

Ex p. 7.6.

$f: \mathbb{R} \rightarrow \mathbb{R} \in C^1$ . Let  $\varphi \in X = C_0([0,1])$  be fixed.

Then, for any  $\psi \in X$

$$\|\widehat{f[\varphi + \psi]} - \widehat{f[\varphi]} - D\widehat{f[\varphi]}\psi\|_\infty$$

$$= \sup_x |f(\varphi(x) + \psi(x)) - f(\varphi(x)) - f'(\varphi(x))\psi(x)|$$

$$= \sup_x \int_0^{\psi(x)} f'(\varphi(x) + z) dz$$

$$\leq \sup_x \int_0^{\psi(x)} |f'(\varphi(x) + z) - f'(\varphi(x))| dz$$

$$\leq \sup_x \int_{-\|\psi\|_\infty}^{\|\psi\|_\infty} |f'(\varphi(x) + z) - f'(\varphi(x))| dz$$

$$\leq \sup_{|y| \leq \|\psi\|_\infty} \int_{-\|\psi\|_\infty}^{\|\psi\|_\infty} |f'(y+z) - f'(y)| dz \cdot \frac{1}{\|\psi\|_\infty} = \|\psi\|_\infty$$

$\rightarrow 0$  as  $\|\psi\|_\infty \rightarrow 0$ , for  $f'$  is unif. cont. on compact.

I see no problem in extending to higher derivatives  $D^{\mathbf{n}}$ .  
Of course we need to know  $f \in C^n$  then.

Ex on p. 20.

a)  $X = C_0(0,1) = \mathbb{Z}$ ,  $Y = \mathbb{R}$ .

b)  $F(x,y) := H(\phi, r) = \phi - g(r \hat{r}(\phi))$

$\rightarrow D_n F = D_n H(0,r) = \mathbb{1} - r g$

So, if  $\phi_k$  is both eigenvector of  $g$ ,

$D_n H(0,r) \phi_k = \phi_k - r \phi_k = \left[ \underbrace{1 - \frac{r}{(2\pi k)^2}} \right] \phi_k$

$\therefore$  If  $r \neq (2\pi k)^2$  for all  $k \in \mathbb{N}$  then

$\mathcal{R}(D_n H(0,r)) = X$ .

c) If  $r_k = (2\pi k)^2$ , then

$\mathcal{N}_k = \text{span } \phi_k$  and  $\mathcal{R}_k = \text{span} \{ \phi_j : j \neq k \}$

$\Rightarrow X = \mathcal{N}_k \oplus \mathcal{R}_k$ .

d)  $X_0^k = \{ \psi \in X : \langle \phi_k, \psi \rangle = 0 \}$ ,  $Z_0^k = \text{span } \phi_k$

e)  $P = |\phi_k\rangle \langle \phi_k|$ ,  $Q = \sum_{n \neq k} |\phi_n\rangle \langle \phi_n|$

f)  $D_{2,1} H(0,r) = -g$ , so if  $\vec{v}_0^k = \phi_k$

$-g \phi_k = -\frac{r_k}{(2\pi k)^2} \phi_k \in \mathcal{N}_k \subseteq \mathcal{R}_k^\perp$ .

Ex. on p. 7.21

a) So we pick  $r_k = (\frac{\pi}{k})^2$ ,  $k \geq 2$  where we have bifurcation from the trivial solution  $\bar{0}$ . The bifurcation, parameterized by  $\psi(s)$ , happens in the direction  $\phi_k$  in the sense (see notes p. 7.18.)

$$\frac{d}{ds} \psi(s) \Big|_{s=0} = \phi_k$$

$$\text{Then } \psi(s, x) = \underbrace{\psi(0, x)}_{=0} + s \phi_k(x) + s \varepsilon(s, x),$$

where  $\varepsilon(s) \rightarrow 0$  as  $s \rightarrow 0$ .

$$\Rightarrow \psi(s) = s (\phi_k + \varepsilon(s)).$$

So pick  $s$  small enough so that  $|\varepsilon(s)| < \frac{1}{2} \|\phi_k\|$  and you can be sure there will be both positive and neg. values.

b) If  $k=1$   $\phi_1(x) = \sin(\pi x) \approx \pi x$  for small  $x$ .

$$\text{Then } \psi(s, x) \approx s(\pi + s \varepsilon(s, x)) < 0 \Leftrightarrow \varepsilon(s, x) < -\frac{\pi}{s}$$

So, it would seem possible, if

$$\partial_x^2 \psi(s, 0) = \varepsilon(s, 0) < 0 = \lim_{x \rightarrow 0} -\frac{\pi}{s}$$

for some  $s$  "pretty small". This is no proof, simply the idea.

2nd. Ex. on p. 721

By def. of  $\tilde{y}(s)$  we have  $\tilde{\Phi}(s, \tilde{y}(s)) = 0, \forall s.$

$$\Rightarrow 0 = \frac{d}{ds} \tilde{\Phi}(s, \tilde{y}(s))$$

$$\int_0^1 \underbrace{D_1 \tilde{\Phi}(s + \hat{v}_0, \tilde{y}(s)) \hat{v}_0}_{\langle z_0^*, D_1 F(s + \hat{v}_0, \tilde{y}(s)) \hat{v}_0 \rangle z_0}$$

$$\Leftrightarrow 0 = \int_0^1 \langle z_0^*, \frac{d}{ds} D_1 F(s + \hat{v}_0, \tilde{y}(s)) \hat{v}_0 \rangle z_0$$

$$+ D_1 F(s + \hat{v}_0, \tilde{y}(s)) [\hat{v}_0, \hat{v}_0] + D_2 F(s + \hat{v}_0, \tilde{y}(s)) \dot{\tilde{y}}(s)$$

$$= \int_0^1 \langle z_0^*, D_{11} F(s + \hat{v}_0, \tilde{y}(s)) [\hat{v}_0, \hat{v}_0] \rangle z_0$$

$$+ \langle z_0^*, D_2 F(s + \hat{v}_0, \tilde{y}(s)) \hat{v}_0 \rangle z_0 \dot{\tilde{y}}(s)$$

So, at  $s=0$  we get, as we perform the <sup>initial</sup> integral

$$0 = \frac{1}{2} \langle z_0^*, D_{11} F(0, y_0) [\hat{v}_0, \hat{v}_0] \rangle z_0 + \langle z_0^*, D_2 F(0, y_0) \hat{v}_0 \rangle z_0 \dot{y}_0$$

from which we get, since  $\langle z_0^*, \frac{D_2 F(0, y_0) \hat{v}_0}{\notin \mathbb{R}} \rangle \neq 0$

$$\frac{d}{ds} \tilde{y}(s) \Big|_{s=0} = -\frac{1}{2} \frac{\langle z_0^*, D_{11} F(0, y_0) [\hat{v}_0, \hat{v}_0] \rangle}{\langle z_0^*, D_2 F(0, y_0) \hat{v}_0 \rangle} \quad \square$$

Ex. on p. 7.23

To apply the previous formula, we must calculate  $D_2 H(\phi, r) \notin D_{11} H(\phi, r)$ . Recall that

$$D_2 H(\phi, r) = \mathbb{1} - rG f'(\phi), \text{ meaning}$$

$$D_1 H(\phi, r)\psi = \psi - rG[f'(\phi), \psi] \xrightarrow{\phi \mapsto \bar{0}} \psi - r f'(\bar{0}) G[\bar{x}].$$

multiplication operator.

Now,  $D_{21} H(\phi, r)\psi = -G[\psi]$ . Finally

$$D_{11} H(\phi, r)(\psi, \psi) = -rG[(D_{\phi\phi}\hat{F})(\psi)(\psi)], \text{ where}$$

$$(D_{\phi\phi}\hat{F})(\psi)(\psi) = D_{\phi}[\hat{F}'(\phi)]\psi(\psi)$$

$$= \hat{F}''(\phi)[\psi, \psi]$$

The very same principle as before, now with function  $F'(\phi) \cdot \psi$ .

Next, we have  $\hat{v}_0 = \phi_k = \sqrt{2} \sin(k\pi \cdot)$ ,  $\hat{z}_0 = \phi_k$  too, and  $\langle \cdot, \cdot \rangle = (\cdot, \cdot)_{L^2[0,1]}$ . Also  $r_k = (k\pi)^2$

Now, since  $\phi_k$  is an eigenvector of  $G$  with e-value  $\frac{1}{(k\pi)^2}$ ,

$$D_y^1(c) = -\frac{1}{2} \frac{(\phi_k, -r_k f''(0) G \phi_k^2)}{(\phi_k, -G \phi_k)} = -\frac{r_k f''(0)}{2} \frac{(\phi_k, G \phi_k^2)}{\frac{1}{(k\pi)^2} \|\phi_k\|^2}$$

Finally, note that  $(\phi_k, G \phi_k^2) = (G \phi_k, \phi_k^2) = \frac{1}{(k\pi)^2} (\phi_k, \phi_k^2)$

$$D_y^1(c) = -\frac{1}{2} (k\pi)^2 f''(0) \underbrace{(\phi_k, \phi_k^2)}_{\int_0^1 (\sqrt{2} \sin(k\pi x))^3 dx}$$

Done.

1-11 on p 7.26

By def. of  $\tilde{w}$  and  $\tilde{\lambda}$ :

$$\underline{F_e(0, y, \tilde{w}(0, y), \tilde{\lambda}(0, y)) = 0, \quad \forall y.}$$

$$D_n F_e(0, y)(\tilde{v}_0 + \tilde{w}(0, y)) - \lambda(0, y)(\tilde{v}_0 + \tilde{w}(0, y))$$

So, applying  $\frac{d}{dy}$  on both sides gives us

$$0 = D_{21} F_e(0, y)(\tilde{v}_0 + \tilde{w}) + D_n F_e(0, y) \partial_y \tilde{w}(0, y) - \partial_y \lambda(0, y)(\tilde{v}_0 + \tilde{w}) \\ - \lambda(0, y) \partial_y \tilde{w}(0, y).$$

Evaluating at  $y = y_0$  (recall  $\tilde{w}(0, y_0) = 0$ )

$$0 = D_{21} F_e(0, y_0) \tilde{v}_0 + D_n F_e(0, y_0) \partial_y \tilde{w}(0, y_0) - \partial_y \lambda(0, y_0) \tilde{v}_0 \\ - \lambda_0 \partial_y \tilde{w}(0, y_0)$$

$$= [D_{21} F_e(0, y_0) - \partial_y \lambda(0, y_0)] \tilde{v}_0 + [D_n F_e(0, y_0) - \lambda_0] \partial_y \tilde{w}(0, y_0)$$

Applying  $\langle z_0^*, \cdot \rangle$  to both sides:

$$0 = \langle z_0^*, D_{21} F_e(0, y_0) \tilde{v}_0 \rangle - \partial_y \lambda(0, y_0) + \underbrace{\langle z_0^*, [D_n F_e(0, y_0) - \lambda_0] \partial_y \tilde{w}(0, y_0) \rangle}_{\text{span}\{\tilde{v}_0\}} = 0$$

$$\therefore \partial_y \lambda(0, y_0) = \lambda'(y_0) = \langle z_0^*, D_{21} F_e(0, y_0) \tilde{v}_0 \rangle.$$

2nd. Ex. on p. 7.26

Well, if  $D_n F(0, y_0) \tilde{v}_0 \notin \mathbb{R}$ , then

$$\langle z_0^*, D_n F(0, y_0) \tilde{v}_0 \rangle \neq 0,$$

which by prev. ex. means  $\lambda_+(y_0) \neq 0$ .

3rd. ——— 7.2.6

Precisely the same principle as in exercises just above:

First differentiate the defining eq. of the non-trivial-branch-

$$\text{eigenvalue } F_c(x(a), y(a), w(x(a), y(a)), \lambda(x(a), y(a))) = 0$$

w.r.t.  $\lambda$ .

Secondly, apply  $\langle z_0^* \rangle$  to the resulting eq. to eliminate terms which have values in  $\mathbb{R}$ .

Done. (Ask the details from me if you really need to see them.  
They won't make you wiser I think.)

4th. ———

By 3rd. ex.

$$\lambda_n'(0) = \lambda_+(y_0) \underbrace{D \tilde{y}(0)} + \langle z_0^*, D_n F(0, y_0) [v_0, v_0] \rangle$$

$$= \frac{-\frac{1}{2} \langle z_0^*, D_n F(0, y_0) [v_0, v_0] \rangle}{\frac{1}{2} \langle z_0^*, D_n F(0, y_0) v_0 \rangle}$$

$$\langle z_0^*, D_n F(0, y_0) v_0 \rangle$$

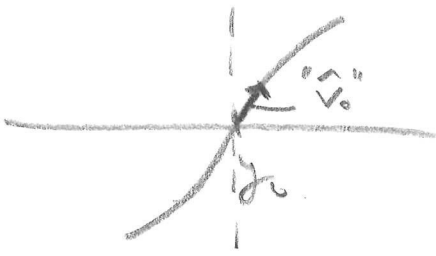
$$D \tilde{y}(0)$$

$$= \frac{1}{2} \langle z_0^*, D_n F(0, y_0) [v_0, v_0] \rangle = -\frac{1}{2} \frac{\langle z_0^*, D_n F(0, y_0) [v_0, v_0] \rangle}{\lambda_+(y_0)} \cdot \lambda_+(y_0)$$

Done.

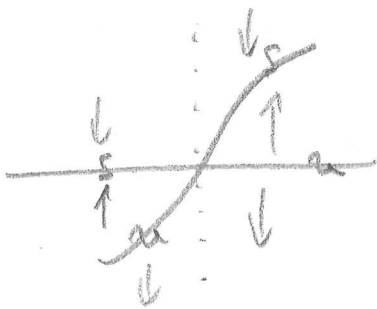
5.26 - u - 7.26

If  $D\hat{y}'(0) > 0$  the situation looks like this:

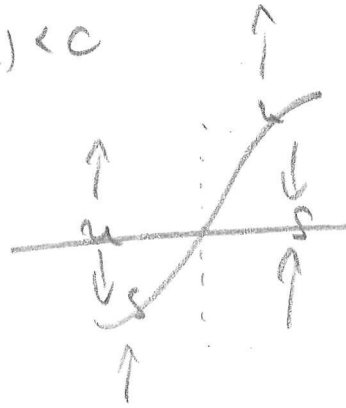


In which case the stability can look like

$\lambda_1(y_0) > 0$

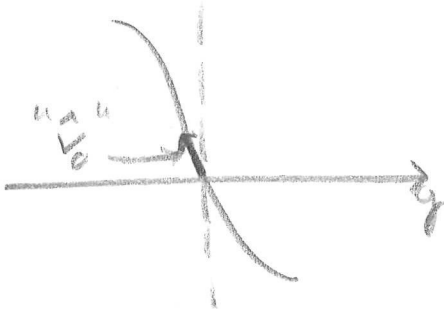


$\lambda_1(y_0) < 0$



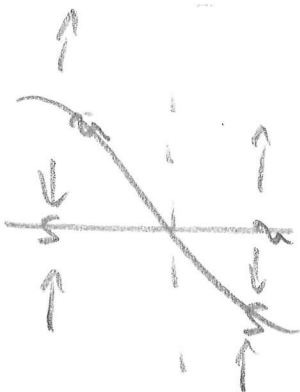
or

If  $D\hat{y}'(0) < 0$  the situation looks like:

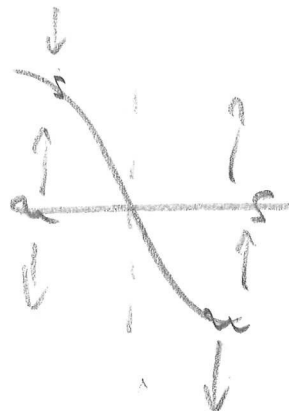


stability can look like:

$\lambda_1(y_0) > 0$



$\lambda_1(y_0) < 0$



or