

Lecture 01: Existence and uniqueness theorems and flows

Introduction

The expounded material can be found in

- ch. 2 of N. Berglund's lecture notes [1]
- ch. 1 of [2]

1 Existence and uniqueness theorems

Let $D \subset \mathbb{R}^n$. For any $\mathbf{x} \in D$ we consider the ordinary differential equation

$$\dot{\mathbf{x}}_t = \mathbf{f}(\mathbf{x}_t) \quad (1.1a)$$

$$\mathbf{x}_0 = \mathbf{x} \quad (1.1b)$$

driven by the vector field

$$\mathbf{f}: D \mapsto \mathbb{R}^n \quad (1.2)$$

We denote by $C^r(D; \mathbb{R}^n) \equiv C^r(D)$ the space of \mathbb{R}^n -valued vector fields with support over D which are r -times continuously differentiable. We then posit $\mathbf{f} \in C^r(D)$ for some $r \geq 0$.

A differential problem (1.1) specified by an *evolution law* (1.1a) and an initial data (1.1b) is also often called a *Cauchy problem*. Equation (1.1b) fixes the *initial data* or *initial condition* for $t = 0$.

1.1 Existence

Theorem 1.1 (Peano–Cauchy). *Let $\mathbf{f} \in C(D)$ i.e. let \mathbf{f} be a continuous vector field over D . For every $\mathbf{x} \in D$, there exists at least one local solution of (1.1) through \mathbf{x} , that is, there is an open interval I such that $0 \in I$ and a function*

$$\mathbf{x}_t: I \mapsto D \quad (1.3)$$

such that (1.1a) and (1.1b) are satisfied for $t \in I$.

The support I of a solution \mathbf{x}_t can be extended as long as \mathbf{x}_t remains finite.

Proposition 1.1. *Every solution*

$$\mathbf{x}_t: I \mapsto D \quad (1.4)$$

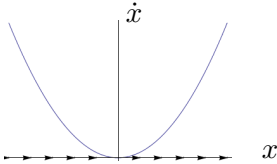
starting from $\mathbf{x}_0 = \mathbf{x}$ can be continued to a maximal interval of existence $I_{\max} = (t_{\min}, t_{\max})$ comprising the time $t = 0$ when the initial condition is assigned. If $t_{\max} < \infty$ or $t_{\min} > -\infty$, then for any compact $K \subset D$ exists a time $t \in I_{\max}$ such that $\mathbf{x}_t \notin K$.

In other words, the proposition states that either there is a time t for which

$$\mathbf{x}_t \in \partial D \tag{1.5}$$

i.e. reaches the boundary of D or the solution diverges. Very simple vector fields (1.2) may bring about explosions.

Example 1.1. For example

$$\dot{x}_t = \alpha x_t^2 \tag{1.6}$$


admits the solution

$$x_t = \frac{1}{\frac{1}{x} - \alpha t} \tag{1.7}$$

For any $x \in \mathbb{R}_+$ the solution diverges in finite time:

$$t_{\text{expl}} = \frac{1}{\alpha x} \tag{1.8}$$

Note the physical dimension of the problem

$$\begin{aligned} [x] &= [space] \\ [t] &= [time] \end{aligned} \Rightarrow [\alpha] = \left[\frac{1}{time \times space} \right] = -[space] - [time] \tag{1.9}$$

The vector field (1.2) is (“time”)-autonomous. This means that it does not depend explicitly upon the “time” variable with respect to which we differentiate \mathbf{x}_t .

Proposition 1.2. *Let*

$$\mathbf{x}_t = \phi(t, \mathbf{x}) \equiv \phi_t(\mathbf{x}) \tag{1.10}$$

be the unique solution of (1.1). Then $\tilde{\mathbf{x}}_t = \mathbf{x}_{t-t_0}$ is the unique solution of the Cauchy problem evolving according to (1.1a) from the initial data

$$\tilde{\mathbf{x}}_{t_0} = \mathbf{x} \tag{1.11}$$

assigned at $t_0 \neq 0$.

Proof. By (1.10)

$$\tilde{\mathbf{x}}_t = \phi(t - t_0, \mathbf{x}) \equiv \phi_{t-t_0}(\mathbf{x}) \tag{1.12}$$

hence for any $t_0 \in \mathbb{R}$

$$\dot{\mathbf{x}}_{t-t_0} = \frac{d\phi}{dt}(t - t_0, \mathbf{x}) = \frac{d\phi}{du}(u, \mathbf{x}) \Big|_{u=t-t_0} = \mathbf{f}(\mathbf{x}_u)|_{u=t-t_0} = \mathbf{f}(\mathbf{x}_{t-t_0}) \tag{1.13}$$

which yields the claim. □

The important consequence of the above proposition is that solutions of autonomous differential equations depend only upon the time *elapsed* from the moment t_0 the initial condition is assigned but not upon t_0 itself. Non-autonomous systems such as

$$\dot{\mathbf{x}}_t = \mathbf{f}(\mathbf{x}_t, t) \tag{1.14}$$

can always be regarded as autonomous systems in a space with one dimension more. Upon defining

$$u = t \quad \Rightarrow \quad \dot{f}(t) = \frac{df}{dt}(t) \equiv \frac{df}{du}(u) \tag{1.15}$$

we can write

$$\dot{\mathbf{x}}_u = \mathbf{f}(\mathbf{x}_t, t) \tag{1.16a}$$

$$\dot{t}_u = 1 \tag{1.16b}$$

Hence if we define $E = D \times I \subset \mathbb{R}^n \times \mathbb{R} = \mathbb{R}^{n+1}$ with

$$\mathbf{y} = \begin{bmatrix} \mathbf{x} \\ t \end{bmatrix} \in E \quad \& \quad \mathbf{g} = \begin{bmatrix} \mathbf{f} \\ 1 \end{bmatrix} : E \mapsto \mathbb{R}^{n+1} \tag{1.17}$$

we arrive to the *autonomous* dynamical system

$$\dot{\mathbf{y}}_u = \mathbf{g}(\mathbf{y}_u) \tag{1.18a}$$

$$\mathbf{y}_0 = \begin{bmatrix} \mathbf{x} \\ 0 \end{bmatrix} \tag{1.18b}$$

1.2 Uniqueness

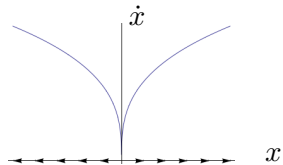
Theorem 1.2 (Picard–Lindelöf). *Suppose that $\mathbf{f} \in C(D)$ and locally Lipschitzian that is, for every compact $K \subset D$ there exists a constant $L_K \in \mathbb{R}_+$ such that*

$$\| \mathbf{f}(\mathbf{x}) - \mathbf{f}(\mathbf{y}) \| \leq L_K \| \mathbf{x} - \mathbf{y} \| \tag{1.19}$$

for all \mathbf{x} and \mathbf{y} belonging to K . Then there is a unique solution of (1.1) for each initial data in D

We notice that a differentiable function satisfies the Lipschitz condition.

Example 1.2. The ordinary differential equation

$$\dot{x}_t = \alpha |x_t|^{1/3} \tag{1.20}$$


is driven by a vector field non differentiable in zero. In general we can integrate (1.20) and obtain

$$x_t = \left(x^{2/3} + \frac{2}{3} \alpha t \right)^{3/2} \tag{1.21}$$

If, however, we assign the initial data to be $x = 0$ we can find the solutions (1.21), and also

$$x_t = 0 \tag{1.22}$$

We can combine these solutions together and generate a one parameter t_* family of solutions taking the form

$$x_t = \begin{cases} 0 & 0 \leq t \leq t_* \\ \left(x^{2/3} + \frac{2}{3} \alpha t \right)^{3/2} & t_* < t \end{cases} \tag{1.23}$$

An important implication of the Picard–Lindelöf theorem is that *for any finite time* different solutions of (1.1) driven by a smooth vector field cannot intersect. Namely we have

Proposition 1.3. *For any $\mathbf{x} \in D$ there exists only one solution of (1.1) with $\mathbf{f} \in C^r(D)$, $r \geq 1$ passing through this point.*

Proof. Let $\mathbf{x}_t^{(i)}$, $i = 1, 2$ solutions of (1.1) satisfying

$$\begin{aligned}\mathbf{x}_{t_1}^{(1)} &= \mathbf{x} \\ \mathbf{x}_{t_2}^{(2)} &= \mathbf{x}\end{aligned}\tag{1.24}$$

Since (1.1) is autonomous

$$\tilde{\mathbf{x}}_t = \mathbf{x}_{t-(t_1-t_2)}\tag{1.25}$$

is also a solution of (1.1) satisfying

$$\tilde{\mathbf{x}}_{t_1} = \mathbf{x}_{t_2}^{(2)} = \mathbf{x} = \mathbf{x}_{t_1}^{(1)}\tag{1.26}$$

We thus see that our starting assumption leads to a contradiction with the Picard–Lindelöf theorem. \square

1.3 Smooth dependence upon initial data

Let us now suppose that the vector field \mathbf{f} in (1.1a) carries a *parametric* dependence upon a set of parameters. Specifically let us suppose that

$$\mathbf{f} \in C^r(D_1 \times D_2; \mathbb{R}^n)\tag{1.27}$$

We will denote by \mathbf{x} coordinates over $D_1 \subset \mathbb{R}^n$ and by \mathbf{y} coordinates over $D_2 \subset \mathbb{R}^m$ so that

$$\mathbf{f}(\mathbf{x}, \mathbf{y}) \in \mathbb{R}^n\tag{1.28}$$

Note that we are not assuming any relation between the integers m and n .

Theorem 1.3 (Picard–Lindelöf). *Suppose that $r \geq 1$ in (1.27). In words: \mathbf{f} is at least once continuously differentiable with respect to all its arguments. Let I the maximal interval of existence of the unique solution of the ordinary differential equation*

$$\dot{\mathbf{x}}_t = \mathbf{f}(\mathbf{x}_t; \mathbf{y})\tag{1.29a}$$

$$\mathbf{x}_0 = \mathbf{x}\tag{1.29b}$$

Under these hypotheses we can write the unique solution of (1.29) as

$$\mathbf{x}_t = \phi(t, \mathbf{x}; \mathbf{y}) \equiv \phi_t(\mathbf{x}; \mathbf{y})\tag{1.30}$$

where

$$\phi_t \in C^r(I \times D_1 \times D_2, D_1)\tag{1.31}$$

Thus, under smoothness assumptions over the vector field \mathbf{f} we are entitled to consider derivatives of solutions (1.29) with respect to *all* its parametric dependencies. Moreover, if we generically denote by \mathbf{x}'_t the derivative of \mathbf{x}_t with respect to any of its parameters the same notation may apply to three different cases:

1. $\xi'_t \in \mathbb{R}^{n \times m}$ if

$$\xi'_t = \partial_{\mathbf{y}} \otimes \xi_t \equiv \partial_{\mathbf{y}} \otimes \Phi_t(\mathbf{x}; \mathbf{y}) \quad (1.32)$$

2. $\xi'_t \in \mathbb{R}^{n^2}$ if

$$\xi'_t = \partial_{\mathbf{x}} \otimes \xi_t \equiv \partial_{\mathbf{x}} \otimes \Phi_t(\mathbf{x}; \mathbf{y}) \quad (1.33)$$

3. $\xi'_t \in \mathbb{R}^n$ if

$$\xi'_t = \dot{\xi}_t \equiv \dot{\Phi}_t(\mathbf{x}; \mathbf{y}) \quad (1.34)$$

For the first case we have that

$$\dot{\xi}'_t = \xi'_t \cdot \partial_{\xi_t} f(\xi_t; \mathbf{y}) + \partial_{\mathbf{y}} \otimes f(\xi_t, \mathbf{y})$$

as the dependence upon the parameters \mathbf{y} is explicit. In the remaining two cases, the dependence upon the parameters is implicit through the trajectory. We find that

$$\dot{\xi}'_t = \xi'_t \cdot \partial_{\xi_t} f(\xi_t) \quad (1.35)$$

For all the above cases it is expedient to define

$$A_t^{ij} = (\partial_{x_t^j} \otimes f^i)(\mathbf{x}_t) \quad (1.36)$$

In other words, A defines an endomorphism of \mathbb{R}^n : with slight abuse of language $A \in \text{End}(\mathbb{R}^n)$.

Proposition 1.4. *Let \mathfrak{t} denote the transposition operation, then*

$$A_t := (\partial_{\mathbf{x}_t} \otimes \mathbf{f})(\mathbf{x}_t) \quad (1.37)$$

and

$$F_t := (\partial_{\mathbf{x}} \otimes \phi_t)(\mathbf{x}) \quad (1.38)$$

then

$$\dot{F}_t = A_t F_t \quad (1.39)$$

Proof. By hypothesis

$$\frac{d\phi_t}{dt}(\mathbf{x}) = \lim_{dt \downarrow 0} \frac{\phi_{t+dt}(\mathbf{x}) - \phi_t(\mathbf{x})}{dt} = \lim_{dt \downarrow 0} \frac{\phi_{dt} \circ \phi_t(\mathbf{x}) - \phi_t(\mathbf{x})}{dt} = \mathbf{f} \circ \phi_t(\mathbf{x}) \quad (1.40)$$

hence it follows

$$\begin{aligned} \partial_{\mathbf{x}} \otimes \dot{\phi}_t(\mathbf{x}) &= \frac{d}{dt} \partial_{\mathbf{x}} \otimes \phi_t(\mathbf{x}) = \\ &(\partial_{\mathbf{x}} \otimes \phi_t)(\mathbf{x}) \cdot \partial_{\phi_t(\mathbf{x})} \mathbf{f} \circ \phi_t(\mathbf{x}) = (\partial_{\phi_t(\mathbf{x})} \otimes \mathbf{f})^{\mathfrak{t}} \circ \phi_t(\mathbf{x}) \cdot (\partial_{\mathbf{x}} \otimes \phi)_t(\mathbf{x}) = A_t F_t \end{aligned} \quad (1.41)$$

□

Of particular interest is the variation with respect to initial data

Definition 1.1. We will refer to the quantity

$$F_t = (\partial_{\mathbf{x}} \otimes \phi_t)(\mathbf{x}) \quad (1.42)$$

as the fundamental solution of the linearized dynamics.

By construction (1.42) satisfies the initial condition

$$F_0 = 1 \quad (1.43)$$

Namely from

$$\phi_0(\mathbf{x}) = \mathbf{x} \quad (1.44)$$

it follows that

$$F_0 \equiv \partial_{\mathbf{x}} \otimes \phi_0(\mathbf{x}) = 1 \quad (1.45)$$

Example 1.3. Let us consider again

$$\dot{x}_t = \alpha x_t^2 \quad \Rightarrow \quad x_t = \frac{1}{\frac{1}{x} - \alpha t} \quad (1.46)$$

We have then

$$\partial_x x_t = \frac{1}{(\frac{1}{x} - \alpha t)^2 x^2} := x'_t \quad (1.47)$$

which satisfies

$$\dot{x}'_t = \frac{2\alpha}{(\frac{1}{x} - \alpha t)^3 x^2} = 2\alpha x_t x'_t \quad (1.48)$$

as expected from direct differentiation of the ordinary differential equation in (1.46).

2 Flows: terminology

Definition 2.1 (Homeomorphism). Let $D \subset \mathbb{R}^n$ be an open domain. A map $\phi: D \mapsto D$ is called an homeomorphism if it is continuous and admits a continuous inverse.

Definition 2.2 (Diffeomorphism). Let $D \subset \mathbb{R}^n$ be an open domain. A map $\phi: D \mapsto D$ is called a diffeomorphism if it is continuously differentiable and admits a continuously differentiable inverse.

Definition 2.3. Similarly, for all $r \geq 1$, a C^r -diffeomorphism is an invertible map $\phi: D \mapsto D$ such that both ϕ and ϕ^{-1} admit continuous derivatives up to order r .

Definition 2.4 (flow). A one parameter family of diffeomorphisms Φ_t

$$\Phi: \mathbb{R} \times \mathbb{R}^n \mapsto \mathbb{R}^n \quad (2.1)$$

such that

$$\begin{aligned} \Phi_0 &= 1 \\ \Phi_t \Phi_s &= \Phi_{t+s} \end{aligned} \quad (2.2)$$

is called flow.

The definition of flow implies $\Phi_{-t} = \Phi_t^{-1}$ and therefore that Φ forms a one parameter t group of transformations. The combined use of the existence and uniqueness theorems for (1.1) allows us to conclude that

Proposition 2.1. The solution of (1.1) regarded as the map $\phi: \mathbb{R} \times D \mapsto D$ is a flow to which we will also refer as the fundamental solution of (1.1).

3 Evolution of volumes

Proposition 3.1. *Let $\mathbf{f} \in C^r(D)$ with $r \geq 1$ and let $M \subset D$. Let us denote by M_t the image of M through the flow ϕ of (1.1) with \mathbf{f} sufficiently regular. Let us define*

$$V_t^{[M]} = \text{Vol}(M_t) \quad (3.1)$$

then we have

$$\dot{V}_t^{[M]} = \int_{M_t} d^n x (\partial_{\mathbf{x}} \cdot \mathbf{f})(\mathbf{x}) \quad (3.2)$$

Proof. The flow being a diffeomorphism means that

1. ϕ is differentiable with the initial data

$$V_t^{[M]} = \int_{M_t} d^n x = \int_M d^n x |\det(\partial_{\mathbf{x}} \otimes \phi_t)(\mathbf{x})| \quad (3.3)$$

2. ϕ is invertible thus

$$\phi_0(\mathbf{x}) = \mathbf{x} \quad (3.4)$$

implies

$$\det(\partial_{\mathbf{x}} \otimes \phi_0)(\mathbf{x}) = \det \mathbf{1} = 1 \quad (3.5)$$

and

$$\det(\partial_{\mathbf{x}} \otimes \phi_t)(\mathbf{x}) > 0 \quad (3.6)$$

The strict inequality holds because otherwise

$$\mathbf{F}_t = (\partial_{\mathbf{x}} \otimes \phi_t)(\mathbf{x}) \quad (3.7)$$

solving the associated linearized dynamics

$$\dot{\mathbf{x}}' = \mathbf{x}' \cdot \partial_{\mathbf{x}_t} \mathbf{f}(\mathbf{x}_t) \equiv (\partial_{\mathbf{x}_t} \otimes \mathbf{f})(\mathbf{x}_t) \cdot \mathbf{x}' \quad (3.8)$$

would not be invertible. This is a contradiction because \mathbf{F}_t is the unique fundamental solution of the linearized dynamics.

We arrived at

$$V_t^{[M]} = \int_{M_t} d^n x = \int_M d^n x \det(\partial_{\mathbf{x}} \otimes \phi_t)(\mathbf{x}) \quad (3.9)$$

We observe that for any strictly positive matrix \mathbf{A} the equality

$$\ln \det \mathbf{A} = \text{tr} \ln \mathbf{A} \quad (3.10)$$

holds true. Then

$$\begin{aligned} \frac{d}{dt} \text{tr} \ln \partial_{\mathbf{x}} \otimes \phi_t &= \text{tr}(\partial_{\mathbf{x}} \otimes \phi_t)^{-1} \frac{d}{dt} (\partial_{\mathbf{x}} \otimes \phi_t) \\ &= \text{tr}(\partial_{\mathbf{x}} \otimes \phi_t)^{-1} (\partial_{\mathbf{x}} \otimes \phi_t) \partial_{\mathbf{x}_t} \otimes \mathbf{f}(\mathbf{x}_t) = \text{tr} \partial_{\mathbf{x}_t} \otimes \mathbf{f}(\mathbf{x}_t) = (\partial_{\mathbf{x}_t} \cdot \mathbf{f})(\mathbf{x}_t) \end{aligned} \quad (3.11)$$

□

References

- [1] N. Berglund. Geometrical theory of dynamical systems. ETH lecture notes, 2001, arXiv:math/0111177.
- [2] S. Wiggins. *Introduction to applied nonlinear dynamical systems and chaos*, volume 2 of *Texts in applied mathematics*. Springer, 2 edition, 2003.