

# Stochastic processes in physics and biology

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These notes are intended to cover the topics of the exercise sessions which are not included in the lecture notes. Please, don't hesitate to point out possible errors and misprints.

## 1 Stochastic calculus with Hermite polynomials

**Proposition 1.** *The transition probability of the Wiener process admits the expansion*

$$\frac{e^{-\frac{(x-y)^2}{2t}}}{\sqrt{2\pi t}} = \sum_{n=0}^{\infty} \frac{e^{-\frac{x^2}{2t}}}{\sqrt{2\pi t}} \left(\frac{y}{t}\right)^n h_n(x, t)$$

where the  $h_n$ 's are Hermite polynomials

$$h_n(x, t) = \frac{(-t)^n}{n!} e^{\frac{x^2}{2t}} \frac{d^n}{dx^n} e^{-\frac{x^2}{2t}} \quad (1)$$

*Proof.* Consider the Taylor expansion of the transition probability as a function of  $y$  at  $y = 0$ :

$$\frac{e^{-\frac{(x-y)^2}{2t}}}{\sqrt{2\pi t}} = \sum_{n=0}^{\infty} \frac{y^n}{n!} \frac{d^n}{dz^n} \frac{e^{-\frac{(x-z)^2}{2t}}}{\sqrt{2\pi t}} \Big|_{z=0} \quad (2)$$

Let us postulate for the  $n$ -th order of the Taylor expansion the form

$$\frac{y^n}{n!} \frac{d^n}{dz^n} \Big|_{z=0} \frac{e^{-\frac{(x-z)^2}{2t}}}{\sqrt{2\pi t}} := \frac{e^{-\frac{x^2}{2t}}}{\sqrt{2\pi t}} \left(\frac{y}{t}\right)^n h_n(x, t) \quad (3)$$

We will show that the  $h_n$ 's in 3 are exactly the Hermite polynomials defined in 1. So, let us start by manipulating the left hand side of 3:

$$\frac{e^{-\frac{x^2}{2t}}}{\sqrt{2\pi t}} \left(\frac{y}{t}\right)^n h_n(x, t) = \frac{y^n}{n!} \frac{d^n}{dz^n} \Big|_{z=0} \int_{\mathbb{R}} \frac{dp}{2\pi} e^{-ipz} \mathbb{E}(e^{ip\xi}) \quad (4)$$

where  $\xi$  is a random variable distributing according to  $N_t(z-x)$  (as a function of  $z$ ). In other words, we are just rewriting  $N_t(z-x)$  as the anti-Fourier transform of its characteristic function, which is by definition the Fourier transform of  $N_t(z-x)$ .

We know that

$$\mathbb{E}(e^{ip\xi}) = \exp \left\{ ip\mathbb{E}\xi - \frac{p^2}{2}\mathbb{E}\xi^2 \right\} = e^{ipx - p^2t/2} \quad (5)$$

hence we get

$$\begin{aligned}
\frac{e^{-\frac{x^2}{2t}}}{\sqrt{2\pi t}} \left(\frac{y}{t}\right)^n h_n(x, t) &= \frac{y^n}{n!} \frac{d^n}{dz^n} \Big|_{z=0} \int_{\mathbb{R}} \frac{dp}{2\pi} e^{-ipz} e^{ipx - p^2 t/2} = \\
&= \frac{y^n}{n!} \frac{d^n}{dz^n} \Big|_{z=0} \sum_{k=0}^{\infty} \int_{\mathbb{R}} \frac{dp}{2\pi} \frac{(-ipz)^k}{k!} e^{ipx - p^2 t/2} = \\
&= \frac{y^n}{n!} \sum_{k=0}^{\infty} \int_{\mathbb{R}} \frac{dp}{2\pi} \frac{(-ipz)^k}{k!} e^{ipx - p^2 t/2} n! \delta_{n,k} = \\
&= \frac{y^n}{n!} \int_{\mathbb{R}} \frac{dp}{2\pi} (-ipz)^n e^{ipx - p^2 t/2} = \\
&= \frac{(-y)^n}{n!} \frac{d^n}{dx^n} e^{-\frac{x^2}{2t}}
\end{aligned} \tag{6}$$

From the last equation we can conclude that indeed

$$h_n(x, t) = \frac{(-t)^n}{n!} e^{\frac{x^2}{2t}} \frac{d^n}{dx^n} e^{-\frac{x^2}{2t}} \tag{7}$$

□

Other two important properties of the Hermite polynomials are stated in the following proposition.

**Proposition 2.** (i) The  $h_n$ 's satisfy the following partial differential equation

$$(x \partial_x + 2t \partial_t) h_n(x, t) = n h_n(x, t)$$

(ii) The expected value of an Hermite polynomial having for argument a Wiener process starting at  $x$  is

$$\mathbb{E} h_n(w_t + x, t) = \frac{x^n}{n!} = h_n(x, 0) \tag{8}$$

*Proof.* (i) It is readily verified by performing the derivatives.

(ii) From the definition of expected value we have

$$\begin{aligned}
\mathbb{E} h_n(w_t + x, t) &:= \int_{\mathbb{R}} dy h_n(y, t) \frac{e^{-\frac{(y-x)^2}{2t}}}{\sqrt{2\pi t}} = \frac{(-t)^n}{\Gamma(n+1)} \int_{\mathbb{R}} dy \frac{e^{-\frac{x^2}{2t} + \frac{xy}{t}}}{\sqrt{2\pi t}} \frac{d^n}{dy^n} e^{-\frac{y^2}{2t}} \\
&= \frac{e^{-\frac{x^2}{2t}}}{\sqrt{2\pi t}} \left\{ \left( \frac{d^{n-1}}{dy^{n-1}} e^{-\frac{y^2}{2t}} \right) e^{\frac{xy}{t}} \Big|_{y=-\infty}^{y=+\infty} - \int dy \left( \frac{d^{n-1}}{dy^{n-1}} e^{-\frac{y^2}{2t}} \right) e^{\frac{xy}{t}} \left( \frac{x}{t} \right) \right\} \tag{9}
\end{aligned}$$

Iterating the integrations by parts one gets

$$\mathbb{E} h_n(w_t + x, t) = \frac{x^n}{n!} \tag{10}$$

In order to see that  $x^n/n! = h_n(x, 0) = \lim_{t \rightarrow 0} h_n(x, t)$ , we notice that the leading term in  $t^{-1}$  of  $\partial_x^n e^{-x^2/2t}$  is  $e^{-x^2/2t} (-x/t)^n$ , hence we get

$$\lim_{t \rightarrow 0} h_n(x, t) = \lim_{t \rightarrow 0} \frac{(-t)^n}{n!} e^{\frac{x^2}{2t}} e^{-\frac{x^2}{2t}} \left( -\frac{x}{t} \right)^n = \frac{x^n}{n!} \tag{11}$$

□

The last proposition shows that the expectation value of the Hermite polynomials is conserved. We can see this fact from a different perspective, namely as a consequence of the following proposition.

**Proposition 3.** *The differential of Hermite polynomials along a realization of the Wiener process is*

$$dh_n(w_t, t) = dw_t \partial_{w_t} h_n(w_t, t)$$

where the differential on the right hand side is an Ito differential.

*Proof.* By Ito lemma we have

$$dh_n(w_t, t) = dt \left( \partial_t + \frac{1}{2} \partial_{w_t}^2 \right) h_n(w_t, t) + dw_t \partial_{w_t} h_n(w_t, t)$$

In order to prove the claim we need to show that

$$\left( \partial_t + \frac{1}{2} \partial_{w_t}^2 \right) h_n(w_t, t) = 0$$

We can use the representation 1 and the property 8 to get for any  $t > 0$

$$\begin{aligned} 0 &= \left( \partial_t - \frac{1}{2} \partial_x^2 \right) \frac{e^{-\frac{(x-y)^2}{2t}}}{\sqrt{2\pi t}} = \\ &= \sum_{n=0}^{\infty} \frac{e^{-\frac{x^2}{2t}}}{\sqrt{2\pi t}} \left( \frac{y}{t} \right)^n \frac{1}{t} \left( -n + t \partial_t - \frac{t}{2} \partial_x^2 + x \partial_x \right) h_n = - \sum_{n=0}^{\infty} \frac{e^{-\frac{x^2}{2t}}}{\sqrt{2\pi t}} \left( \frac{y}{t} \right)^n \left( \partial_t + \frac{1}{2} \partial_x^2 \right) h_n \end{aligned}$$

which implies

$$\left( \partial_t + \frac{1}{2} \partial_x^2 \right) h_n = 0$$

as each of these multiply positive definite terms of different order in  $y$ . □

Now we can easily recover 8: consider

$$h_n(w_t + x, t) = h_n(x, 0) + \int_0^t dw_t \partial_{w_t} h_n(w_t, t)$$

From the property of the Ito integral we get

$$\mathbb{E} h_n(w_t + x, t) = \mathbb{E} h_n(x, 0) + \int_0^t \mathbb{E} [dw_t \partial_{w_t} h_n(w_t, t)] = h_n(x, 0) = \frac{x^n}{n!}$$

## 2 Recursion relation over the Wiener process

**Proposition 4.** *Stochastic integrals over Hermite polynomials satisfy the simple recursion relation*

$$\int_0^t dw_s h_n(w_s, s) = h_{n+1}(w_t, t)$$

*Proof.* Consider the exponential process:

$$\xi_t = e^{\lambda w_t - \frac{\lambda^2}{2} t} \tag{12}$$

Applying to it Ito lemma yields

$$d\xi_t = \lambda \xi_t dt + \lambda dw_t \tag{13}$$

or equivalently

$$e^{\lambda w_t - \frac{\lambda^2}{2} t} = 1 + \lambda \int_0^t dw_s e^{\lambda w_s - \frac{\lambda^2}{2} s}$$

Now consider the  $n$ -th derivative with respect to  $\lambda$  on both sides:

$$\left. \frac{d^n}{d\lambda^n} \right|_{\lambda=0} e^{\lambda w_t - \frac{\lambda^2 t}{2}} = \left. \frac{d^n}{d\lambda^n} \right|_{\lambda=0} \lambda \int_0^t dw_s e^{\lambda w_s - \frac{\lambda^2 s}{2}} \quad (14)$$

Contrasting the left-hand side of 14 with the expression for the Hermite polynomials in 3, we conclude

$$\left. \frac{d^n}{d\lambda^n} \right|_{\lambda=0} e^{\lambda w_t - \frac{\lambda^2 t}{2}} = t^n \left. \frac{d^n}{dz^n} \right|_{z=0} e^{\frac{z}{t} w_t - \frac{z^2}{2t}} = n! h_n(w_t, t) \quad (15)$$

where we set  $\lambda = z/t$ . The right hand side of 14 is

$$\begin{aligned} \left. \frac{d^n}{d\lambda^n} \right|_{\lambda=0} \lambda \int_0^t dw_s e^{\lambda w_s - \frac{\lambda^2 s}{2}} &= \left. \frac{d^{n-1}}{d\lambda^{n-1}} \frac{d}{d\lambda} \lambda \int_0^t dw_s e^{\lambda w_s - \frac{\lambda^2 s}{2}} \right|_{\lambda=0} = \\ &= \left. \frac{d^{n-1}}{d\lambda^{n-1}} \left[ \int_0^t dw_s e^{\lambda w_s - \frac{\lambda^2 s}{2}} + \lambda \frac{d}{d\lambda} \int_0^t dw_s e^{\lambda w_s - \frac{\lambda^2 s}{2}} \right] \right|_{\lambda=0} = \\ &= \left. \frac{d^{n-1}}{d\lambda^{n-1}} \int_0^t dw_s e^{\lambda w_s - \frac{\lambda^2 s}{2}} \right|_{\lambda=0} + \\ &\quad + \left. \frac{d^{n-2}}{d\lambda^{n-2}} \frac{d}{d\lambda} \left[ \lambda \frac{d}{d\lambda} \int_0^t dw_s e^{\lambda w_s - \frac{\lambda^2 s}{2}} \right] \right|_{\lambda=0} \end{aligned}$$

By iterating this procedure we get finally

$$\left. \frac{d^n}{d\lambda^n} \right|_{\lambda=0} \lambda \int_0^t dw_s e^{\lambda w_s - \frac{\lambda^2 s}{2}} = n \left. \frac{d^{n-1}}{d\lambda^{n-1}} \right|_{\lambda=0} \int_0^t dw_s e^{\lambda w_s - \frac{\lambda^2 s}{2}} \quad (16)$$

By looking at 15 we can rewrite the right hand side of the last equation as follows:

$$\begin{aligned} \left. \frac{d^n}{d\lambda^n} \right|_{\lambda=0} \lambda \int_0^t dw_s e^{\lambda w_s - \frac{\lambda^2 s}{2}} &= n \int_0^t dw_s \left. \frac{d^{n-1}}{d\lambda^{n-1}} \right|_{\lambda=0} e^{\lambda w_s - \frac{\lambda^2 s}{2}} = \\ &= n \int_0^t dw_s (n-1)! h_{n-1}(w_s, s) = \\ &= n! \int_0^t dw_s h_{n-1}(w_s, s) \end{aligned} \quad (17)$$

We have therefore proved that

$$h_n(w_t, t) = \int_0^t dw_s h_{n-1}(w_s, s)$$

□

An important consequence is the following. Since

$$h_0(w_t, t) = 1$$

we have that

$$\int_0^t dw_s = \int_0^t dw_s h_0(w_s, s) = h_1(w_t, t)$$

and

$$\int_0^t dw_{s_1} \int_0^{s_1} dw_{s_0} = h_2(w_t, t)$$

or in full generality

$$\int_0^t dw_{s_1} \prod_{i=1}^{n-1} \int_0^{s_{i-1}} dw_{s_{i-1}} = h_n(w_t, t)$$