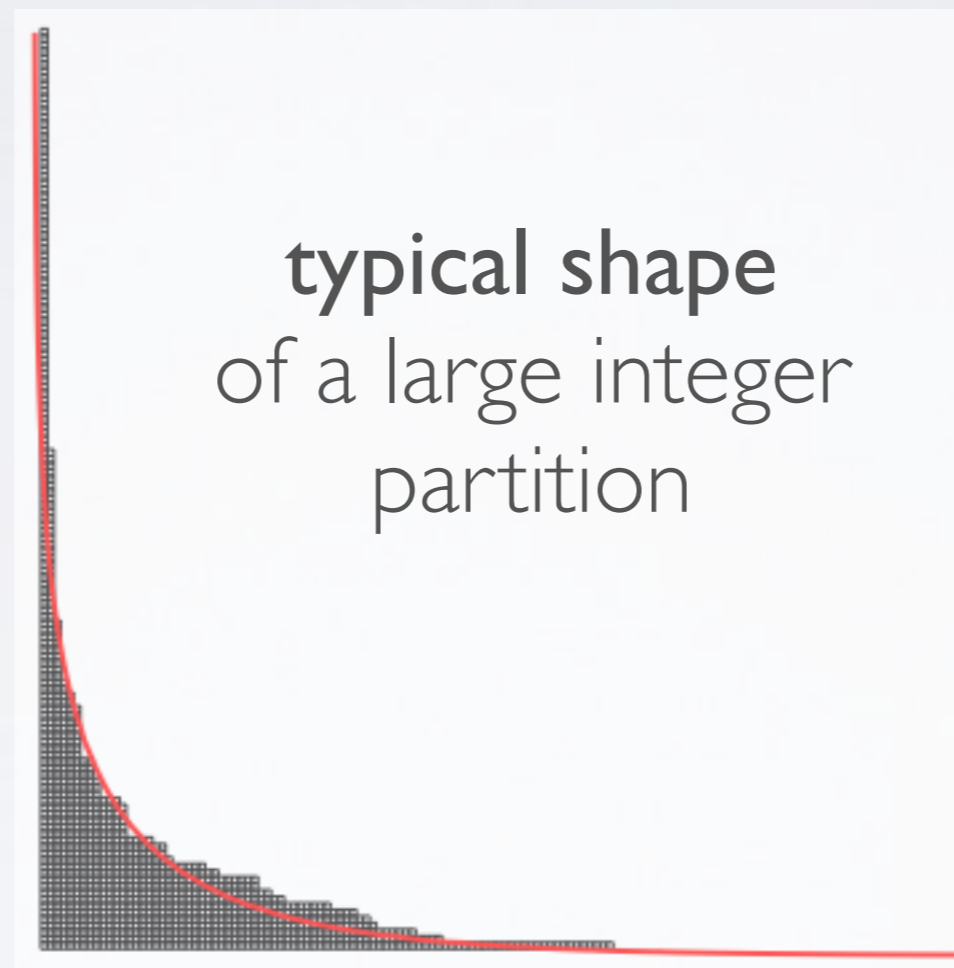
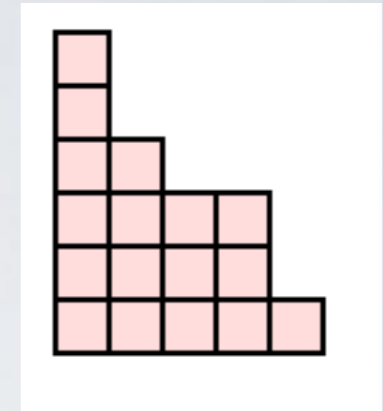
The genus-zero five-vertex model
[arXiv:2101.04195](#), 2021

LIMIT SHAPE OF 2D YOUNG DIAGRAMS

Vershik

integer partitions

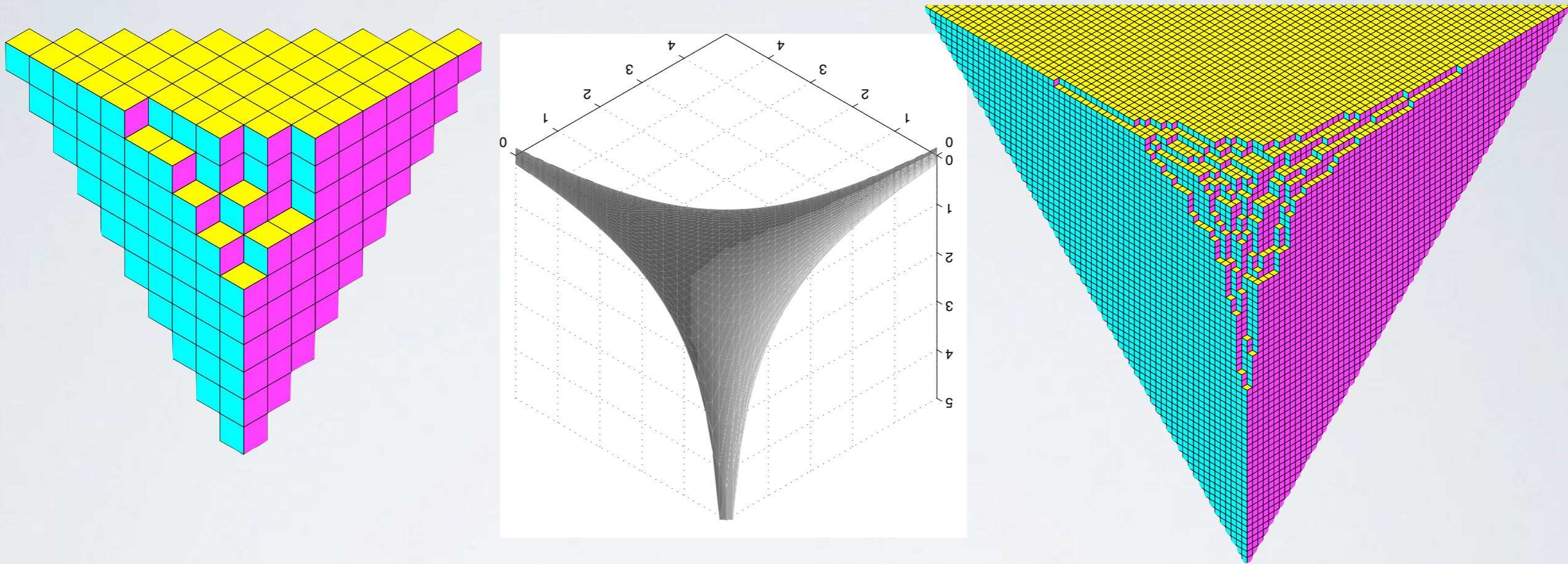
$$17 = 5 + 4 + 4 + 2 + 1 + 1$$



$$e^{-cx} + e^{-cy} = 1$$

3D YOUNG DIAGRAM LIMIT SHAPE

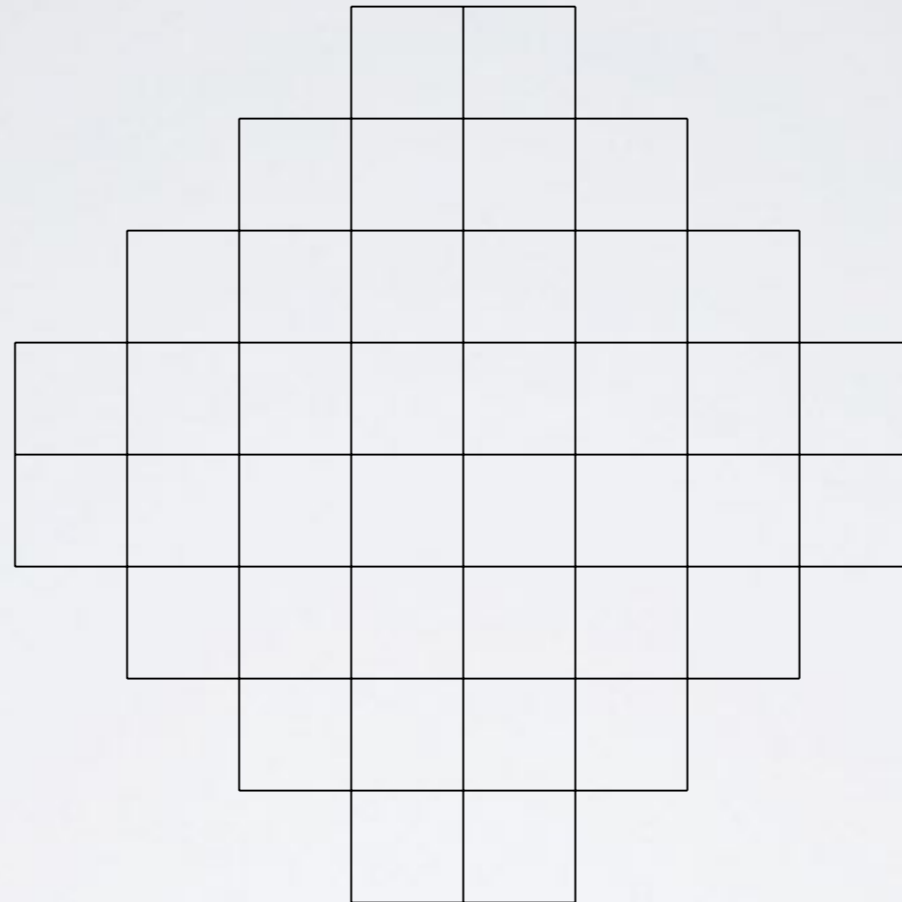
Cerf-Kenyon, Okounkov-Reshetikhin



“melting” equilibrium **crystal**

planar projection: lozenge tilings, algebraic geometry

THE AZTEC DIAMOND



consider a **uniformly random** tiling by
dominos (2×1)

How does it look like?

THE ARCTIC CIRCLE

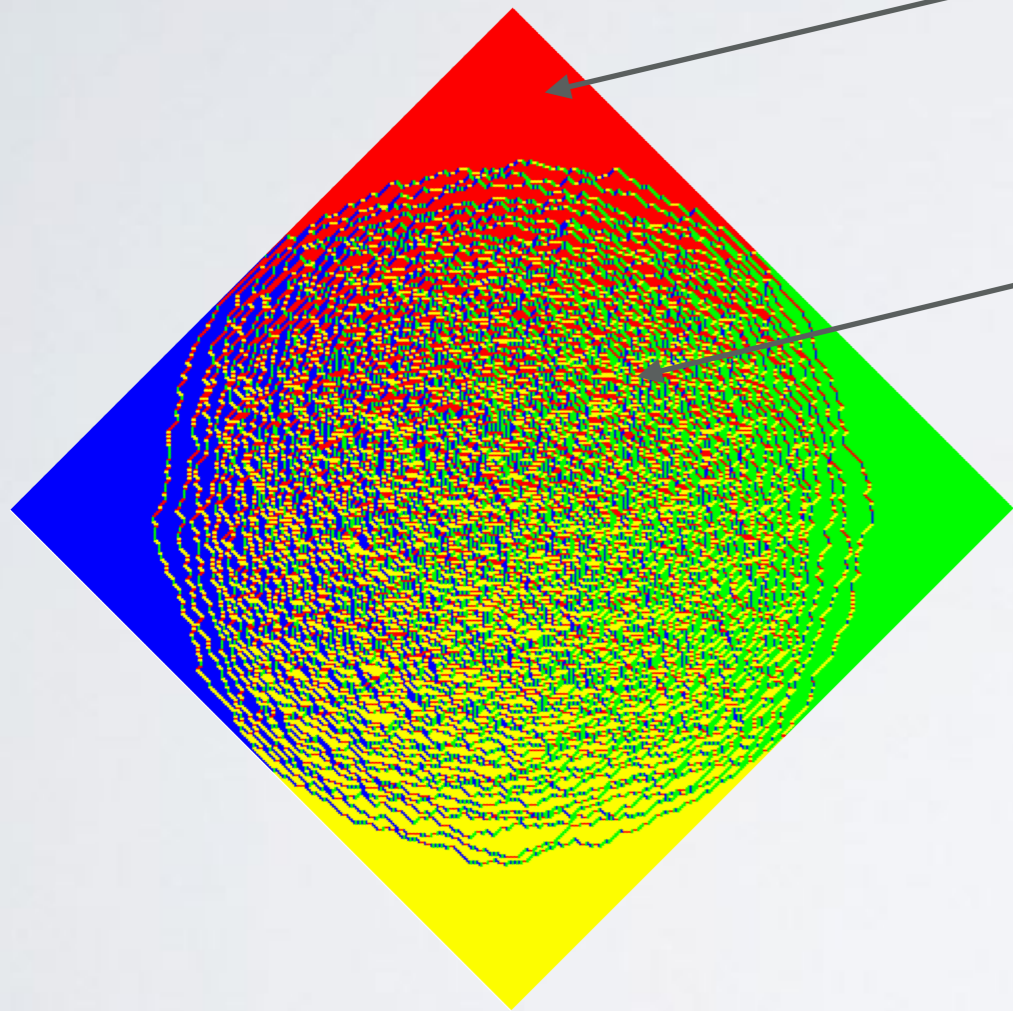
Jockusch-Propp-Shor

Cohn-Larsen-Propp

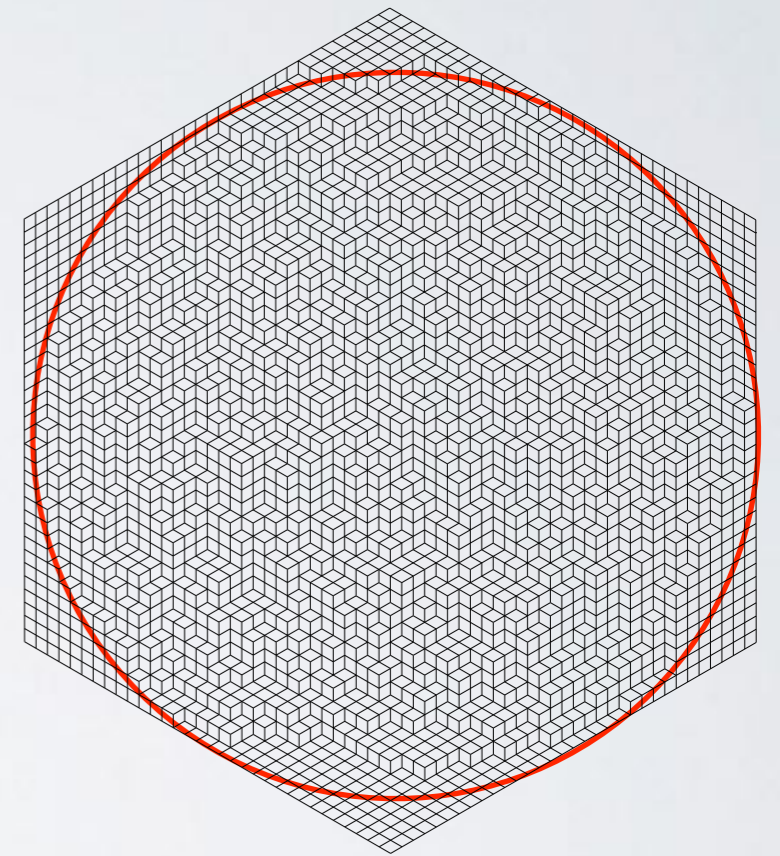
dominos

frozen

liquid

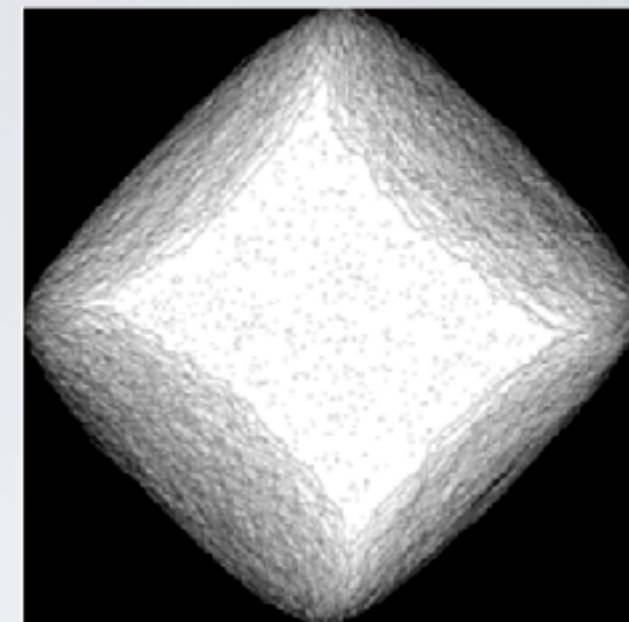
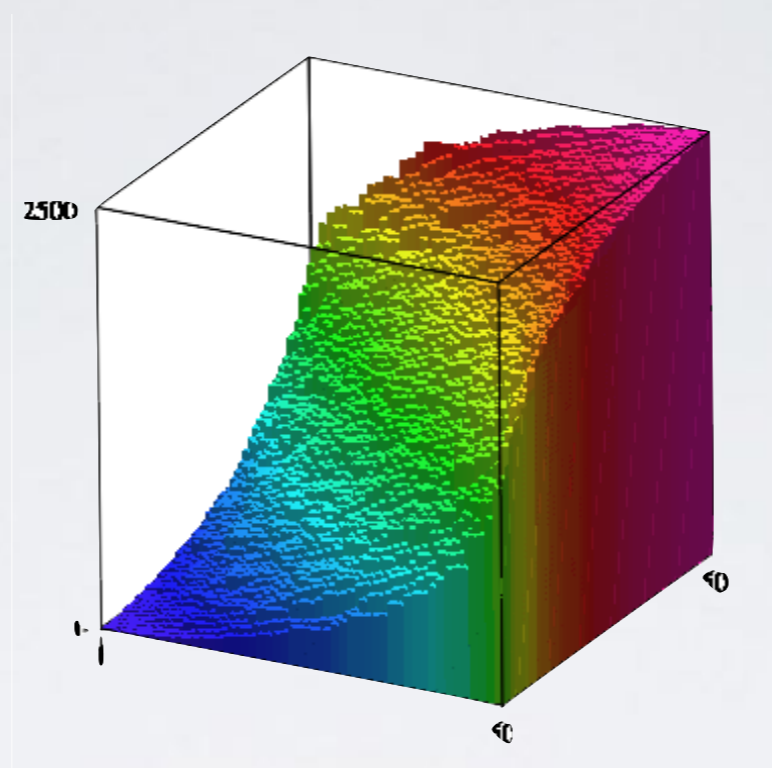
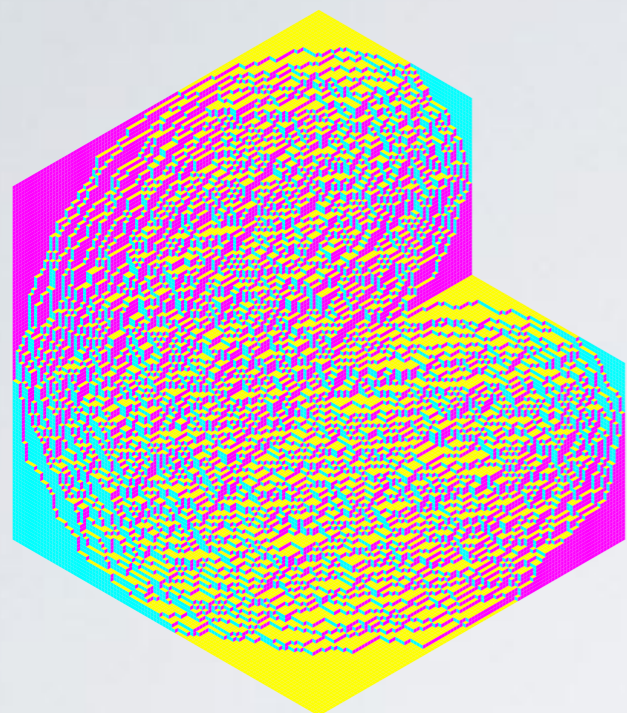


lozenges



harmonic functions ?

ZOO OF LIMIT SHAPES



general boundary conditions & a variety of models

variational approach

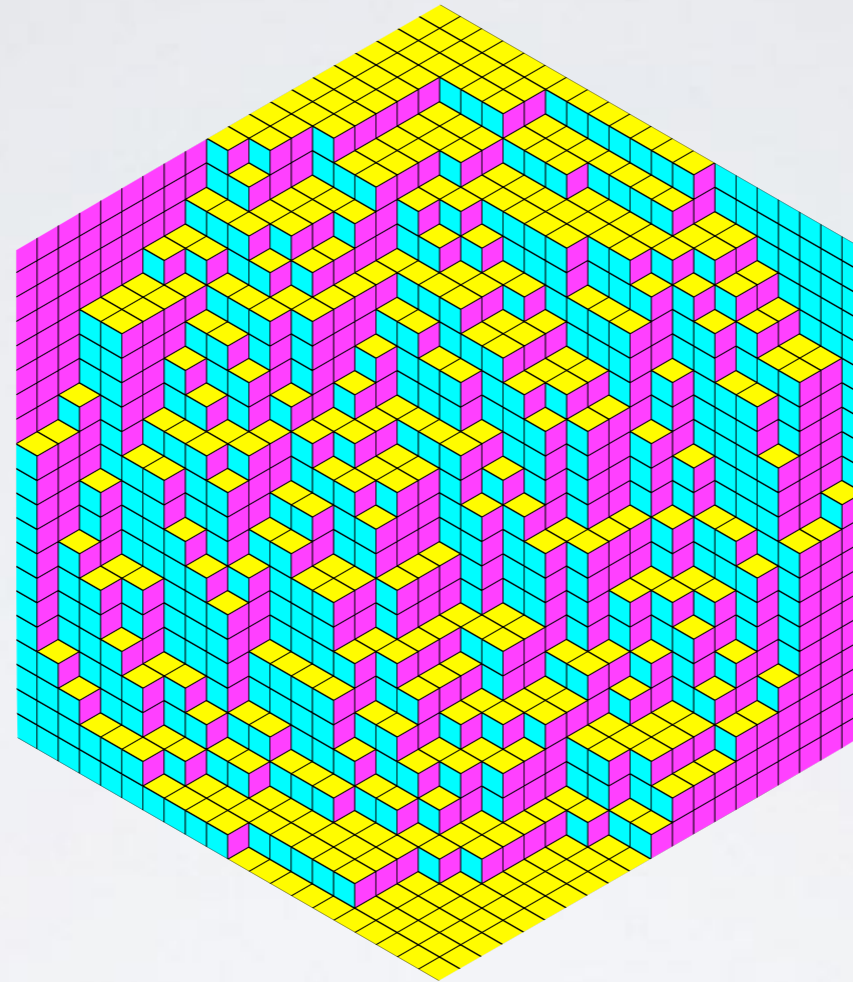
dimer model
(domino/lozenge tilings
etc)

|
random
Young tableaux

five-vertex model

HEIGHT FUNCTION

Thurston



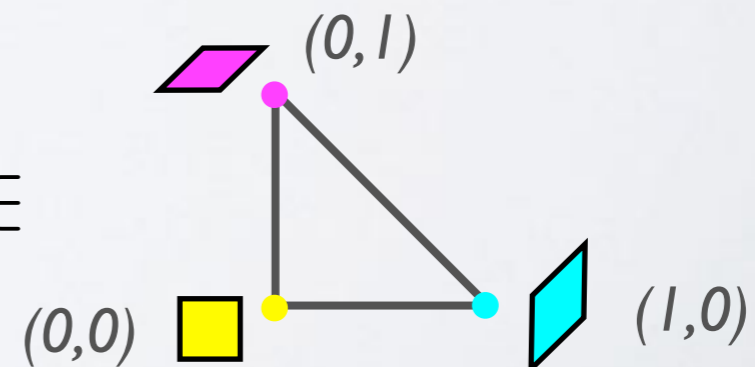
3d surface
(simulations
by A. & M. Borodin)

random stepped surface

$\mathcal{N} = \text{unit triangle}$

height function

$\nabla h \in$



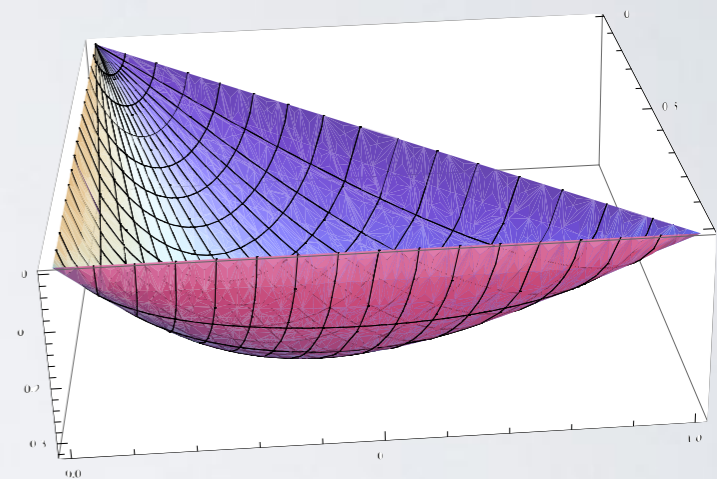
LIMIT SHAPE THEOREM

Cohn-Kenyon-Propp

The random surfaces as mesh size $\rightarrow 0$
concentrate around a deterministic surface,
called **limit shape**

The limit shape is a minimizer of a
variational problem

‘minimal surface’ spanning a wire-frame

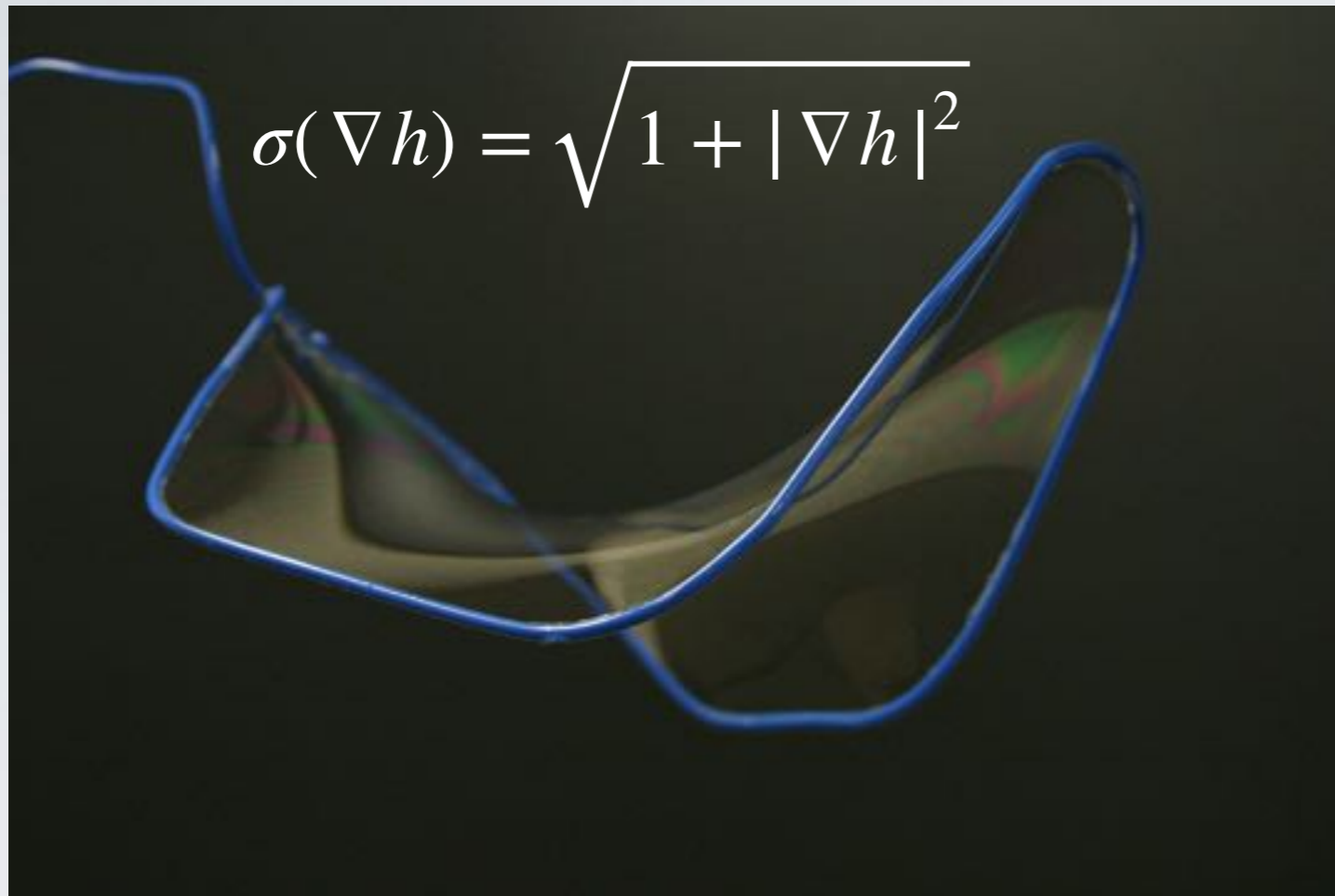


analytic, strictly
convex surface tension
in the interior
singular and **degenerates**
on the boundary

$h: \Omega \rightarrow \mathbb{R}$ Lipschitz

$$\min_h \int_{\Omega} \sigma(\nabla h), \quad \nabla h \in \mathcal{N}$$
$$h|_{\partial\Omega} = h_0$$

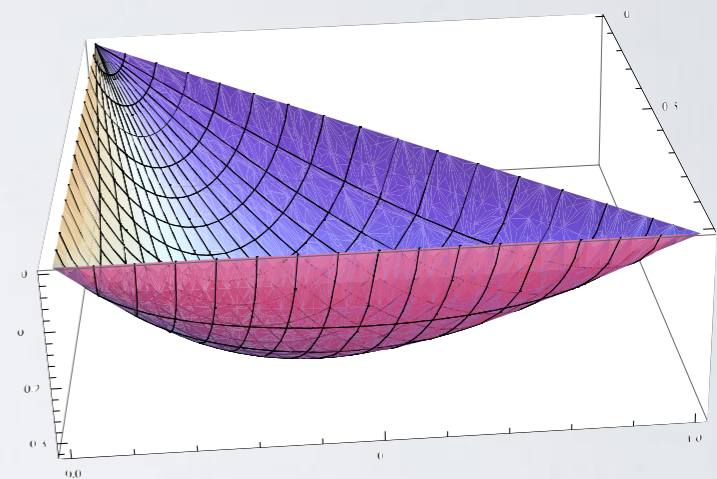
LIMIT SHAPE THEOREM



pp

h size $\rightarrow 0$
 stochastic surface,

er of a



‘minimal surface’ spanning a wire-frame

analytic, strictly
 convex surface tension
 in the interior
singular and *degenerates*
 on the boundary

$$h: \Omega \rightarrow \mathbb{R} \quad \text{Lipschitz}$$

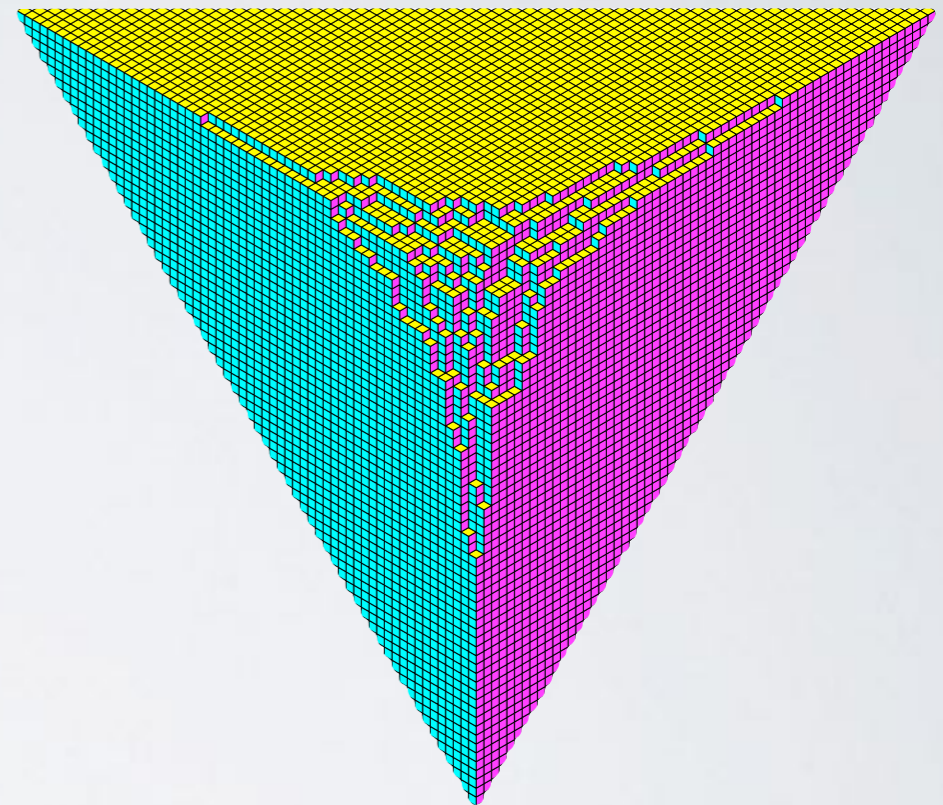
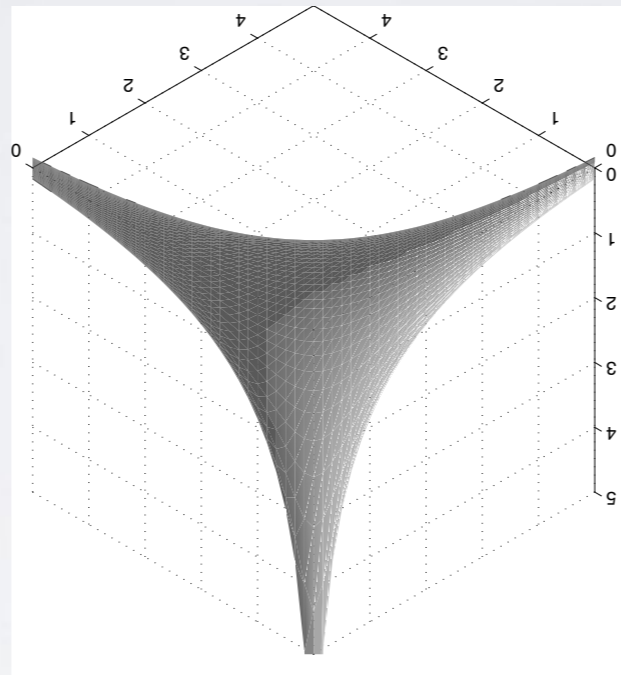
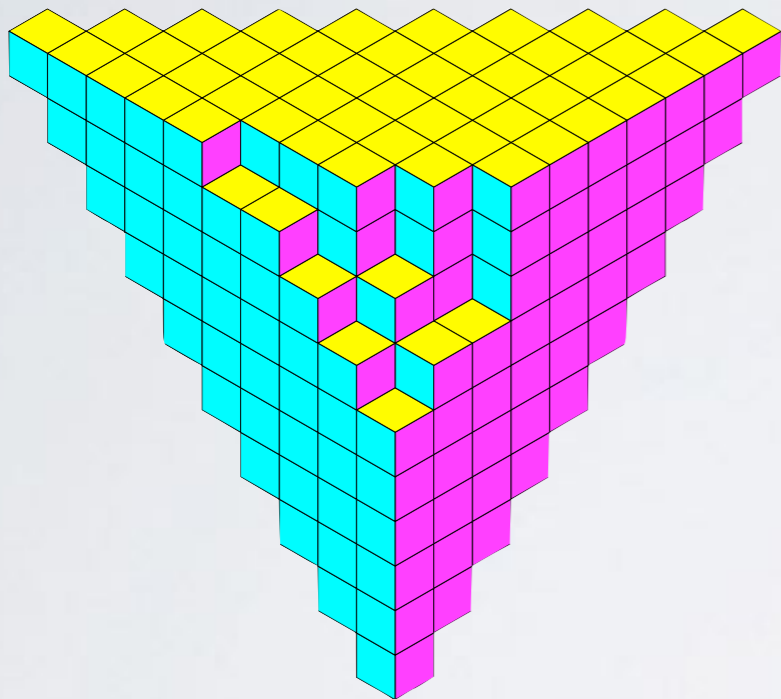
$$\min_h \int_{\Omega} \sigma(\nabla h), \quad \nabla h \in \mathcal{N}$$

$$h|_{\partial\Omega} = h_0$$

WULFF SHAPE

Wulff shape - Legendre dual of surface tension
itself a **limit shape**

(lozenge tilings \leftrightarrow 3D Young diagram)



“fundamental solution”

(facets, facet-rough transition, phases, algebraic boundary etc)

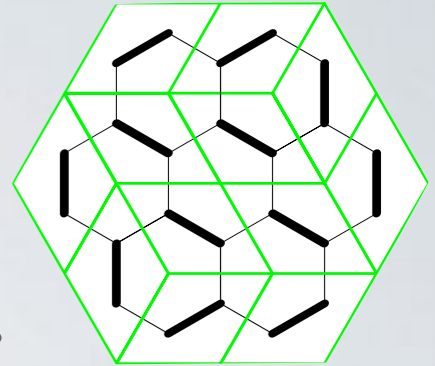
Kenyon-Okounkov-Sheffield

$\det D^2 \sigma \equiv \pi^2$ for the dimer model (determinantal)

DETERMINANTAL MODEL

Kasteleyn

The number of dimer covers for G (a subgraph of hexagonal lattice) is $|\det(K)|$, where K is the (bi)adjacency matrix of the bipartite graph G .



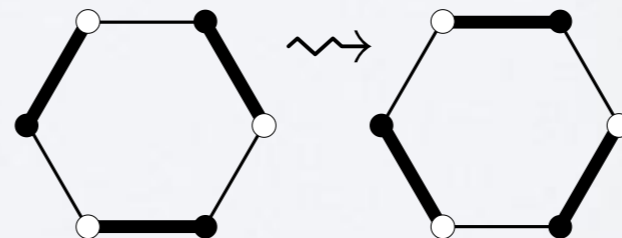
$$\det(K) = \sum_{\sigma \in S_n} \left(\text{sgn}(\sigma) \prod_{i=1}^n K(b_i, w_{\sigma(i)}) \right)$$

determinant counts *perfect matchings* with **signs**

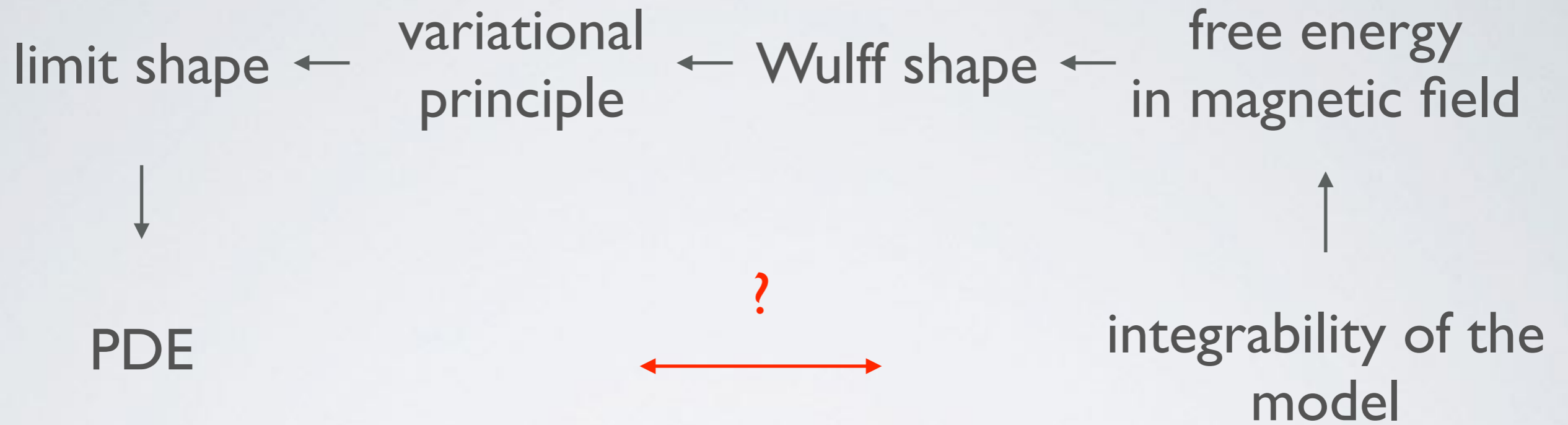
key point: signs are **consistent**

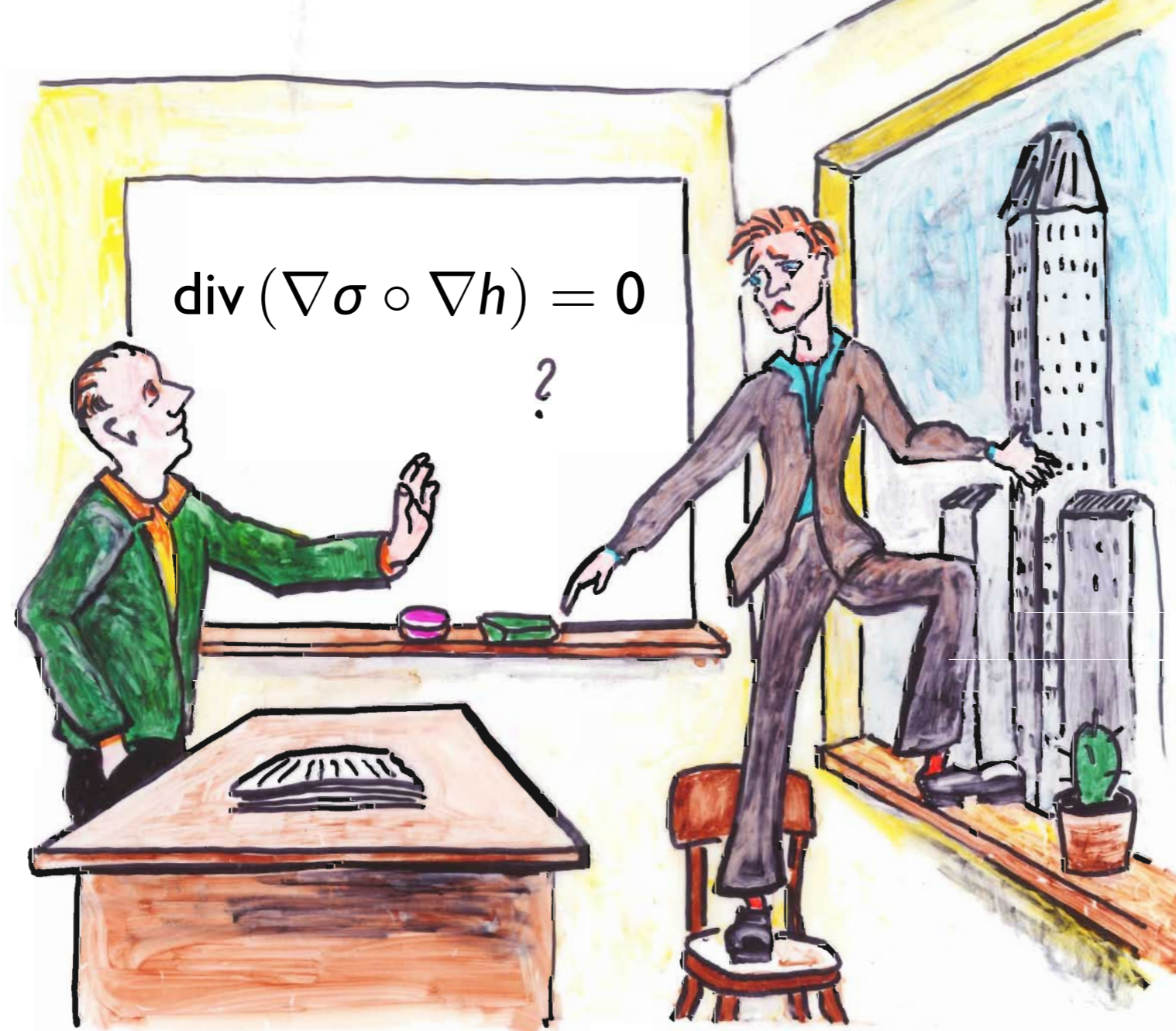
a hexagon flip move changes σ by a 3-cycle (even permutation)

adds/removes cubes



INTEGRABLE PDE ?





Don't jump, “complexify to simplify” !

a not-so-“hidden” *complex variable*

(fluctuations, integrability, isothermal)

ISOTHERMAL COORDINATE

Gauss

Riemannian metric associated to the Wulff shape

let z be the **isothermal** coordinate

$$\sigma_{ss} ds^2 + 2\sigma_{st} ds dt + \sigma_{tt} dt^2 = \rho |dz|^2$$

$$(s, t) \in \mathcal{N} \leftrightarrow z \in \mathbb{C}$$

$$(x, y) \in \mathcal{L} \mapsto \nabla h(x, y) \mapsto z(x, y)$$

liquid

this is the **conformal** coordinate of the model

“Write everything in terms of z ”

κ -HARMONIC ENVELOPE

Kenyon-Prause

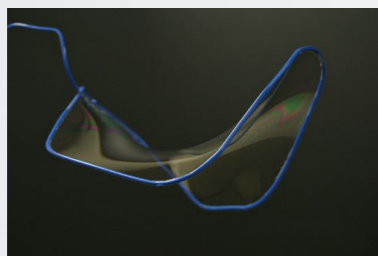
$$\kappa(\mathbf{z}) = \sqrt{\det D^2 \sigma} \text{ as a function of } z$$

$$\nabla \cdot \kappa \nabla u = 0$$

Thm: s, t and $h-(sx+ty)$ are all κ -harmonic(z) in the liquid region
(multi-valued in z)

Corollary:

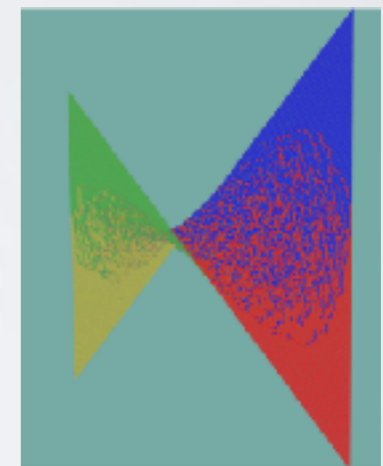
Limit shapes in the dimer model are *envelopes* of harmonically moving planes in \mathbb{R}^3



minimal surfaces

$$(x, y, h(x, y))$$

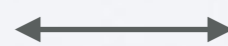
$$\mu_z = 0$$



dimer limit shapes

$$(h_x, h_y, h - \nabla h \cdot (x, y))$$

$$\mu_{\bar{z}} = 0$$



κ -HARMONIC ENVELOPE

Kenyon-Prause

$\kappa(\mathbf{z}) = \sqrt{\det D^2\sigma}$ as a function of \mathbf{z}

$$\nabla \cdot \kappa \nabla u = 0$$

Thm: s, t and $h-(sx+ty)$ are all κ -harmonic(\mathbf{z}) in the liquid region
(multi-valued in \mathbf{z})

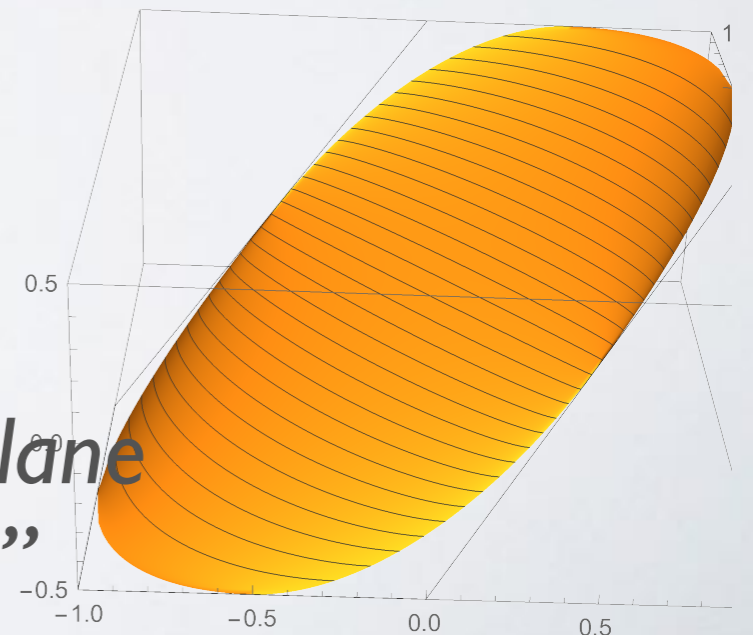
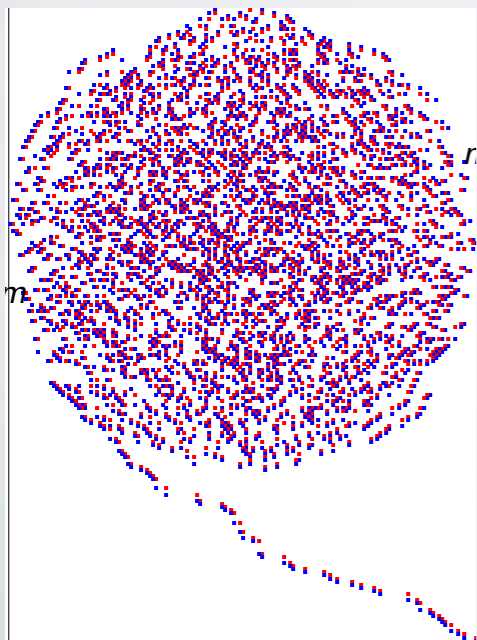
Corollary:

Limit shapes in the dimer model
are *envelopes* of harmonically
moving planes in \mathbb{R}^3

limit shape
algorithm

tangent method
Colomo-Sportiello

“tangent plane
method”



BACK TO THE ARCTIC CIRCLE

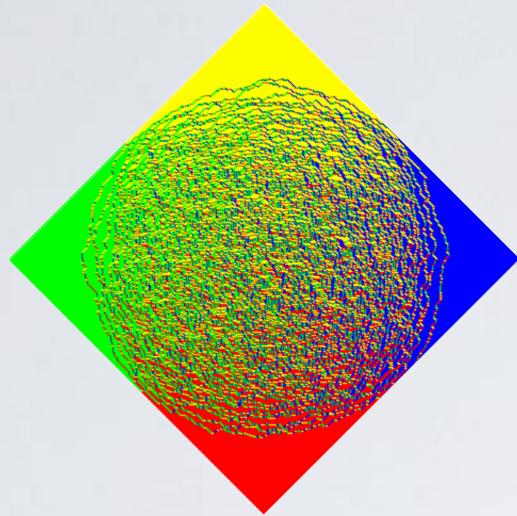
envelope of planes $x_3 = sx + ty + c$

$$s_z x + t_z y + c_z = 0$$

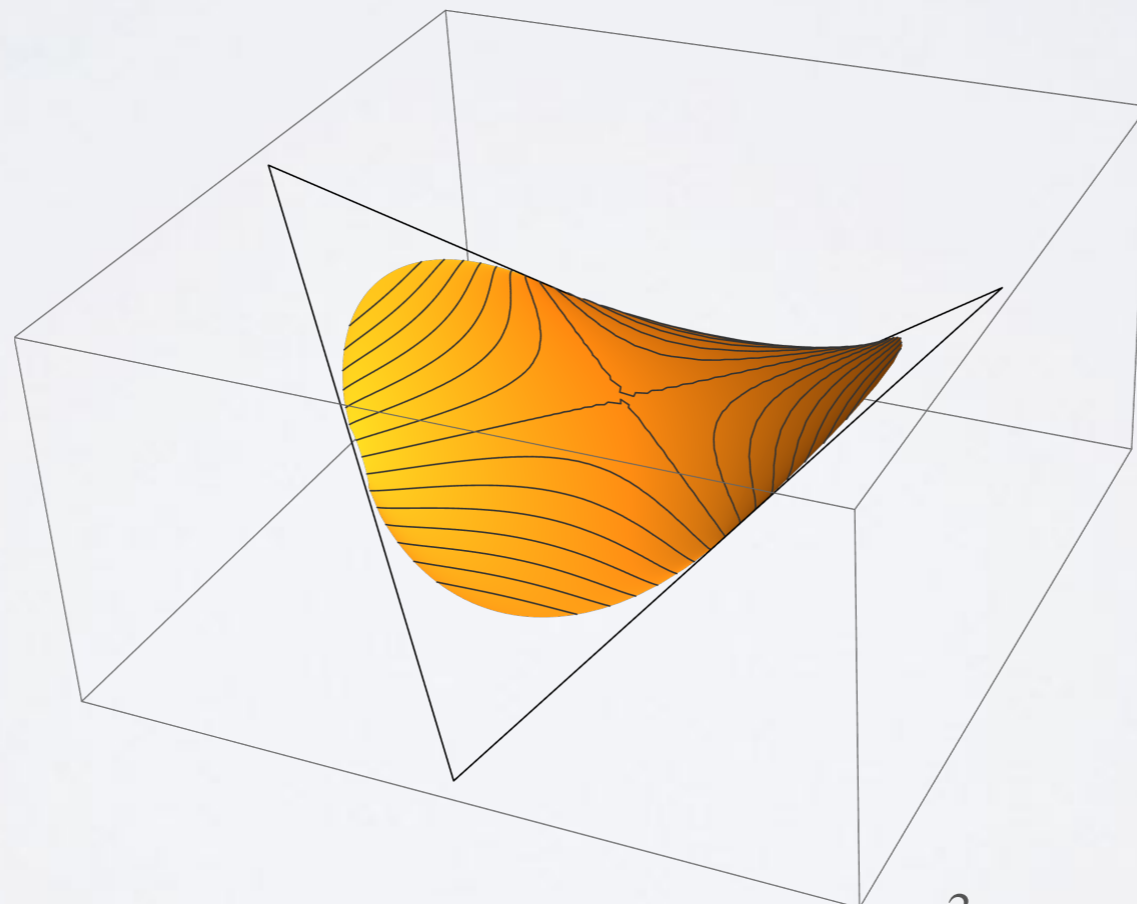
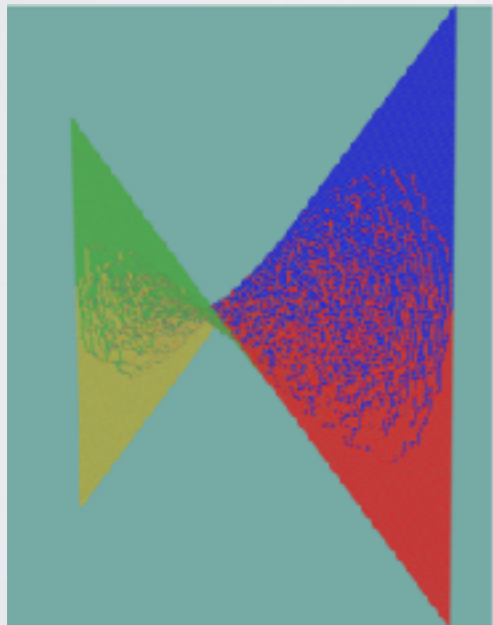
$$s = \omega - \omega$$

$$t = \omega - \omega$$

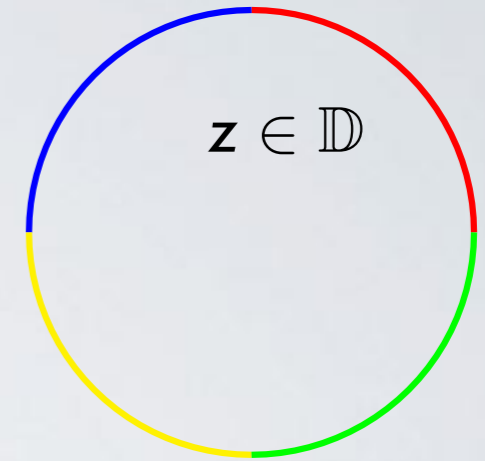
$$c = \omega - \omega + \omega - \omega$$



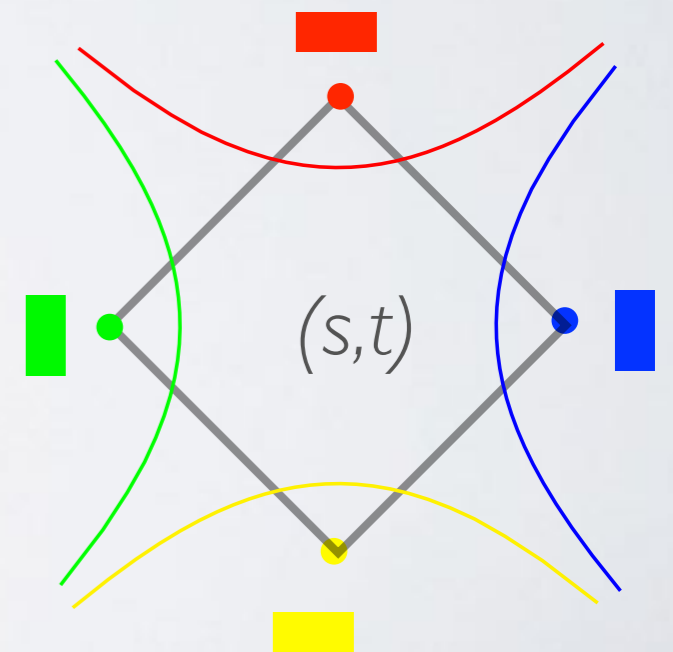
$(x, y) \in \mathcal{L}$



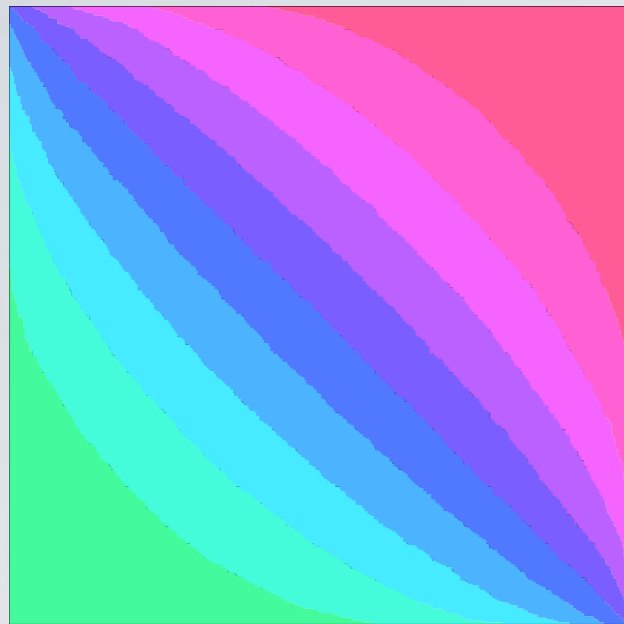
$$z \mapsto (x(z), y(z), x_3(z)) \subset \mathbb{R}^3$$



↓ harmonic

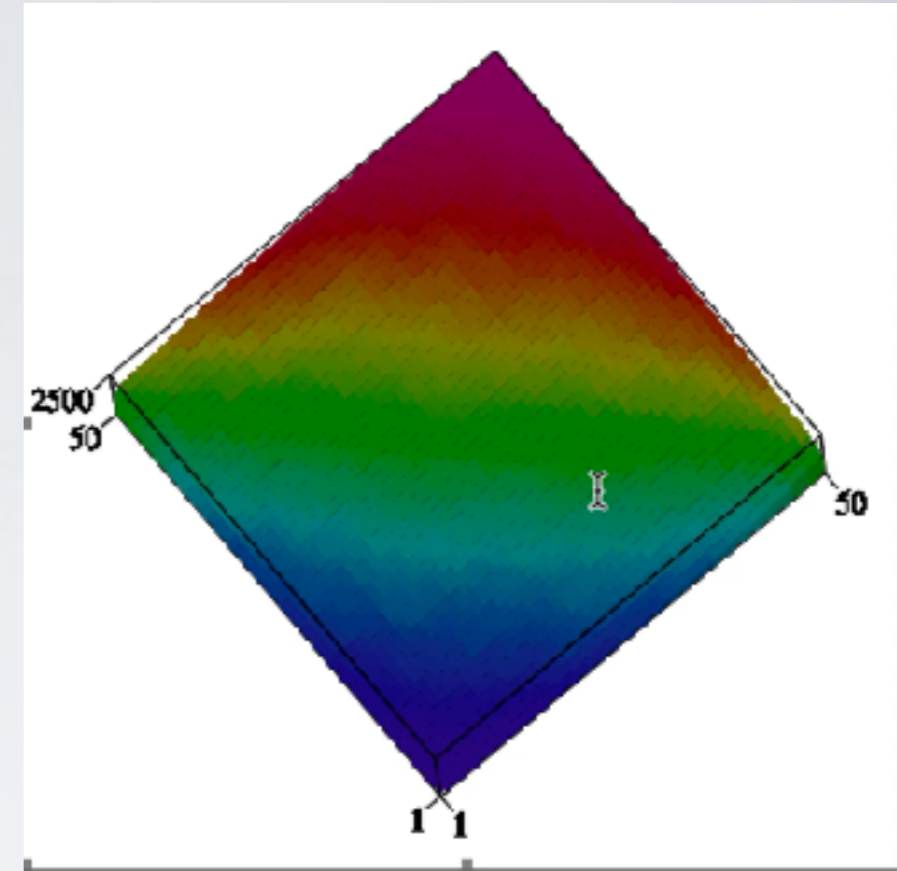


YOUNG TABLEAUX EXAMPLE



Dan Romik's
MacTableaux

square shape
Pittel-Romik
Biane

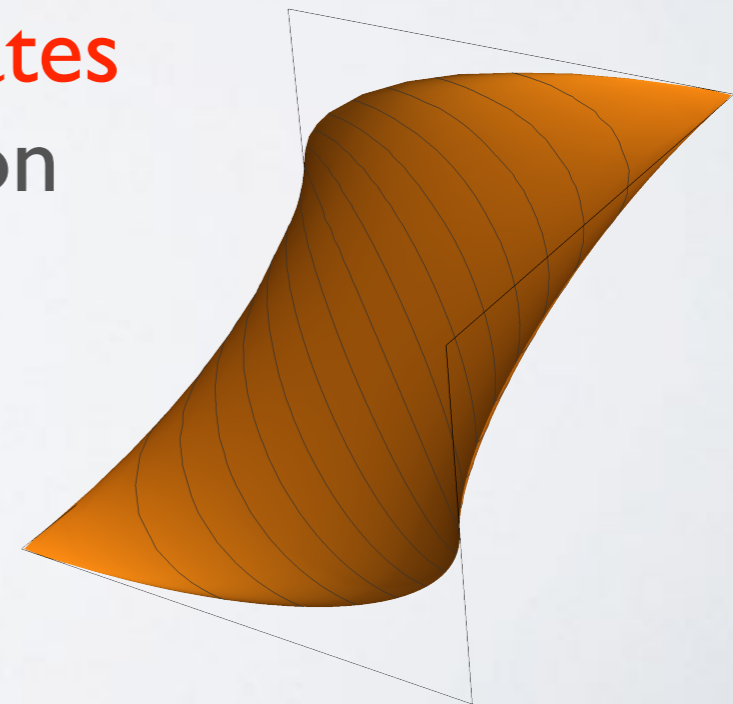


Cyril Banderier et al.'s
YoungPackage

W. Sun variational principle

$\kappa \equiv 1 \rightarrow$ harmonic coordinates
for surface tension

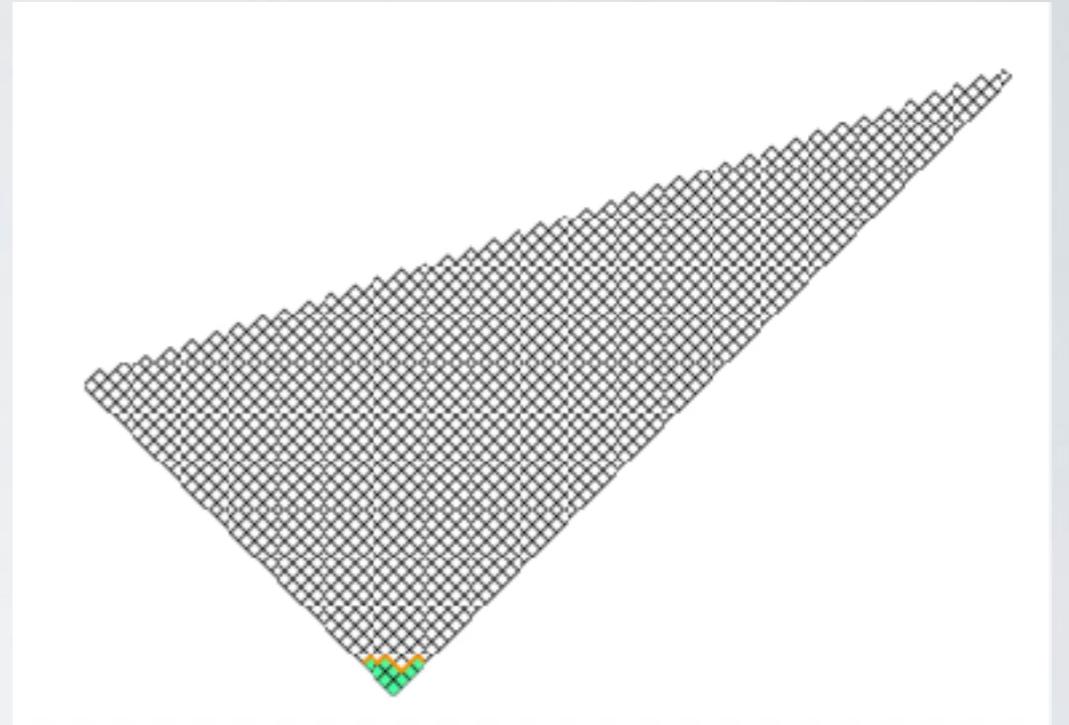
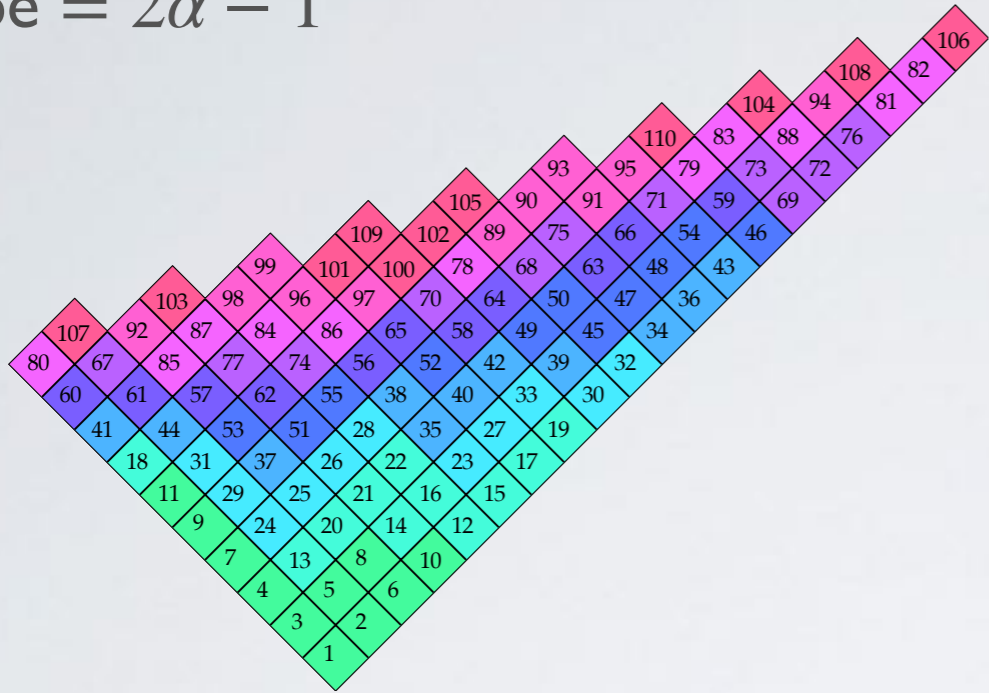
limit surface as harmonic envelope
(use harmonic extension of boundary facets)



SLANTED STAIRCASE

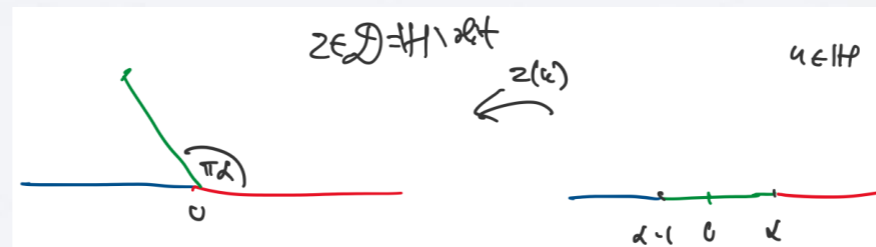
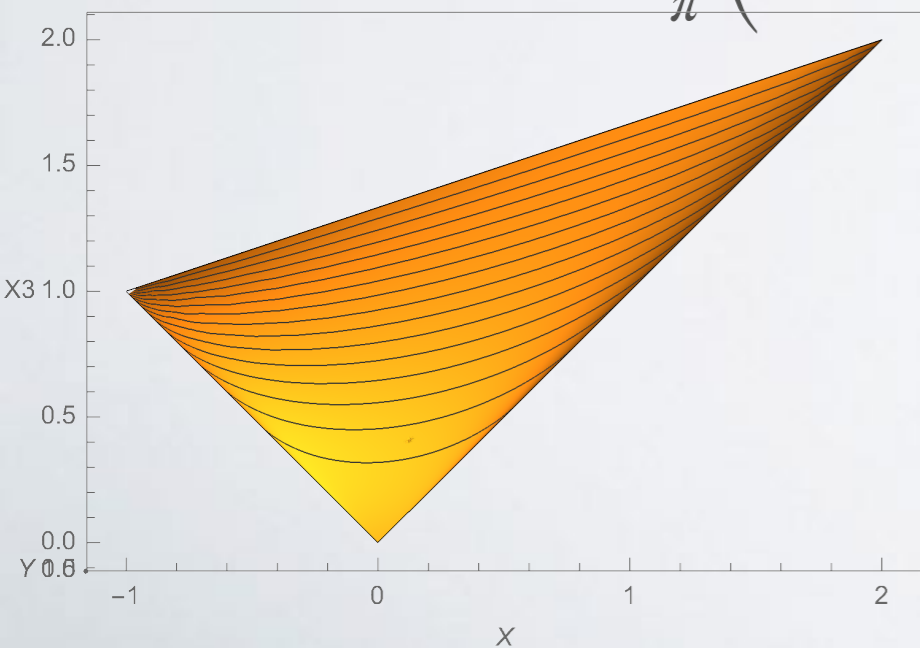
slope = $2\alpha - 1$

P



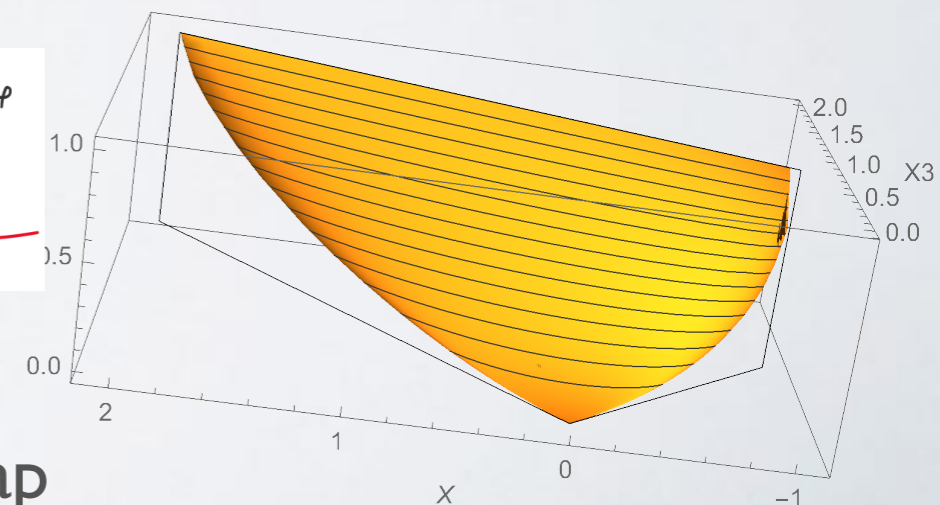
$$x(u) = -\frac{1}{1-\alpha} \frac{\operatorname{Im} \frac{u-2\alpha+1}{z(u)}}{\operatorname{Im} \frac{1}{z(u)}}, \quad y(u) = 1 - \frac{1}{1-\alpha} \frac{\operatorname{Im}(u-2\alpha+1)}{\operatorname{Im} z(u)} \quad u \in \mathbb{H}$$

$$x_3(u) = \frac{2}{\pi} \left((\arg(z(u)) - 1)x(u) + \operatorname{Im}(z(u))(y(u) - 1) + \frac{1}{1-\alpha} \operatorname{Im} u + \alpha \arg \frac{u-\alpha}{u-\alpha+1} \right)$$

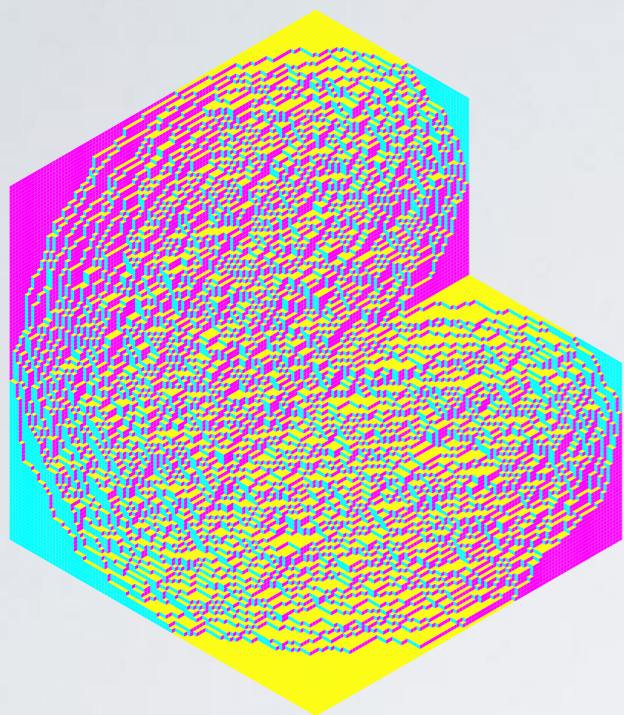


$u \mapsto z(u)$

explicit (slit) conformal map



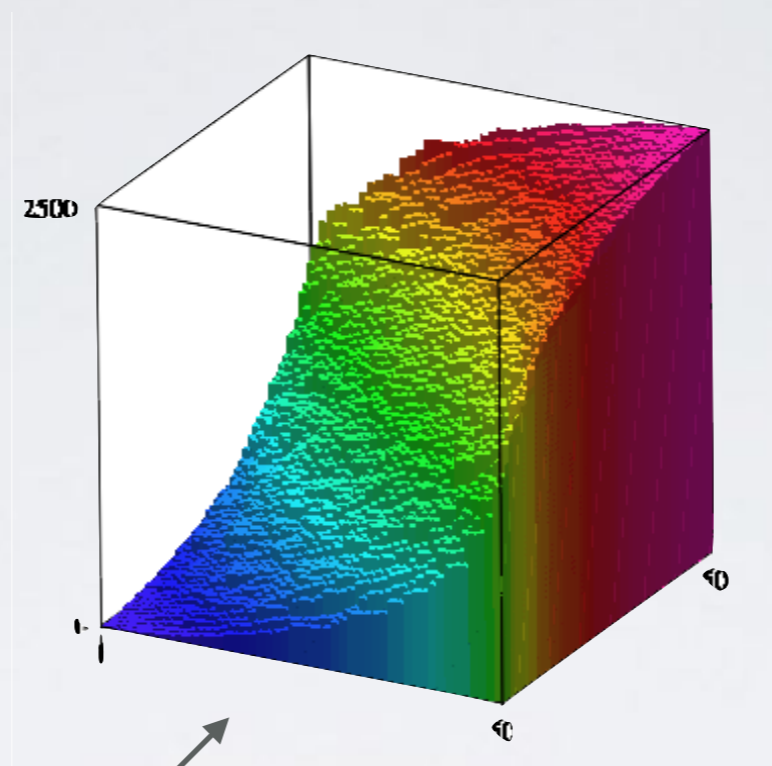
ZOO OF LIMIT SHAPES



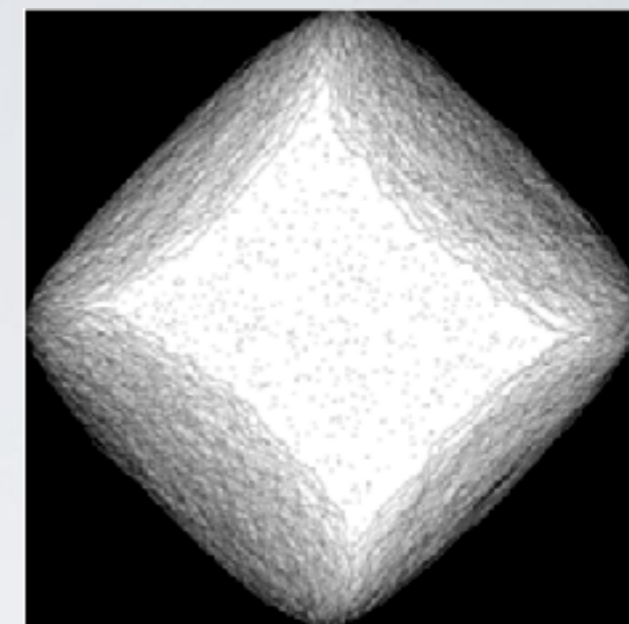
determinantal

$$\kappa \equiv 1$$

dimer model
(domino/lozenge tilings
etc)



random
Young tableaux



non-determinantal

$$\kappa = ?$$

five-vertex model

TRIVIAL POTENTIAL

Kenyon-Prause

$$\nabla \cdot \kappa \nabla u = 0$$

reduction to Schrödinger equation

$$(-\Delta + q)(\kappa^{1/2}u) = 0 \quad q = \frac{\Delta \kappa^{1/2}}{\kappa^{1/2}} \quad \text{potential}$$

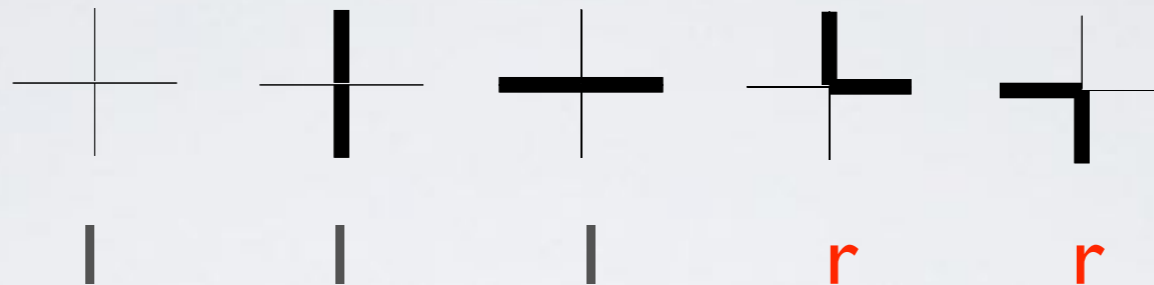
Def: a surface tension has *trivial potential* if $\sqrt[4]{\det D^2\sigma}$ is a **harmonic** function of the intrinsic coordinate z

Then κ -harmonic: $\frac{\text{harmonic}(z)}{\sqrt[4]{\det D^2\sigma}} \quad (q=0)$

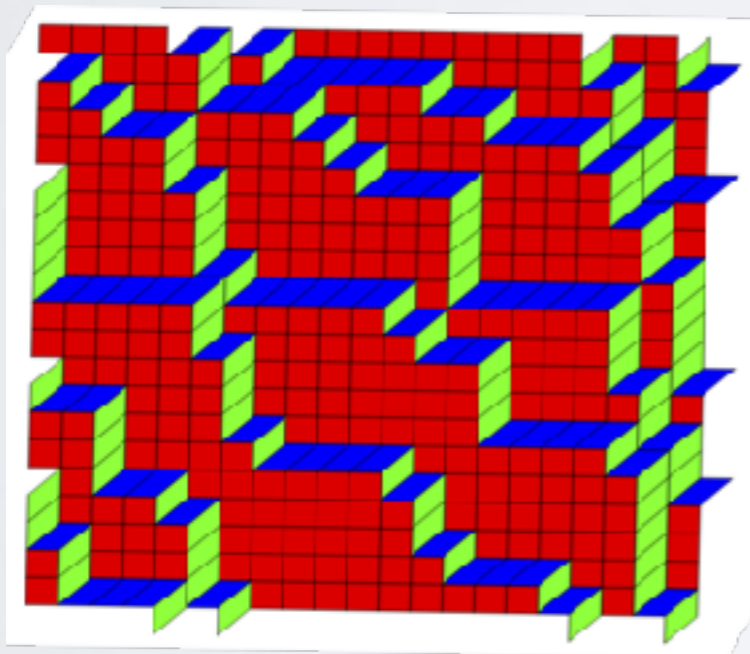
5-VERTEX MODEL

$r \neq 1$ (non-determinantal) “interacting fermions”

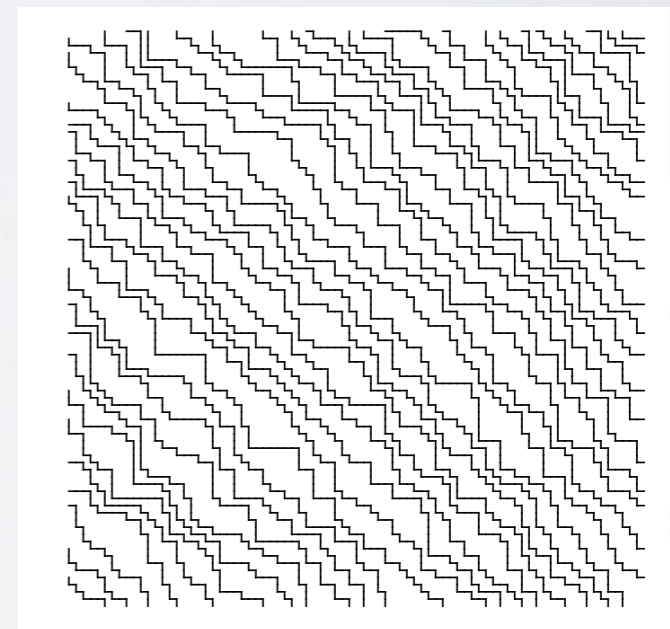
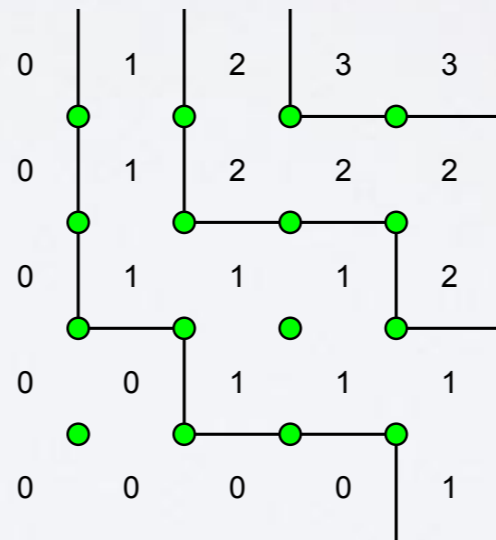
Bethe ansatz



$$P(\text{configuration}) \propto r^{\text{\#corners}}$$

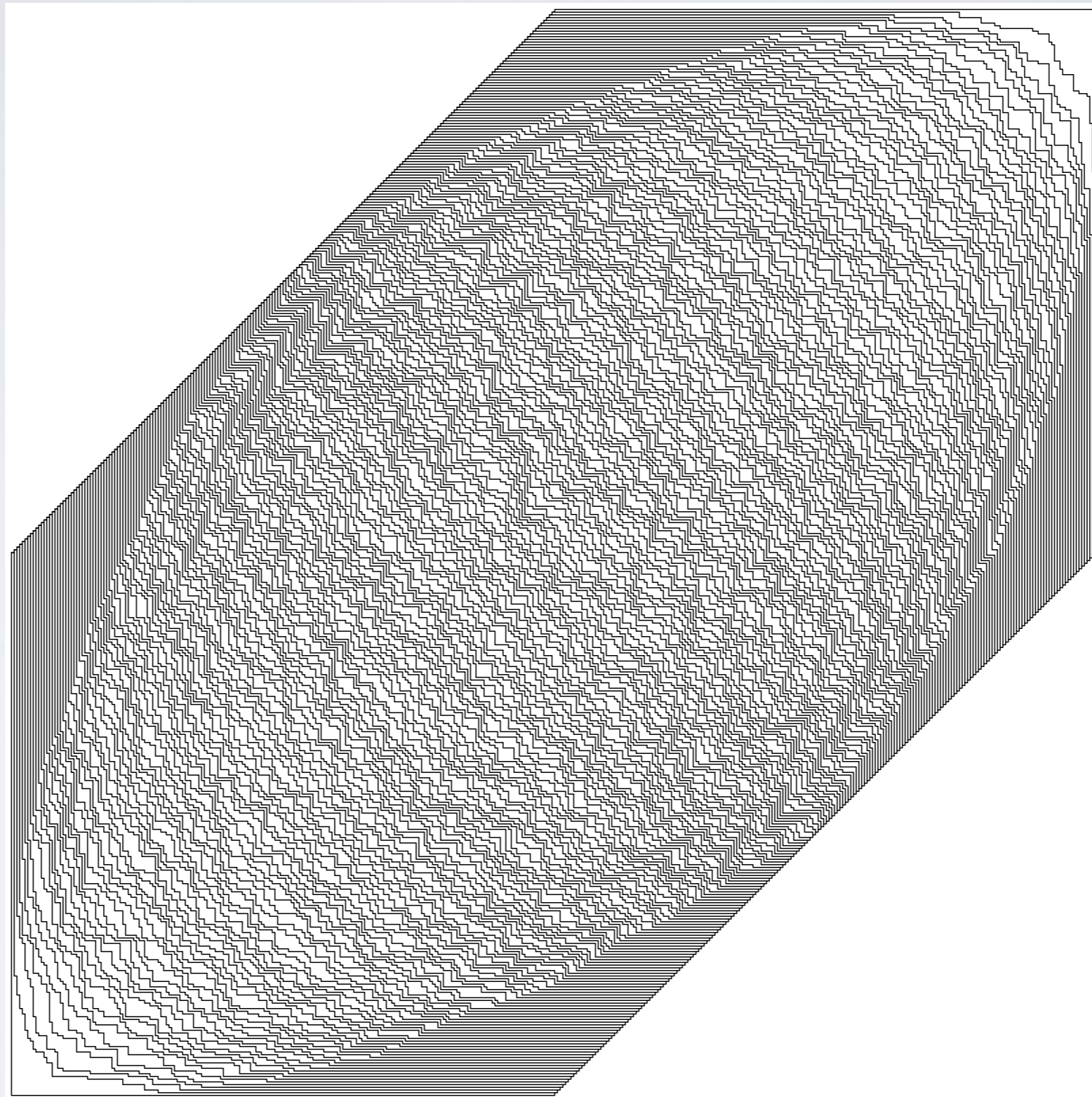


lozenge tilings with
(blue-green) *interaction*



monotone non-intersecting
lattice paths
with corners *penalized*

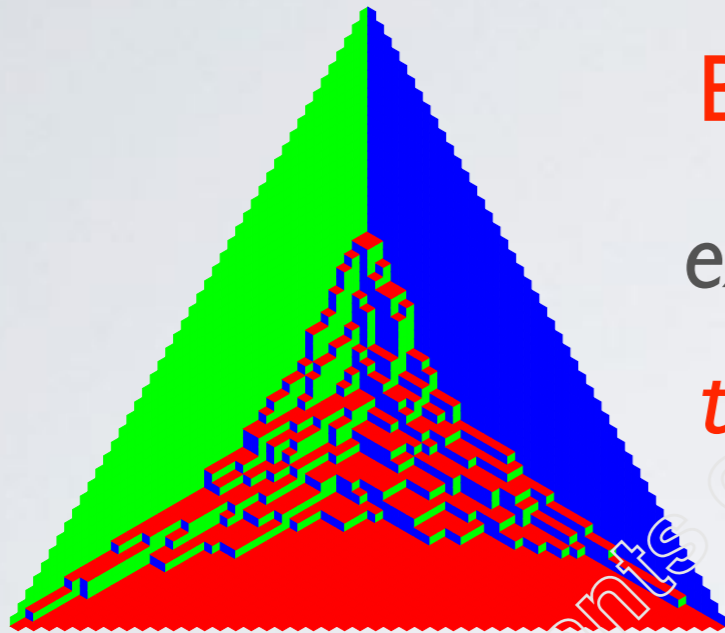
5-VERTEX BOXED PLANE PARTITION



$r=0.6$

5-VERTEX SURFACE TENSION

de Gier-Kenyon-Watson

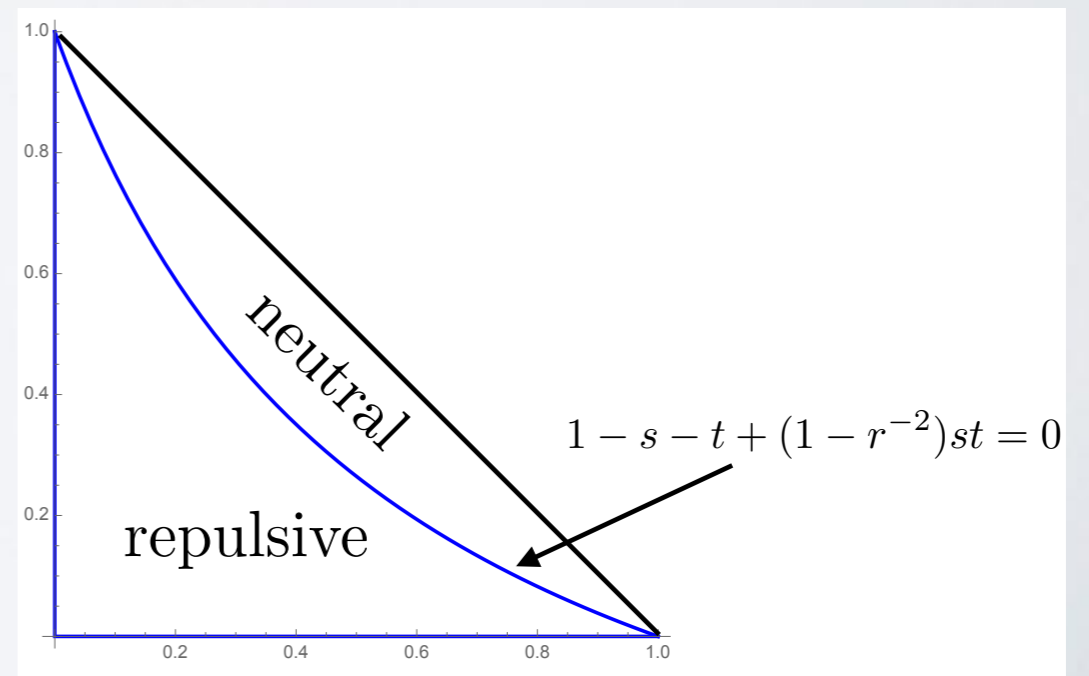
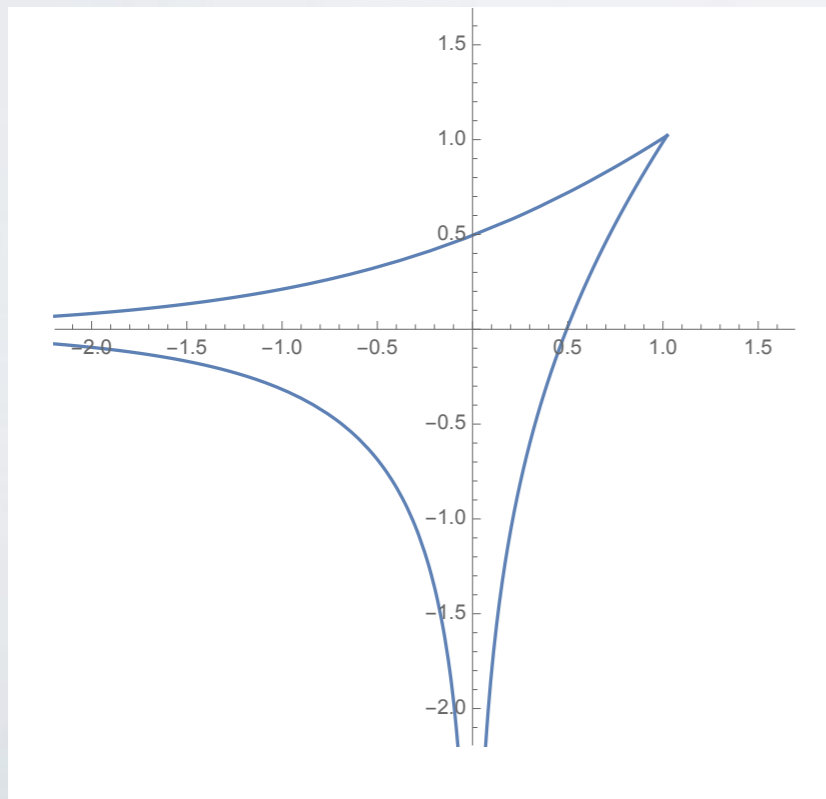
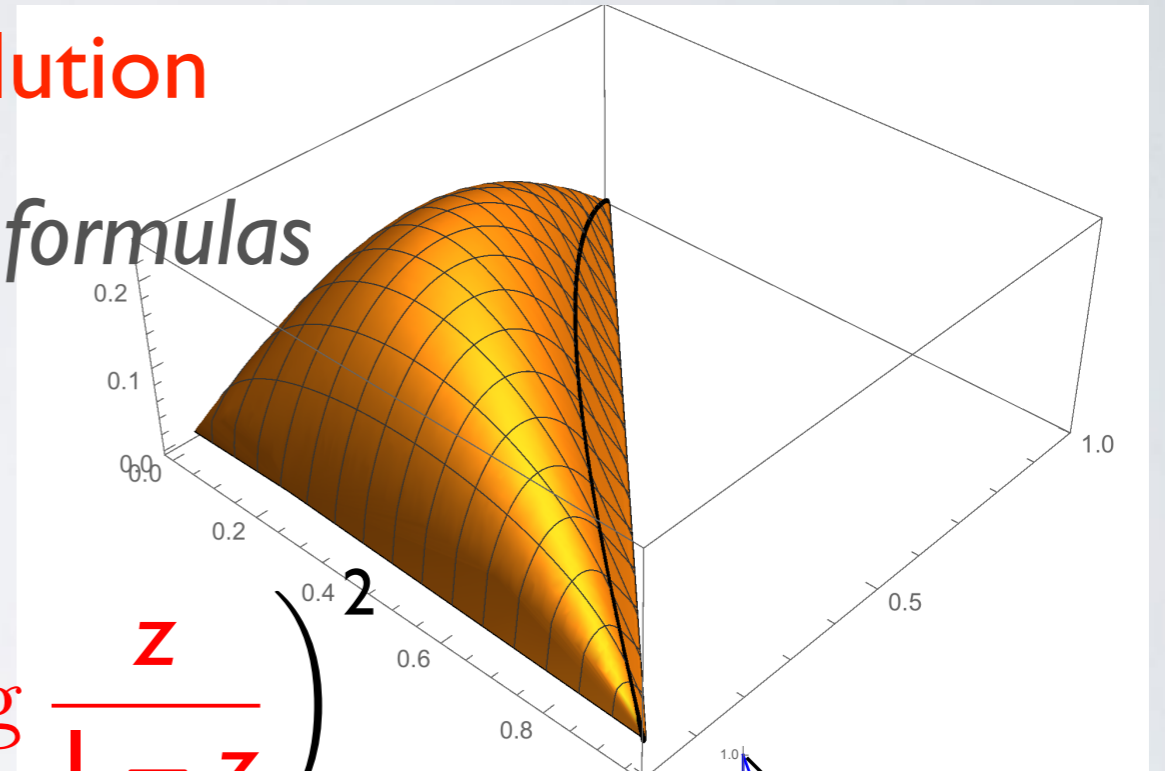


Bethe ansatz solution

explicit (involved) formulas

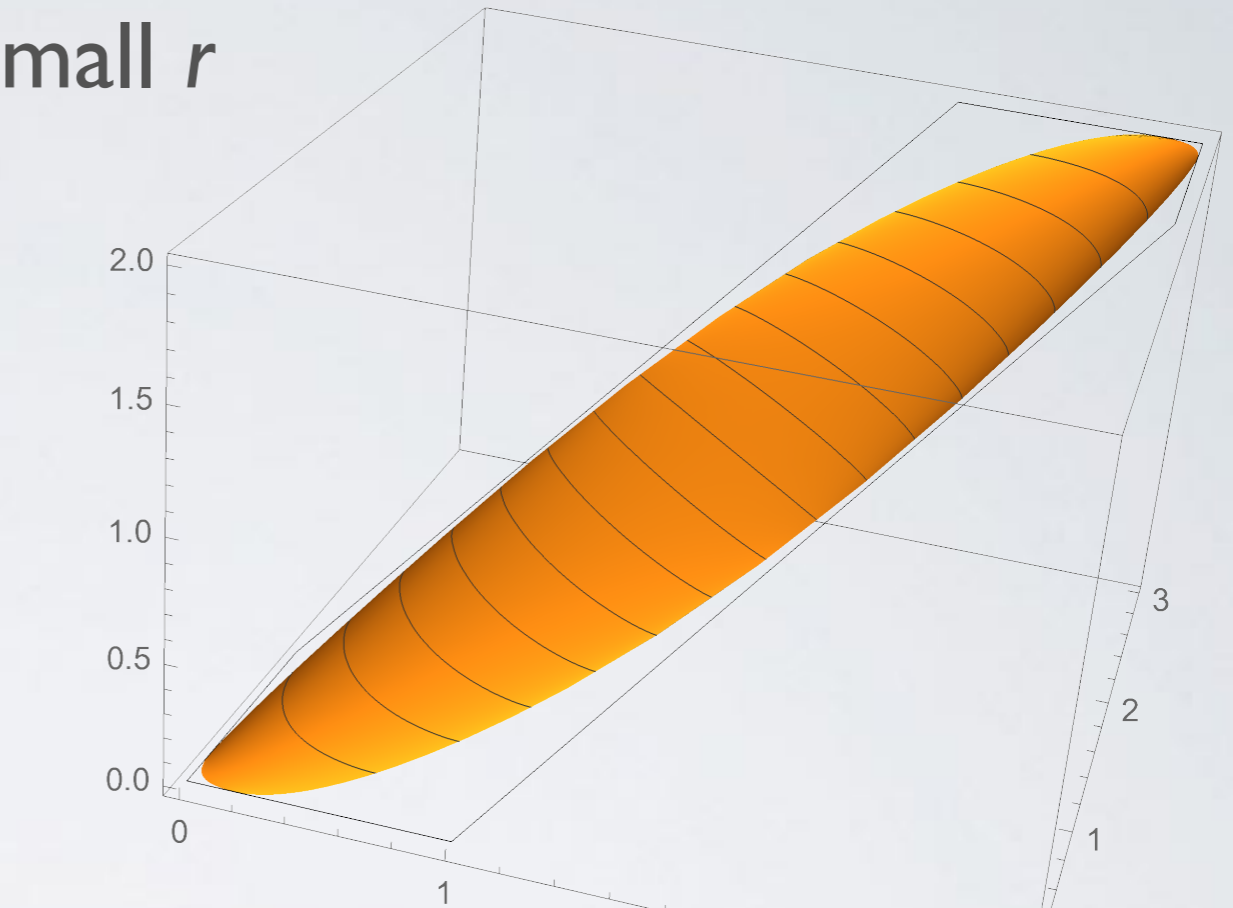
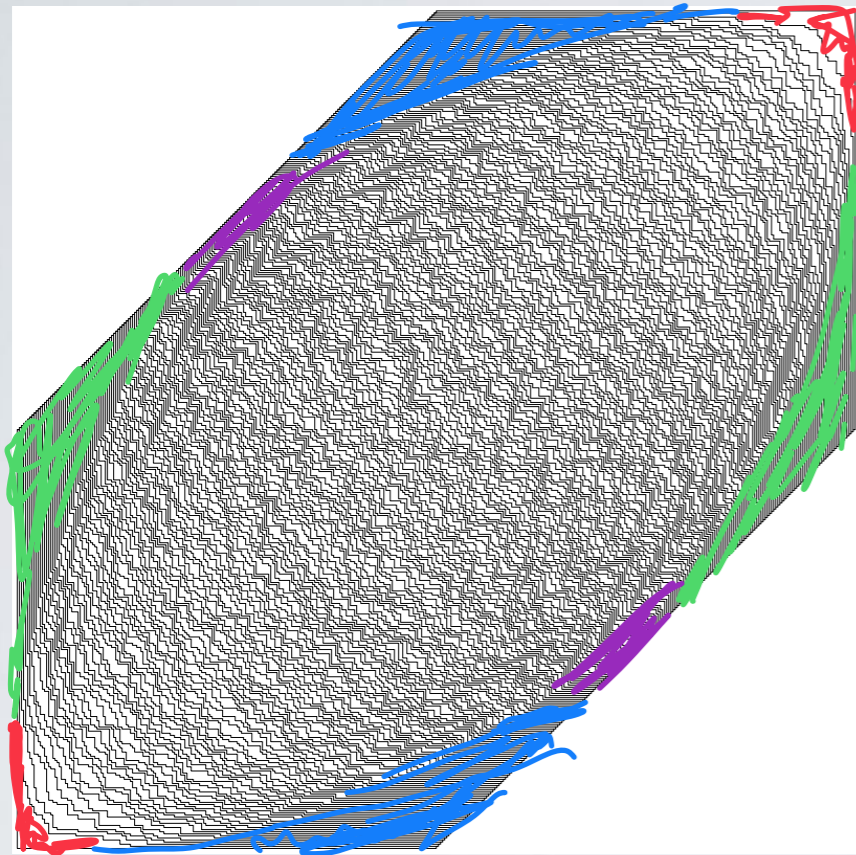
trivial potential !

$$\pi \sqrt{\det D^2 \sigma} = \left(\arg \frac{z}{1-z} \right)$$



BPP EXAMPLE

small r



envelope of planes² $\zeta \in \mathbb{H}_3$
 ($u(\zeta)$ degree 2 cover of \mathbb{H})

$$x_3 = s(\zeta)x + t(\zeta)y + c(\zeta)$$

are all *ratios* of linear combinations of *harmonic measures*

6 facets+ 2 neutral regions

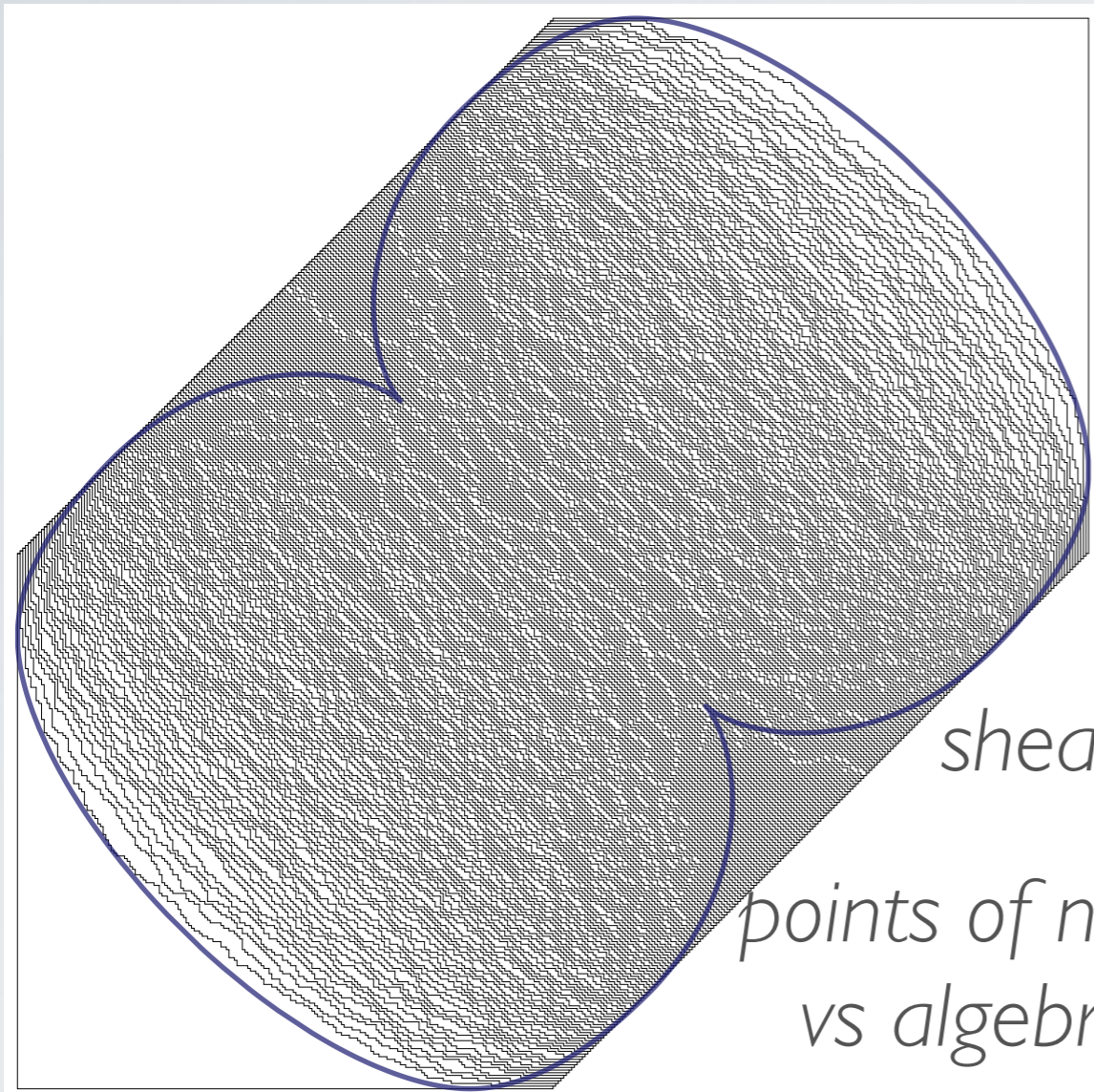


8 intervals on $\partial\mathbb{H}$

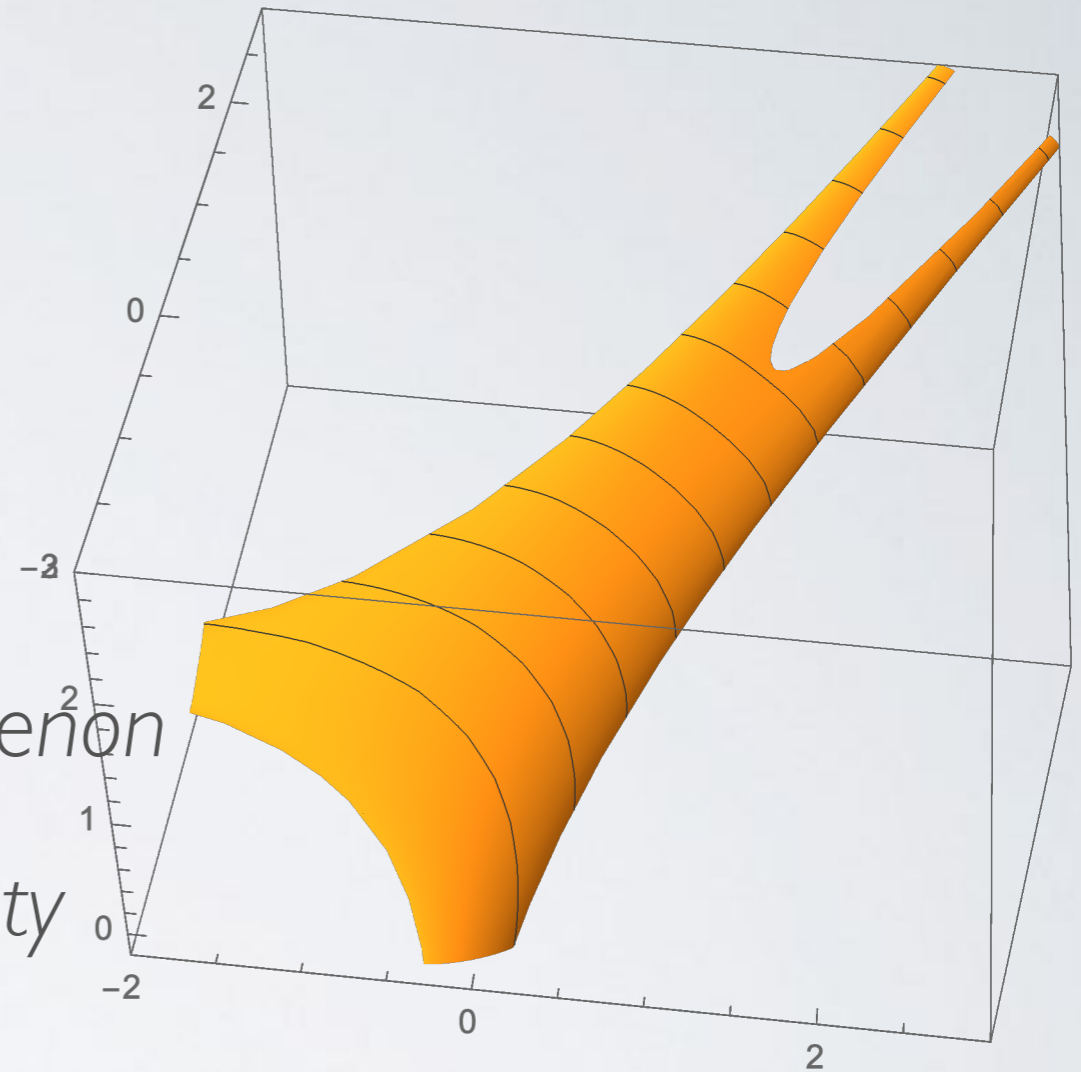
full boundary information

BPP EXAMPLE

large r

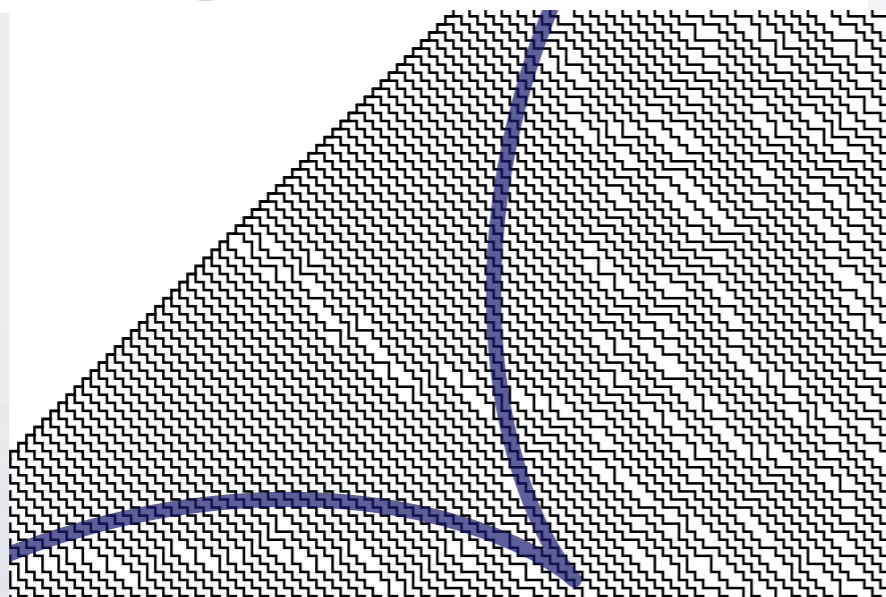


shear phenomenon
points of non-analyticity
vs algebraic curves



free energy/Wulff shape

$r=2.5$



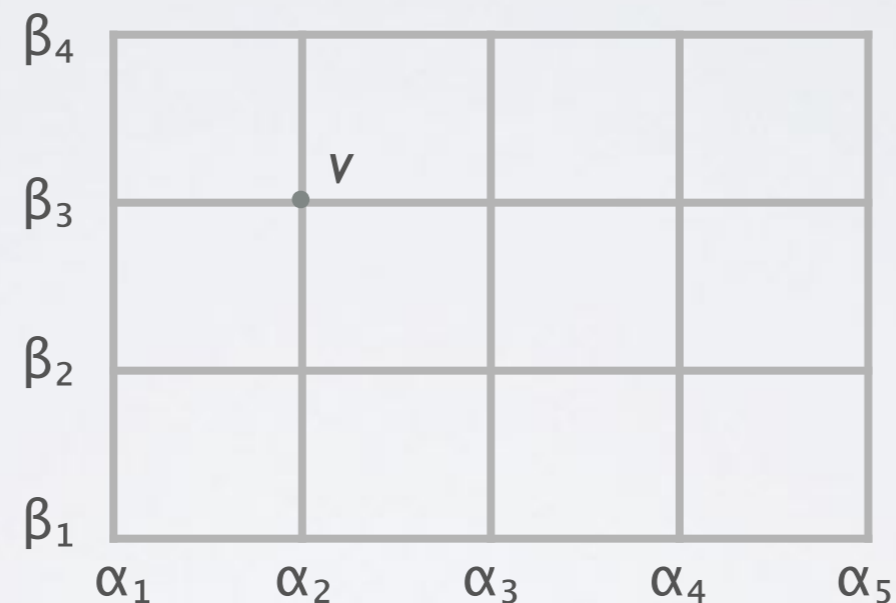
GENUS-ZERO 5-VERTEX MODEL

Kenyon-Prause

staggered model

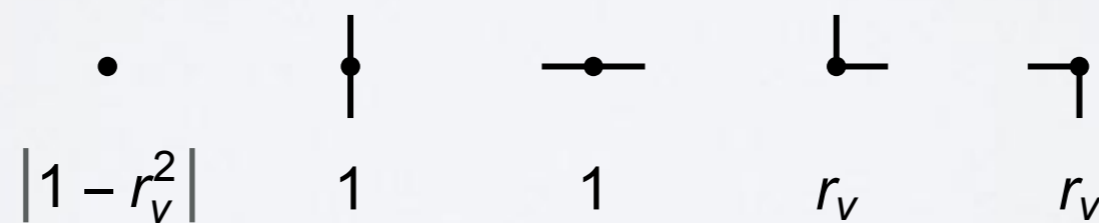
$m_1 \times m_2$ fundamental domain

$$\vec{\alpha} = (\alpha_1, \dots, \alpha_{m_1}) \quad \vec{\beta} = (\beta_1, \dots, \beta_{m_2})$$



$$v = (i, j)$$

$$r_v = \alpha_i \beta_j$$



“small r ”
(repulsive)

all $\alpha_i \beta_j < 1$

or

all $\alpha_i \beta_j > 1$

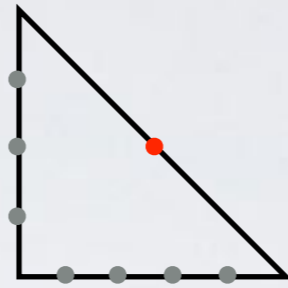
“large r ”
(attractive)

SURFACE TENSION

(small r)

convex in \mathcal{N}
 strictly convex in $\mathring{\mathcal{N}}$
 piecewise linear on $\partial\mathcal{N}$
 $\sigma|_{\mathcal{N}\setminus\mathring{\mathcal{N}}} \equiv 0$

non-uniqueness



$u \in \mathbb{H}$

conformal coordinate

$$\sqrt{\det D^2\sigma} = \frac{1}{\pi} (\arg u)^2$$

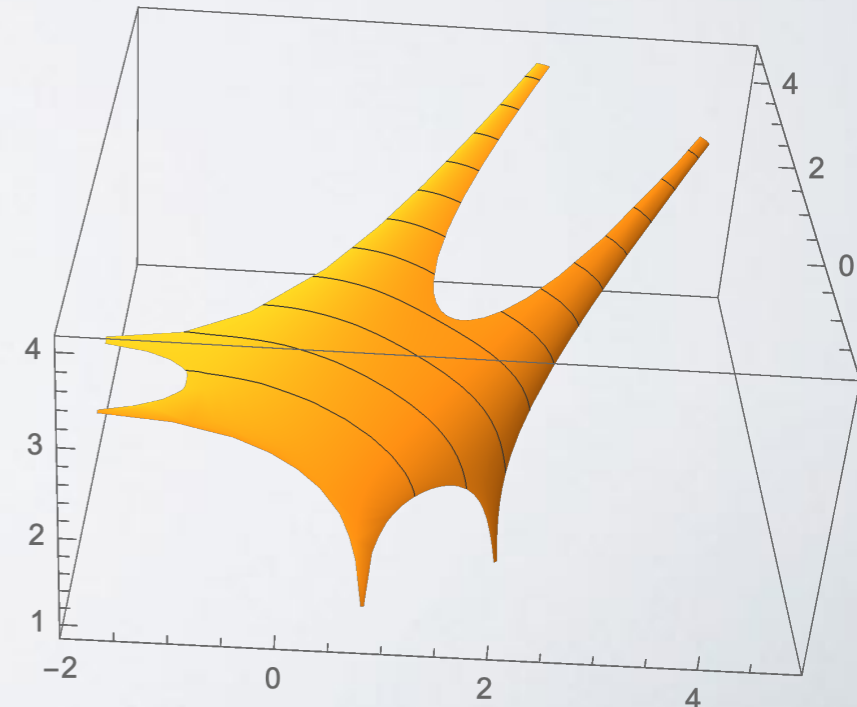
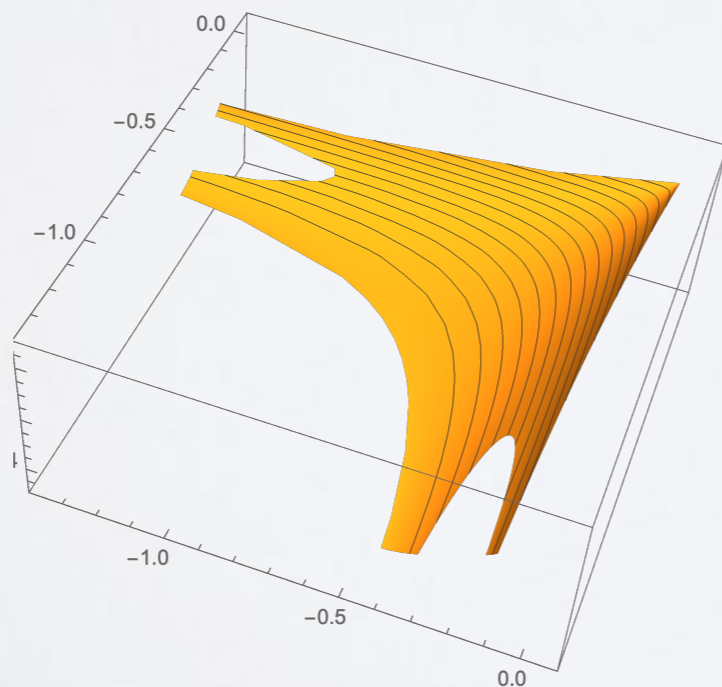
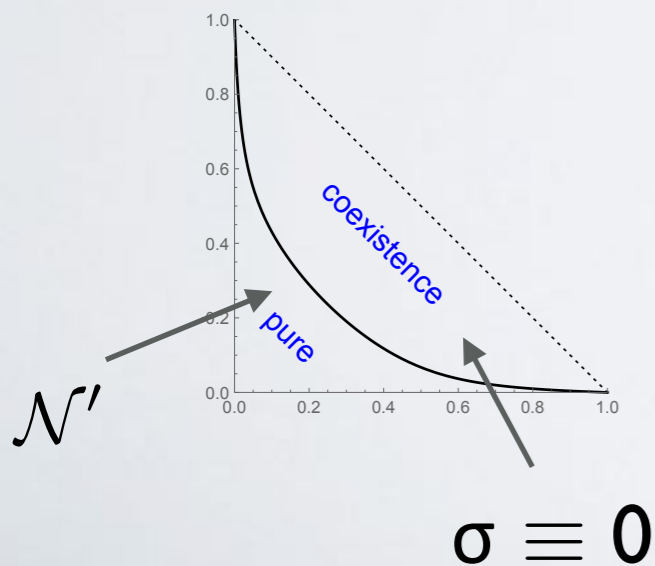
trivial potential

$$\sqrt{\det D^2\sigma} = \frac{1}{\pi} (2\pi - \arg u)^2$$

unique minimizer

(large r)

convex in \mathcal{N}
 strictly convex in $\mathring{\mathcal{N}}$
 piecewise linear on $\partial\mathcal{N}$
 slope discontinuity at $(1/2, 1/2)$



DARBOUX INTEGRABILITY

Thm: In any component of the liquid region the tangent planes to the limit shape can be parametrized by a complex ζ

$$\nabla h(x, y) = (s(u), t(u)) \quad h(x, y) - \nabla h(x, y) \cdot (x, y) = G(\zeta)/\theta(u)$$

with $u(\zeta)$ holomorphic, $G(\zeta)$ **harmonic**, $\theta(u) = \begin{cases} \arg u & (r < 1) \\ 2\pi - \arg u & (r > 1) \end{cases}$

$$\sqrt{\kappa} = \sqrt[4]{\det D^2 \sigma}$$

Corollary: In any component of the liquid region

$$(s\theta)_\zeta x + (t\theta)_\zeta y + G_\zeta - \theta_\zeta h(x, y) = 0$$

(shear phenomenon)

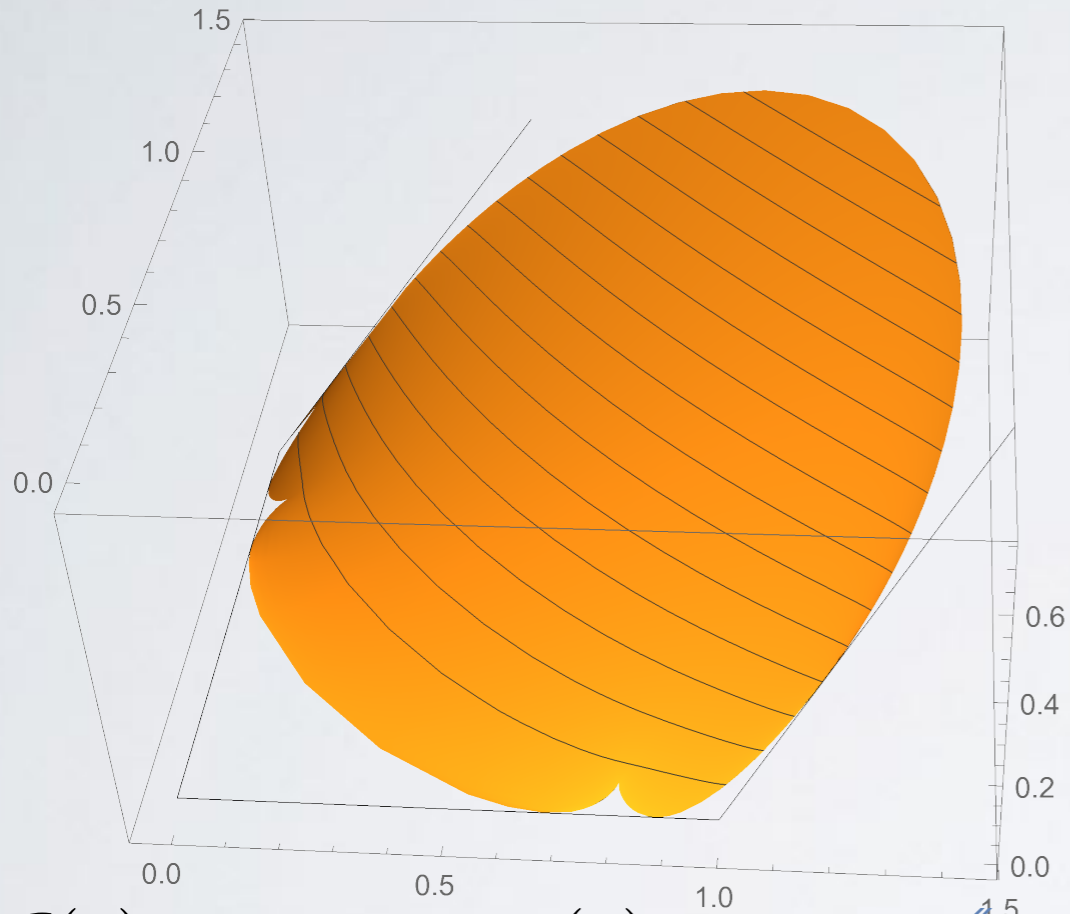
all **holomorphic** functions

2x2 EXAMPLES

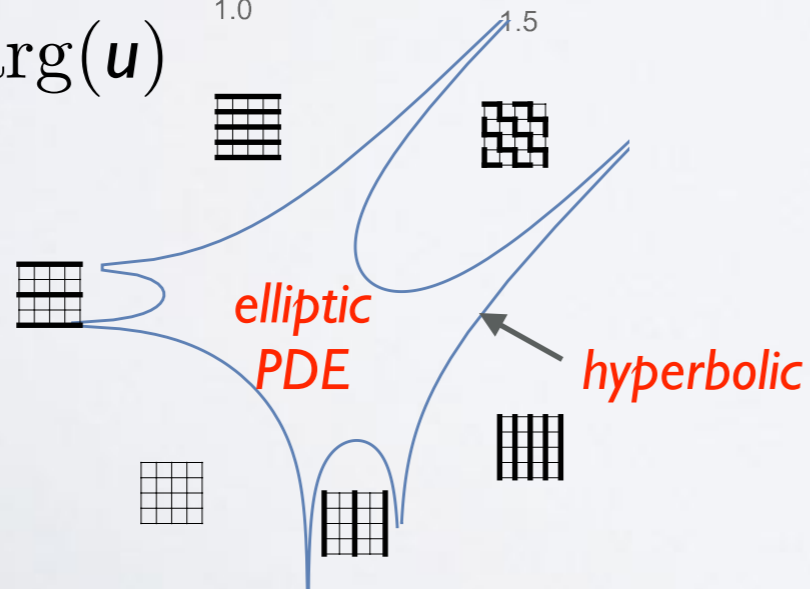
“semi-boxed plane partition”

large r

$$(\alpha_1, \alpha_2) = (\beta_1, \beta_2) = (2, 5/4)$$



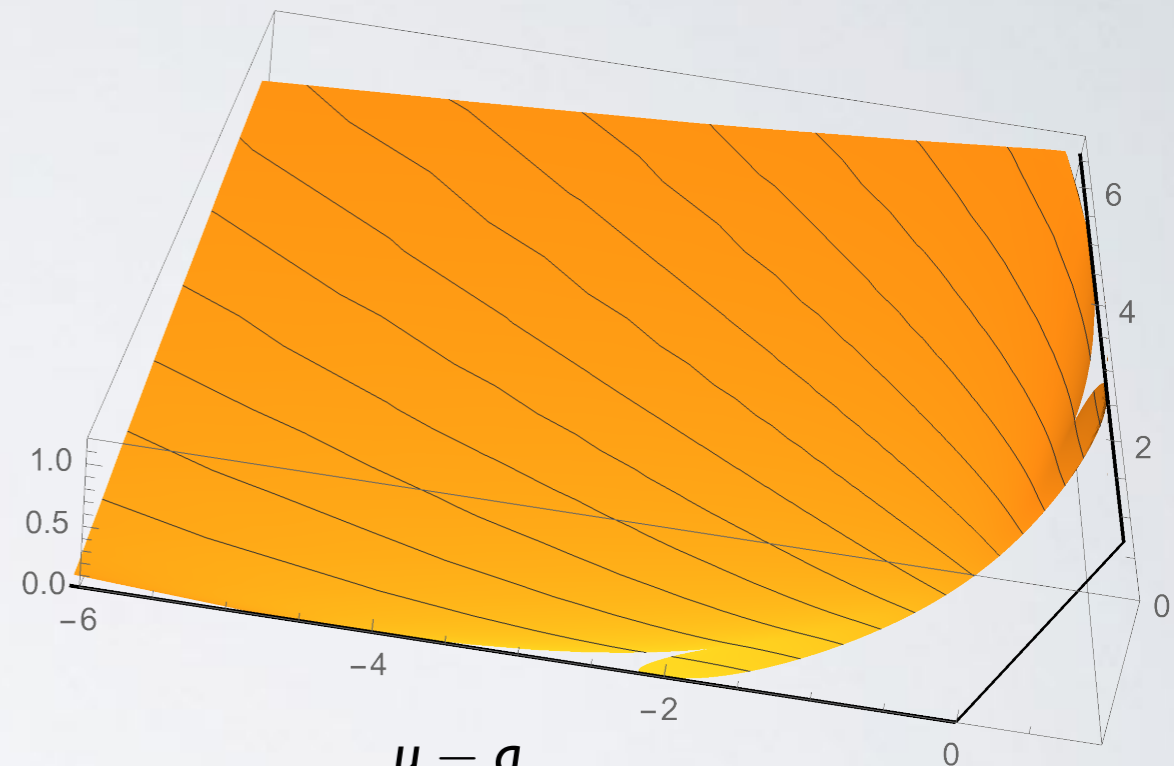
$$G(u) = -\pi - \arg(u)$$



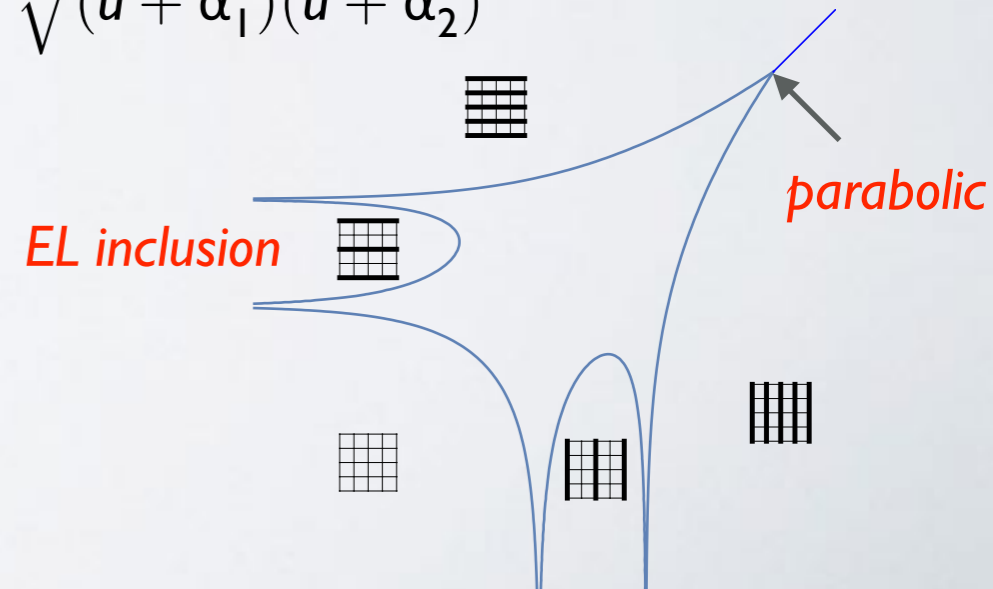
$$u = \zeta$$

small r

$$(\alpha_1, \alpha_2) = (\beta_1, \beta_2) = (4/5, 1/4)$$



$$G(u) = \arg\left(\frac{u - a}{\sqrt{(u + \alpha_1^2)(u + \alpha_2^2)}}\right)$$



DARBOUX HIERARCHY

free fermionic	constant Hessian det
5-vertex	trivial potential
six-vertex	?
...	?

LIFE BEYOND THE ARCTIC CIRCLE

