

Semigroups & Delay eq., Solutions

Ex 1. Suppose $\lambda \in \rho(A)$ is such that

$$(\lambda - A)^{-1} \mathcal{Y} \subseteq \mathcal{X}.$$

Let $y \in \mathcal{Y}$. Then, since $R(\lambda - A) = \mathcal{X} \supseteq \mathcal{Y}$,
we can find $x \in D(A)$ such that

$$(\lambda - A)x = y \Leftrightarrow x = (\lambda - A)^{-1}y \in \mathcal{Y},$$

since $(\lambda - A)^{-1} \mathcal{Y} \subseteq \mathcal{Y}$. Thus two things follow:

$$1^\circ x \in D(A_1)$$

$$2^\circ (\lambda - A_1)^{-1}y = (\lambda - A)^{-1}y$$

$\therefore \lambda \in \rho(A_1)$ and $R(\lambda, A_1) = R(\lambda, A)$, \square

Ex 2. By the very definition of the generator of a semigroup

$$x^0 \in D(A^0) \Leftrightarrow \exists \lim_{h \rightarrow 0} \frac{T^0(h)x^0 - x^0}{h} \text{ in } \mathcal{X}^0$$

But since $T^0 = T^*|_{\mathcal{X}^0}$ this is equivalent to

$$\exists \lim_{h \rightarrow 0} \frac{T^*(h)x^0 - x^0}{h} =: Ax^0 \text{ and } Ax^0 \in \mathcal{X}^0$$

$\Leftrightarrow x^0$ as an element of \mathcal{X}^* belongs to $D(A^*)$ and $A^*x^0 \in \mathcal{X}^0$.

$\Leftrightarrow x^0 \in D(A^*|_{\mathcal{X}^0})$. Done.

$\rho(A^*) \subseteq \rho(A^0)$: Follows from Ex. 1 with the observation $(\lambda - A^*)^{-1}x^0 \in \mathcal{X}^0$ for $\lambda \in \rho(A^*)$

$\rho(A^0) \subseteq \rho(A^*)$: Let $\lambda \in \rho(A^0)$ and fix $y \in \mathcal{X}$.

For small enough $\varepsilon > 0$ $\exists (\lambda - \varepsilon A^*)^{-1}$ too, and $(\lambda - \varepsilon A^*)^{-1}y \xrightarrow{\|\cdot\|} y$ as $\varepsilon \rightarrow 0$.

Now, since $(\lambda - \varepsilon A^*)^{-1}y \in D(A^*) \subseteq \mathcal{X}^0$ we find $x_\varepsilon \in D(A^0)$ such that $x_\varepsilon = (\lambda - A^0)^{-1}(\lambda - \varepsilon A^*)^{-1}y$. Then, by boundedness of $(\lambda - A^0)^{-1}x_\varepsilon$ is Cauchy so $\exists \lim_{\varepsilon \rightarrow 0} x_\varepsilon =: x$

... Ex 2: Now we have that

$$(I - A^*)x_\varepsilon = (I - A^0)x_\varepsilon \text{ converges (to } y)$$

$$\text{and } x_\varepsilon = (I - A^0)^{-1}(I - \varepsilon A^*)^{-1}y \text{ converges to } x$$

as $\varepsilon \rightarrow 0$.

Therefore, by closedness of A^* , we have $x \in \mathcal{D}(A^*)$ and $(I - A^*)x = y$

$\therefore 1 \in \rho(A^*) \quad \square$

Ex 3: I use slightly different notation:

vector ψ is denoted by $|\psi\rangle$ when it's considered as an element of a "original space", and by $\langle\psi|$ when it's acting as a functional on e.g. \mathcal{X}^* .

With that convention B takes the form

$$B = \sum_j |r_j^{0*}\rangle \langle r_j^*| : \mathcal{X} \rightarrow \mathcal{X}^{0*}$$

It's also very easy to check that $B^* := B^*|_{\mathcal{X}^0}$ is

$$B^* = \sum_j |r_j^*\rangle \langle r_j^{0*}| : \mathcal{X}^0 \rightarrow \mathcal{X}^*$$

For warm up let us take the case " $M(A)c = c$ ":

$$A|\psi\rangle = (A_c^{0*} + B)|\psi\rangle = \lambda|\psi\rangle$$

for $\lambda \in \rho(A_c^{0*})$ if and only if

$$(A - A_c^{0*})^{-1} B|\psi\rangle = |\psi\rangle. \quad \textcircled{*}$$

This holds iff $|\psi\rangle \in \mathcal{R}(A - A_c^{0*})^{-1} B = \underline{\text{span}} \{|r_j(A)\rangle\}_j$,

which implies the representation

$$|\psi\rangle = \sum_j |r_j(A)\rangle \langle r_j(A)|\psi\rangle,$$

and $\textcircled{*}$ holds in the new representation...

... Ex 3: ...

$$(A - A^{0*})^{-1} \sum_i |r_i^{0*}\rangle \langle r_i^{0*}| \sum_j |r_j^0\rangle \langle r_j^0| \psi \rangle =: c_j$$

$$= \sum_i \underbrace{(A - A^{0*})^{-1} |r_i^{0*}\rangle}_{=: |r_i^0\rangle} \sum_j \underbrace{\langle r_i^{0*}| r_j^0\rangle}_{=: M(A)_{ij}} c_j$$

$$= \sum_i |r_i^0\rangle (M(A) \bar{c})_i$$

"right side of \oplus " = $|\psi\rangle = \sum_j |r_j^0\rangle c_j$.

Since $|r_j^0\rangle$ are independent, this is true iff

$$M(A) \bar{c} = \bar{c}. \text{ Done}$$

Now the dual case $A^* |\psi^0\rangle = \lambda |\psi^0\rangle$ is similar. As before it's easily checked

$$A^* |\psi^0\rangle = (A_0^0 + B^*) |\psi^0\rangle = \lambda |\psi^0\rangle$$

$$\text{iff } (\lambda - A_0^0)^{-1} B^* |\psi^0\rangle = |\psi^0\rangle. \quad (*)$$

Again this means $|\psi^0\rangle \in \mathcal{R}((\lambda - A_0^0)^{-1} B^*) = \text{span} \{ |r_j^0(A)\rangle \}$
 so that $|\psi^0\rangle$ is a linear comb.

$$|\psi^0\rangle = \sum_i |r_i^0(A)\rangle \langle r_i^0(A) | \psi^0 \rangle$$

and then (*) becomes...

Ex 3: ...

$$(\lambda - A^0)^{-1} \sum_i |r_i^*\rangle \langle r_i^{0*}| \sum_j |r_j^0\rangle \langle r_j^0| \psi^0\rangle$$

$=: d_j$

$$= \sum_i \frac{(\lambda - A^0)^{-1} |r_i^*\rangle}{\langle r_i^0|} \sum_j \frac{\langle r_i^{0*}| r_j^0\rangle}{M(\lambda)_{ji}} d_j$$

$$= \sum_i |r_i^0\rangle (\bar{d}^T M(\lambda))^{-1}$$

"right side of (+)" = $|\psi^0\rangle = \sum |r_i^0\rangle d_i$

which holds if $\bar{d}^T M(\lambda) = \bar{d}^T \square$

Remarks: a) For me \bar{d} is not a row vector.

b) I defined $M(\lambda)_{ij}$ as $\langle r_i^* | r_j^0\rangle$ which is the reversed order of i and j as compared to the lectures.

c) $\langle r_i^{0*} | r_j^0\rangle$ is $M(\lambda)_{ji}$ because, formally,

$$\begin{aligned} \langle r_i^{0*} | (\lambda - A^0)^{-1} r_j^* \rangle &= \langle (\lambda - A^{0*})^{-1} r_i^{0*} | r_j^* \rangle \\ &= \langle r_i^0 | r_j^* \rangle \\ &= \langle r_j^* | r_i^0 \rangle \\ &= M(\lambda)_{ji} \quad (\text{in my book}), \end{aligned}$$

Ex. 4 First define

$$g_{ij}(t) := \left\langle \gamma^{-1} \int_0^t T_0^{0*}(t-s) r_j^{0*} ds, r_i^* \right\rangle, \quad (t \geq 0).$$

Then g_{ij} is locally Lipschitz, for

$$\begin{aligned} |g_{ij}(t) - g_{ij}(\bar{t})| &\leq \int_0^t \underbrace{\|T_0^{0*}(t-s) r_j^{0*}\|}_{\leq M e^{\omega(t-s)} \|r_j^{0*}\|} ds \|r_i^*\| \\ &\leq M e^{\omega t} - e^{\omega \bar{t}} \|r_j^{0*}\| \|r_i^*\| \end{aligned}$$

(pick $\omega \neq 0$)
$$\leq M \frac{e^{\omega \bar{t}} - e^{\omega t}}{\omega} \|r_j^{0*}\| \|r_i^*\|$$

Thus g_{ij} is abs. cont. so $\exists k_{ij} \in L^1_{loc}$ s.t.

$$g_{ij}(t) = \int_0^t k_{ij}(s) ds,$$

and in fact k_{ij} must be L^∞_{loc} or otherwise g_{ij} is not Lipschitz.

Now let $\eta \in L^1_{loc}$. We show that

$$\left\langle \gamma^{-1} \int_0^t T_0^{0*}(t-\tau) r_j^{0*} \eta(\tau) d\tau, r_i^* \right\rangle = \int_0^t k_{ij}(t-\tau) \eta(\tau) d\tau$$

by showing the integrals $\int_0^t \dots ds$ of both sides coincide for all $t \geq 0$:

Right side:

$$\begin{aligned} \int_0^t ds \int_0^s k_{ij}(s-\tau) \eta(\tau) d\tau &= \int_0^t d\tau \left(\int_\tau^t k_{ij}(s-\tau) ds \right) \eta(\tau) \\ &= \int_0^t d\tau g_{ij}(t-\tau) \eta(\tau). \end{aligned}$$

... Ex 4: Left side:

$$\begin{aligned} & \int_0^+ ds \left\langle \gamma^{-1} \int_0^+ T_0^{0*}(s-\tau) r_j^{0*} \eta(\tau) d\tau, r_i^* \right\rangle \\ &= \left\langle \gamma^{-1} \int_0^+ d\tau \eta(\tau) \int_{\tau}^+ ds T_0^{0*}(s-\tau) r_j^{0*}, r_i^* \right\rangle \\ &= \int_0^+ d\tau \eta(\tau) g_{ij}(t-\tau) \quad \text{Done.} \end{aligned}$$

Lastly, take the Laplace \mathcal{T} :

$$\begin{aligned} & \int_0^{\infty} ds e^{-\lambda s} \left\langle \gamma^{-1} \int_0^+ ds T_0^{0*}(s-\tau) r_j^{0*} \eta(\tau), r_i^* \right\rangle \\ &= \left\langle \gamma^{-1} \int_0^{\infty} ds e^{-\lambda s} \eta(s) \int_{\tau}^{\infty} ds e^{-\lambda(s-\tau)} T_0^{0*}(s-\tau) r_j^{0*}, r_i^* \right\rangle \\ &= \underbrace{\int_0^{\infty} ds e^{-\lambda s} \eta(s)}_{\hat{\eta}(\lambda)} \left\langle \gamma^{-1} (\lambda - A^{0*})^{-1} r_j^{0*}, r_i^* \right\rangle \end{aligned}$$

On the other hand, the Laplace \mathcal{T} of $K_{ij}^* \eta$ is $\hat{K}_{ij}^*(\lambda) \eta(\lambda)$, so, since $\eta \in L^1_{loc}$ is arbitrary,

$$\hat{K}_{ij}^*(\lambda) = \left\langle \underbrace{\gamma^{-1} (\lambda - A^{0*})^{-1} r_j^{0*}}_{r_j(\lambda)}, r_i^* \right\rangle = M(\lambda)_{ij}$$

(using the notation convention of the lectures) \square

Ex 5: Given φ , the solution to

$$u(t) = T_0(t)\varphi + \int_0^t T_0^{0*}(t-s)(B u(s) + f(s)) ds \quad (1)$$

is unique. For, if u_1, u_2 are two solutions then

$$\begin{aligned} \|u_1(t) - u_2(t)\| &\leq \int_0^t \|T_0^{0*}(t-s)(B(u_1(s) - u_2(s)))\| ds \\ &\leq \int_0^t M e^{\alpha(t-s)} \|B\| \|u_1(s) - u_2(s)\| \end{aligned}$$

so $u_1(t) \equiv u_2(t)$ by Grönwall's lemma.

Therefore, all we have to do is show that solution u to

$$u(t) = T(t)\varphi + \int_0^t T^{0*}(t-s)f(s) ds \quad (2)$$

satisfies (1) too.

Plugging (2) to the left of (1) gives

$$\underline{T_0(t)\varphi} + \int_0^t \underline{T_0^{0*}(t-s)B T(s)\varphi} + \int_0^t T_0^{0*}(t-s)f(s) ds$$

(Recall $T(t)\varphi = T_0\varphi + \int_0^t T_0^{0*}(t-s)B T(s)\varphi$).

Plugging (2) to right of (2) gives

$$\underline{T_0(t)\varphi} + \int_0^t \underline{T_0^{0*}(t-s)B \left\{ T(s)\varphi + \int_0^s T(s-\tau)f(\tau) d\tau \right\}} + \int_0^t T_0^{0*}(t-s)f(s)$$

Making the underlined comparisons and cancellations it comes down to showing...

... Ex 5 ... that following holds

$$\int_0^t T^{0*}(t-s) f(s) ds = \int_0^t T_0^{0*}(t-s) B \int_0^s T_0^{0*}(s-\tau) f(\tau) d\tau + \int_0^t T_0^{0*}(t-s) f(s) ds$$

$$\int_0^t T_0^{0*}(t-s) B \int_0^{t-s} T_0^{0*}(t-s-\tau) f(\tau) d\tau$$

Since for each $x \in \mathcal{X}$ and $t \geq 0$ $T(t)$ obeys

$$T(t)x = T_0(t)x + \int_0^t T_0^{0*}(t-s) B T(s)x ds,$$

so it looks "formally right". The trouble, of course, is that $f(s) \in \mathcal{X}^{0*}$.

Remedy is to show instead that the integrals $\int_0^t dt$ over both sides match for all $t \geq 0$, and then use integrated semigroups. For example left side becomes

$$\int_0^t \int_0^s T_0^{0*}(s-\tau) f(\tau) d\tau ds = \int_0^t W(t-s) f(s) ds,$$

where $W(t) := \int_0^t T_0^{0*}(s) ds$.

This approach works because W have a "smoothing effect" that $W: \mathcal{X}^{0*} \rightarrow \mathcal{X}$. The details are somewhat tedious I'm afraid. The remaining part of the proof/calculation can be found on page 75, prop. 2.5, in Perturbation Theory for dual Semigroups III.