One-dimensional deconvolution will serve as a basic example throughout Part I of the book. Two-dimensional deconvolution is a project topic in Section 10.

### 2.1.1 Continuum model for one-dimensional convolution

We build a computational model for one-dimensional convolution with periodic boundary conditions. We consider 1-periodic functions $f: \mathbb{R} \rightarrow \mathbb{R}$ satisfying $f(x)=f(x+n)$ with any integer $n \in \mathbb{Z}$. Essentially the function $f$ is defined on an interval of length 1 such as $[0,1]$ or $\left[-\frac{1}{2}, \frac{1}{2}\right]$ with the endpoints identified; another way of thinking about this is to consider $f(x)$ defined on a circle with radius $(2 \pi)^{-1}$ and $x$ being the arc length variable.

The reason for considering periodic functions is that we can avoid some technicalities related to boundary conditions that would obscure the main message about ill-posedness. Also, the Fourier transform and the wavelet transform are easily defined and implemented in the periodic setting.

The continuum measurement model concerns a 1-periodic signal $f: \mathbb{R} \rightarrow$ $\mathbb{R}$ blurred by a 1-periodic point spread function (PSF) $\psi$. Other common names for the point spread function include device function, impulse response, blurring kernel, convolution kernel and transfer function.

Let us first construct the PSF using a building block $\psi_{0}$ defined in the interval $[-a, a] \subset \mathbb{R}$ with some constant $0<a<1 / 2$ :

$$
\begin{equation*}
\psi_{0}(x)=C_{a}(x+a)^{2}(x-a)^{2}, \text { for }-a \leq x \leq a \tag{2.1}
\end{equation*}
$$

where the constant $C_{a}:=\left(\int_{-a}^{a}(x+a)^{2}(x-a)^{2} d x\right)^{-1}$ is chosen to enforce the following normalization:

$$
\begin{equation*}
\int_{-a}^{a} \psi_{0}(x) d x=1 \tag{2.2}
\end{equation*}
$$

The periodic point spread function is defined by copying $\psi_{0}(x)$ to every interval $[n-a, n+a]$ with $n \in \mathbb{Z}$ and setting $\psi(x)$ to zero outside those intervals. The resulting $\psi$ is a non-negative and even function:

$$
\begin{equation*}
\psi(x) \geq 0 \quad \text { and } \quad \psi(x)=\psi(-x) \quad \text { for all } x \in \mathbb{R} \tag{2.3}
\end{equation*}
$$

See Figure 2.1 for a plot of the point spread function with $a=0.04$.
We remark that instead of (2.2) one often requires $\int_{-a}^{a} \psi_{0}(x)^{2} d x=1$. However, we prefer (2.2) since then constant functions remain unchanged in convolution with $\psi$; this will be convenient below when we compare plots of reconstructions to the plot of the true signal by showing them in the same figure.


Figure 2.1: Point spread function according to (2.4) with $a=0.04$ for onedimensional convolution. Left: the continuously differentiable building block $\psi_{0}(x)$ used for constructing the periodic PSF. Right: the periodic PSF $\psi(x)$.



Figure 2.2: Effect of convolution on a piecewise continuous function. Left: target function $f(x)$. Right: the function $(\psi * f)(x)$.

Definition 2.1.1 The continuum model of convolution, or blurring, is given by the following integral:

$$
\begin{equation*}
(\psi * f)(x)=\int_{-a}^{a} \psi\left(x^{\prime}\right) f\left(x-x^{\prime}\right) d x^{\prime} \tag{2.4}
\end{equation*}
$$

Note that formula (2.4) is not of the form (1.1) since the left hand side is not a $k$-dimensional vector. However, suppose the function $f$ is defined on an interval $[b, b+1]$, and assume that we have a device that measures the values of the convolution function $(\psi * f)(x)$ at a collection of $k$ equally spaced points $\tilde{x}_{1}=b, \tilde{x}_{2}=b+\frac{1}{k}, \tilde{x}_{3}=b+\frac{2}{k}, \ldots, \tilde{x}_{k}=b+\frac{k-1}{k}$ and define

$$
\begin{equation*}
\mathbf{m}:=\left[(\psi * f)\left(\tilde{x}_{1}\right),(\psi * f)\left(\tilde{x}_{2}\right), \ldots,(\psi * f)\left(\tilde{x}_{k}\right)\right]^{T} \in \mathbb{R}^{k} \tag{2.5}
\end{equation*}
$$

Then $\mathcal{A} f=\mathbf{m}$ is of the form (1.1).

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### 2.1.2 Discrete convolution model

Next we need to discretize the continuum model to arrive at a finite-dimensional measurement model of the form (1.3). Define

$$
\begin{equation*}
x_{j}=b+\frac{j-1}{n} \quad \text { for } j=1,2 \ldots, n \text {; } \tag{2.6}
\end{equation*}
$$

then the 1-periodic real-valued function $f(x)$ is represented by a vector $\mathbf{f}$ containing values at the grid points:

$$
\begin{equation*}
\mathbf{f}=\left[\mathbf{f}_{1}, \mathbf{f}_{2}, \ldots, \mathbf{f}_{n}\right]^{T}=\left[f\left(x_{1}\right), f\left(x_{2}\right), \ldots, f\left(x_{n}\right)\right]^{T} \in \mathbb{R}^{n} \tag{2.7}
\end{equation*}
$$

Furthermore, denote $\Delta x:=x_{2}-x_{1}=1 / n$.
We can approximate the integral appearing in (2.4) by numerical quadrature. For any reasonably well-behaved function $g:[b, b+1] \rightarrow \mathbb{R}$ we have

$$
\begin{equation*}
\int_{b}^{b+1} g(x) d x \approx \Delta x \sum_{j=1}^{n} g\left(x_{j}\right) \tag{2.8}
\end{equation*}
$$

the approximation becoming better as $n$ increases.
For convenience, let us take $k=n$ and measure the convolution at the same points (2.6) as where the unknown function $f$ is sampled. This is not necessary in general, but it will lead to a square-shaped matrix $A$, making it easy to illustrate naïve reconstructions and inverse crimes.

Let us construct an $n \times n$ matrix $A$ so that $A \mathbf{f} \in \mathbb{R}^{k}$ approximates $\mathcal{A} f$ defined by (2.4). We define a discrete PSF denoted by

$$
\mathbf{p}=\left[\mathbf{p}_{-\nu}, \mathbf{p}_{-\nu+1}, \ldots, \mathbf{p}_{-1}, \mathbf{p}_{0}, \mathbf{p}_{1}, \ldots, \mathbf{p}_{\nu-1}, \mathbf{p}_{\nu}\right]^{T}
$$

as follows. Recall that $\psi_{0}(x) \equiv 0$ for $|x|>a>0$. Take $\nu>0$ to be the smallest integer satisfying the inequality $(\nu+1) \Delta x>a$ and set

$$
\widetilde{\mathbf{p}}_{j}=\psi_{0}(j \Delta x) \quad \text { for } \quad j=-\nu, \ldots, \nu .
$$

For example, with $a=0.04$ as in Figure 2.1 and $n=64$, we get $\nu=2$. By (2.8) the normalization condition (2.2) almost holds: $\Delta x \sum_{j=-\nu}^{\nu} \widetilde{p}_{j} \approx 1$. However, in practice it is a good idea to normalize the discrete PSF explicitly by the formula

$$
\begin{equation*}
\mathbf{p}=\left(\Delta x \sum_{j=-\nu}^{\nu} \widetilde{\mathbf{p}}_{j}\right)^{-1} \widetilde{\mathbf{p}} ; \tag{2.9}
\end{equation*}
$$

then it follows that

$$
\begin{equation*}
\Delta x \sum_{j=-\nu}^{\nu} \mathbf{p}_{j}=1 \tag{2.10}
\end{equation*}
$$

Now

$$
\begin{aligned}
\int_{-a}^{a} \psi\left(x^{\prime}\right) f\left(x_{j}-x^{\prime}\right) d x^{\prime} & \approx \Delta x \sum_{\ell=-\nu}^{\nu} \psi\left(x_{\ell}\right) f\left(x_{j}-x_{\ell}\right) \\
& \approx \Delta x \sum_{\ell=-\nu}^{\nu} \mathbf{p}_{\ell} \mathbf{f}_{j-\ell}
\end{aligned}
$$

Hence discrete convolution is defined by the formula

$$
\begin{equation*}
(\mathbf{p} * \mathbf{f})_{j}=\sum_{\ell=-\nu}^{\nu} \mathbf{p}_{\ell} \mathbf{f}_{j-\ell} \tag{2.11}
\end{equation*}
$$

where $\mathbf{f}_{j-\ell}$ is defined using periodic boundary conditions for the cases $j-\ell<$ 1 and $j-\ell>n$. Then

$$
\begin{equation*}
\Delta x(\mathbf{p} * \mathbf{f}) \approx \mathcal{A} f \tag{2.12}
\end{equation*}
$$

and we define the measurement vector $\mathbf{m}=\left[\mathbf{m}_{1}, \ldots \mathbf{m}_{k}\right]^{T}$ by

$$
\begin{equation*}
\mathbf{m}_{j}=\Delta x(\mathbf{p} * \mathbf{f})_{j}+\varepsilon_{j} \tag{2.13}
\end{equation*}
$$

We would like to write formula (2.13) using a matrix $A$ so that we would arrive at the desired model (1.3). To this end, set

$$
\left[\begin{array}{c}
\mathbf{m}_{1} \\
\vdots \\
\mathbf{m}_{k}
\end{array}\right]=\left[\begin{array}{ccc}
a_{11} & \cdots & a_{1 n} \\
\vdots & \ddots & \vdots \\
a_{k 1} & \cdots & a_{k n}
\end{array}\right]\left[\begin{array}{c}
\mathbf{f}_{1} \\
\vdots \\
\mathbf{f}_{n}
\end{array}\right]+\left[\begin{array}{c}
\varepsilon_{1} \\
\vdots \\
\varepsilon_{k}
\end{array}\right]
$$

The answer is to build a circulant matrix having the elements of $\mathbf{p}$ appearing systematically on every row of $A$.

Let us illustrate the structure of the convolution matrix $A$ by an example in the case $n=64$. As observed above, if $a=0.04$ then $\nu=2$, and the PSF takes the form $p=\left[\begin{array}{lllll}p_{-2} & p_{-1} & p_{0} & p_{1} & p_{2}\end{array}\right]^{T}$. According to (2.11) we have

$$
\begin{aligned}
(\mathbf{p} * \mathbf{f})_{1} & =\mathbf{p}_{0} \mathbf{f}_{1}+\mathbf{p}_{-1} \mathbf{f}_{2}+\mathbf{p}_{-2} \mathbf{f}_{3}+\mathbf{p}_{2} \mathbf{f}_{n-1}+\mathbf{p}_{1} \mathbf{f}_{n} \\
(\mathbf{p} * \mathbf{f})_{2} & =\mathbf{p}_{1} \mathbf{f}_{1}+\mathbf{p}_{0} \mathbf{f}_{2}+\mathbf{p}_{-1} \mathbf{f}_{3}+\mathbf{p}_{-2} \mathbf{f}_{4}+\mathbf{p}_{2} \mathbf{f}_{n} \\
(\mathbf{p} * \mathbf{f})_{3} & =\mathbf{p}_{2} \mathbf{f}_{1}+\mathbf{p}_{1} \mathbf{f}_{2}+\mathbf{p}_{0} \mathbf{f}_{3}+\mathbf{p}_{-1} \mathbf{f}_{4}+\mathbf{p}_{-2} \mathbf{f}_{5} \\
& \vdots \\
(\mathbf{p} * \mathbf{f})_{n} & =\mathbf{p}_{-1} \mathbf{f}_{1}+\mathbf{p}_{-2} \mathbf{f}_{2}+\mathbf{p}_{2} \mathbf{f}_{n-2}+\mathbf{p}_{1} \mathbf{f}_{n-1}+\mathbf{p}_{0} \mathbf{f}_{n}
\end{aligned}
$$

Consequently the matrix $A$ looks like this:

$$
A=\Delta x\left[\begin{array}{lllllllll}
\mathbf{p}_{0} & \mathbf{p}_{-1} & \mathbf{p}_{-2} & 0 & 0 & 0 & \cdots & \mathbf{p}_{2} & \mathbf{p}_{1}  \tag{2.14}\\
\mathbf{p}_{1} & \mathbf{p}_{0} & \mathbf{p}_{-1} & \mathbf{p}_{-2} & 0 & 0 & \cdots & 0 & \mathbf{p}_{2} \\
\mathbf{p}_{2} & \mathbf{p}_{1} & \mathbf{p}_{0} & \mathbf{p}_{-1} & \mathbf{p}_{-2} & 0 & \cdots & 0 & 0 \\
0 & \mathbf{p}_{2} & \mathbf{p}_{1} & \mathbf{p}_{0} & \mathbf{p}_{-1} & \mathbf{p}_{-2} & \cdots & 0 & 0 \\
\vdots & & & & \ddots & & & & \\
\vdots & & & & & \ddots & & & \\
0 & 0 & \cdots & \mathbf{p}_{2} & \mathbf{p}_{1} & \mathbf{p}_{0} & \mathbf{p}_{-1} & \mathbf{p}_{-2} \\
\mathbf{p}_{-2} & 0 & \cdots & 0 & \mathbf{p}_{2} & \mathbf{p}_{1} & \mathbf{p}_{0} & \mathbf{p}_{-1} \\
\mathbf{p}_{-1} & \mathbf{p}_{-2} & \cdots & 0 & 0 & \mathbf{p}_{2} & \mathbf{p}_{1} & \mathbf{p}_{0}
\end{array}\right] ;
$$

note the systematic band-diagonal structure, which characterizes $A$ as a circulant matrix. Linear systems involving circulant matrices can be quickly solved using Fast Fourier Transforms, a topic we will return to later.

Returning to the general case of $\mathbf{p}$ defined by (2.9), the approximation formula (2.12) can be written in the form

$$
\begin{equation*}
A \mathbf{f} \approx \mathcal{A} f \tag{2.15}
\end{equation*}
$$

Figure 2.3 shows data computed by the discrete model $A \mathbf{f}$ and compares the result to the continuous data $(\psi * f)(x)$ defined by (2.4).

Now let's add a little noise to the data. For example, we might take $k=64=n$ and construct the measurement noise in a probabilistic manner by taking a realization of a random vector with 64 independently distributed Gaussian elements having standard deviation $\sigma=0.01 \cdot \max |f(x)|$. This corresponds to a relative noise level of $1 \%$.

### 2.1.3 Naïve deconvolution and inverse crimes

We illustrate numerically the failure of the following naïve reconstruction attempt:

$$
\begin{equation*}
\mathbf{f} \approx A^{-1} \mathbf{m} \approx A^{-1}(A \mathbf{f}+\varepsilon)=\mathbf{f}+A^{-1}(\varepsilon) \tag{2.16}
\end{equation*}
$$

In the case of no added noise $(\varepsilon=0)$ we use the data shown in the left plot of Figure 2.4 and get the left plot in Figure 2.5. The naïve reconstruction seems perfect! However, there is a catch. This apparently accurate reconstruction is not to be trusted; it is an example of an inverse crime. We will show how to avoid inverse crime in Section 2.1.4.

If we apply naïve reconstruction (2.16) to the slightly noisy data shown in the right plot of Figure 2.4, we get the result shown in the right plot

