

We get, since  $S(t)$  contracts in  $L^2$ :

$$0 \leq \hat{S}(t, k) \leq \lim_{\Lambda} (f_{\Lambda}, f_{\Lambda}) = \hat{S}(0, k)$$

Spectral theorem gives

$$(f_{\Lambda, k}, S(t) f_{\Lambda, k}) = \int_{\mathbb{R}^d} d\nu(\lambda) e^{-\lambda t}$$

where

$$(f_{\Lambda, k}, L f_{\Lambda, k}) = \int \lambda d\nu_{\Lambda}(\lambda)$$

Thus

$$(f_{\Lambda, k}, S(t) f_{\Lambda, k}) = S(0, k) \frac{\int d\nu_{\Lambda}(\lambda) e^{-\lambda t}}{\int d\nu_{\Lambda}(\lambda)}$$

Jensen's inequality  $\Rightarrow$

$$e^{-\lambda t} > \exp(-t) \frac{\int \lambda d\nu_{\Lambda}}{\int d\nu_{\Lambda}} = \exp(-t) \frac{(f_{\Lambda, k}, L f_{\Lambda, k})}{(f_{\Lambda, k}, f_{\Lambda, k})}$$

$$\text{Now } \lim_{\Lambda} (f_{\Lambda, k}, L f_{\Lambda, k}) = \sum_x e^{ikx} \underbrace{E^{\Lambda} (f_0, L f_x)}_{= g(x)}$$

$$\text{where } f_x(y) = \eta(x-y), \quad g(x) = (f_0, L f_x) =$$

$$= (L f_0, f_x) \quad \text{by } \underline{\text{reversibility}}.$$

$$\text{Also } g(x) = (f_{-x}, L f_0) \quad \text{by } \underline{\text{translation invariance}}$$

$$\text{Hence } g(x) = \overline{g(-x)}.$$

$$\text{Since } L S = 0, \quad g(x) = (\eta(0), L \eta)(x)$$

$$\text{We have } \sum_x g(x) = 0, \quad \text{by conservation of}$$

particle number. Indeed,

$$\begin{aligned}
 L\eta(x) &= \frac{1}{2} \sum_{u,v} c(u,v,\eta) \underbrace{(\eta^{uv}(x) - \eta(v))}_{=0 \text{ unless } x=u \text{ or } x=v} \\
 &= \frac{1}{2} \sum_v c(x,v,\eta) (\eta(v) - \eta(x)) \\
 &\quad + \frac{1}{2} \sum_u c(u,x,\eta) (\eta(u) - \eta(x))
 \end{aligned}$$

(finite sum,  $|x-v| \in \mathbb{R}$ ,  $|v-u| \in \mathbb{R}$ )

Thus ( $g(x)$  decays exponentially or absolutely summable)

$$\begin{aligned}
 \sum_x g(x) &= \sum_{x,u} (\eta(x)) (c(x,u,\eta) + c(u,x,\eta)) (\eta(u) - \eta(x)) \\
 &= 0
 \end{aligned}$$

We get, after  $\Lambda \rightarrow \infty$ :

$$\hat{S}(\ell, k) \geq \hat{S}(\ell, k) \geq \hat{S}(0, k) \exp\left[-\ell \frac{A(k)}{\hat{S}(0, k)}\right]$$

$$A(k) = \frac{1}{2} \sum_x (1 - \cos kx) g(x)$$

$$g(v) = (\eta(0), L\eta(v)) = \sum_y (\eta(0), c(x,y,\eta) (\eta(y) - \eta(x)))$$

$$\text{We have } A(k) = \frac{1}{4} \sum_{i,j} a_{ij} k_i k_j + o(k^2) \text{ as } k \rightarrow 0$$

$$a_{ij} = \sum_x x_i v_j g(x)$$

Also

$$\int_{(-\pi, \pi)^d} \frac{dk}{(2\pi)^d} \hat{S}(\ell, k) = E^\eta(\eta_\ell(0) \eta_0(0)) - S^2$$

$$\hat{S}(0, 0) = \sum_x E^\eta(\eta(x) \eta(0)) - S^2 = X$$

(Since  $E^\eta(\eta(0)\eta(0)) = E^\eta(\eta(0)) = S$ ,  $\hat{S}(0,0) \geq S - S^2 > 0$ )

$X$  is "static compressibility",

Hence

$$\Gamma^M(\gamma_x(t), \gamma_y(t)) - \xi^2 \geq \int \frac{dk}{(2\pi)^d} \hat{S}(t, k) e^{-\xi \frac{A(k)}{\hat{S}(t, k)}}$$

Exercise compute  $a_{ij}$  for symmetric

$$a_{ij} = \begin{cases} 1 & |x-y|=1 \\ 0 & \text{otherwise} \end{cases}$$

and show  $a > 0$  (as a matrix)

$$\text{Typically, } a > 0 \text{ so } \frac{A(k)}{\hat{S}(t, k)} \sim k^2 \text{ as } k \rightarrow 0$$

get

$$\Gamma^M(\gamma_x(t), \gamma_y(t)) - \xi^2 \geq \text{const} \frac{1}{t^{d/2}}$$

(i.e. correlation in time decaying only as a power law (not exponentially)).

The lower bound

$$\hat{S}(t, k) \geq c e^{-\lambda_0 k^2}$$

is consistent with

$$S(t, x) \leq c t^{-d/2} e^{-x^2/4Dt}$$

which is diffusive behaviour.

A challenging conjecture is:

$$\lim_{L \rightarrow \infty} L^d S(-t, Lx) = \chi \frac{1}{\sqrt{\det D}} \frac{1}{(4\pi t)^{d/2}} e^{-\frac{(x, D^{-1} x)}{4t}}$$

$D$  positive matrix

Green-Kubo formula

Define, whenever exists, the diffusion matrix

$$D_{\alpha\beta} = \lim_{t \rightarrow \infty} \frac{1}{t} \frac{1}{X} \sum_{X \in \mathbb{Z}^d} X_\alpha X_\beta S(t, X)$$

where

$$X = \sum_x S(0, x)$$

We will show the limit exists and we will

relate  $D_{\alpha\beta}$  to a current correlation function.

$$\text{We have } Lf(\eta) = \frac{1}{2} \sum_{x,y} c(x,y,\eta) (f(\eta^{xy}) - f(\eta))$$

$\eta_e(x)$  changes in time due to jumps of particles to  $x$  or from  $x$ .

$$\text{Rate of jumps } x \rightarrow y = c(x,y,\eta) \eta(x) (1 - \eta(y))$$

$$\text{---} \quad y \rightarrow x = c(y,x,\eta) \eta(y) (1 - \eta(x))$$

$$\text{Flux of particles from } x \rightarrow y := \underline{j(x,y,\eta)} =$$

$$c(x,y,\eta) [\eta(x)(1 - \eta(y)) - \eta(y)(1 - \eta(x))] = \underline{c(x,y,\eta) (\eta(x) - \eta(y))}$$

Let  $\eta_t \in \mathcal{D}[0, \infty]$  be a trajectory of our process

$t \mapsto \eta_t(x) \in \{0, 1\}$  consists of jumps

$$\text{Let } J_{xy}([s, t]) = \# \text{ jumps } x \rightarrow y - \# \text{ jumps } y \rightarrow x \text{ during time } [s, t]$$

We can define it as a process coupled to

$\eta$  as follows.

Define Markov process  $(\eta_t, N_t)$ ;  $\eta_t \in \{0, 1\}^{\mathbb{Z}^d}$ ,

$N_t \in \{0, 1\}$  by generator

$$L f(\eta, N) = \frac{1}{2} \sum_{(x', y') \neq (x, y)} c(x', y', \eta) (f(\eta^{x', y'}, N) - f(\eta, N))$$

$$+ c(x, y, \eta) \eta(x)(1-\eta(y)) [f(\eta^{x, y}, N+1) - f(\eta, N)]$$

$$+ c(x, y, \eta) (1-\eta(x)) \eta(y) [f(\eta^{x, y}, N-1) - f(\eta, N)]$$

We set  $J_{xy}[S, A] = N_t - N_s$ .

Recall that for  $f \in \mathcal{D}(L)$

$$f(\eta_t, N_t) - \int_0^t (L f)(\eta_s, N_s) ds$$

is a martingale. Take  $f(\eta, N) \equiv N$

and  $N_0 = 0$ . Then  $L f = c(x, y, \eta) [\eta(x)(1-\eta(y))$

$$- \eta(y)(1-\eta(x))] = j(x, y, \eta) \text{ or}$$

$$M_{xy}(t) = J_{xy}[0, t] - \int_0^t j(x, y, \eta_s) dt'$$

is a martingale.

Prop Let  $\mu$  be stationary measure. Then

$$E^\mu M_{xy}(t) M_{x'y'}(t) = t (\delta_{xx'} \delta_{yy'} - \delta_{xy'} \delta_{yx'}) \\ - E^\mu (c(x, y, \eta))$$

Pf We need a Lemma

Lemma Let  $X_t \in Y$  be a Markov process with generator  $L$ . Let  $f \in \mathcal{D}(L)$  be s.d.

$f^2 \in \mathcal{D}(L)$ . Then

$$M(t) := f(X_t) - \int_0^t Lf(X_s) ds$$

$$R(t) := M(t)^2 - \int_0^t ((Lf)^2(X_s) - 2f(X_s)Lf(X_s)) ds$$

are Martingales.

Assuming the Lemma, let us define

$$N(t) = \int_{x,y} (cov(X)) \quad \text{and} \quad N'(t) = \int_{x,y'} (cov(X'))$$

as above (add to the generator above also  $x', y'$  variables i.e.  $S(\mathcal{Z}, N, N')$ ).

Apply lemma to  $F_1, F_2, F_1 + F_2$ . Get

$$M_1(t)M_2(t) - \int_0^t (L(F_1F_2) - F_1LF_2 - F_2LF_1) ds$$

is Martingale. So

$$0 = E(M(t)M'(t)) - \int_0^t E(L(NN') - NN' - N'LN)$$

So by stationarity, if  $\{x, y\} \neq \{x', y'\}$ ,

$$E(M'(t)M'(t)) = \mathbb{1} E^N(L(NN') - NN' - N'LN)$$

We compute  $L(NN') = (LN)N' + N(LN') \Rightarrow \text{get } = 0$ .

$$E M^2 = \mathbb{1} E^N(LN^2 - 2NLN) = \mathbb{1} E^{(x,y,z)} [\eta(x)(1-\eta(y))$$

$$((W(t))^2 - N^2) + 2N(N'-N) + \eta(y)(1-\eta(x))((W(t))^2 - N^2) + 2N(N'-N)]$$

$$= \lambda E [c(x, y, \eta) (\eta(x)(1-\eta(y)) + \eta(y)(1-\eta(x)))]$$

Let us define  $c(x, y, \eta) \equiv 0$  if  $\eta(x) = \eta(y)$

(this is no loss!). Since  $\eta(x) = 1$  if  $\eta(x) \neq \eta(y)$

$$\text{get } = \lambda E [c(x, y, \eta)] \quad \square$$

We can now compute the current-current correlation function:

Exercise Show  $E(\sum_{x,y} \delta_{x,y}(0, t) \int_0^t \int_0^t j(x, y, \eta_s) ds) = 0$   
using reversibility.

Then

$$E^M(\sum_{x,y} \delta_{x,y}(0, t) \sum_{x,y} \delta_{x,y}(0, t)) = E(\sum_{x,y} \delta_{x,y} \delta_{x,y} - \delta_{x,y} \delta_{y,x}) \\ = E^M(c(0,0,\eta)) - \int_0^t ds \int_0^s ds' E^M(j(x, y, \eta_s) + j(y, x, \eta_{s'}))$$

Define

$$j_\alpha(\eta) = \frac{1}{2} \sum_{y \neq x} v_{xy} j(x, y, \eta)$$

Then we have

Prop. The limit

$$D_{\alpha\beta} = \lim_{A \rightarrow \infty} \frac{1}{A} \frac{1}{2X} \sum_x X_\alpha X_\beta S(x, x)$$

exists and is given by

$$D_{\alpha\beta} = \frac{1}{2X} \left( \frac{1}{2} \delta_{\alpha\beta} \sum_x X_\alpha^2 E^M(c(0, x, \eta)) \right. \\ \left. - 2 \int_0^\infty dt \sum_x E^M(j_\alpha S(t) \tau_x j_\beta) \right)$$

when the sums and integrals converge

Pf. We have from conservation of particles

$$\eta_\alpha(x) - \eta_\alpha(0) = - \sum_y J_{xy} [0, A]$$

so

$$\sum_x x_\alpha x_\beta (S(A, x) - S(0, x)) = -\frac{1}{2} \sum_{\alpha, \beta} E(\eta_\alpha(x) - \eta_\alpha(0))(\eta_\beta(0) - \eta_\beta(x))$$

$$= -\frac{1}{2} \sum_x \sum_{y, y'} x_\alpha x_\beta E(J_{xy} [0, A] J_{y'y'} [0, A])$$

$$= \frac{1}{2} \sum_x x_\alpha x_\beta E(C(x, 0, \eta)) - \frac{1}{2} \sum_{x, y, y'} \int_0^x ds \int_0^{y'} ds'$$

$$x_\alpha x_\beta E(j(x, y, \eta(s)) j(0, y', \eta(s')))$$

Recalling  $j(x, y, \eta) = C(x, y, \eta) (\eta(x) - \eta(y)) = -j(y, x, \eta)$

some algebra gives (exercise)

$$x = -2 \int_0^x ds \int_0^{y'} ds' \sum_x E(j_\alpha(\eta(s)) j_\beta(\eta(s')))$$

Now

$$E(j_\alpha(\eta(s)) j_\beta(\eta(s'))) = (j_\alpha, S(|s-s'|) \tau_x j_\beta) \\ \equiv G_{\alpha\beta}(|s-s'|, x) = \text{symmetric matrix.}$$

Let  $v \in \mathbb{R}^d$ ,  $G_{\alpha\beta} = \sum_{\alpha, \beta} v_\alpha v_\beta G_{\alpha\beta}$ ,  $j = \sum j_\alpha v_\alpha$

$$G_v(T, x) = (j, S(T) \tau_x j)$$

This decays exponentially in  $x$  for fixed  $T$  (proof

as for  $S(t, x)$ , use  $Ej = 0$ )



Thus we may write

$$\sum_x G_T(T, x) = \lim_{\Lambda \uparrow \mathbb{Z}^d} (j_\Lambda, S(T) j_\Lambda)$$

$$j_\Lambda \equiv \frac{1}{\sqrt{|\Lambda|}} \sum_{x \in \Lambda} j(x)$$

We have obtained

$$\frac{1}{\lambda} \sum_x (v \cdot x)^2 S(\lambda, x) = \frac{1}{\lambda} \sum_x (v \cdot x)^2 S(\lambda, x)$$

$$+ \frac{1}{2} \sum_x (v \cdot x)^2 E(C(v, x, \lambda)) - \frac{1}{\lambda} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} \partial_i^2 (j_\Lambda, S(\lambda, v \cdot x) j_\Lambda) \lim_{\Lambda \uparrow \infty}$$

Lemma LHS  $> 0$ .

PF

$$\text{LHS} = -\frac{1}{\lambda} \sum_{\alpha, \beta} v_\alpha v_\beta \frac{\partial^2 \hat{S}(\lambda, k)}{\partial k_\alpha \partial k_\beta} \Big|_{k=0}$$

But  $\hat{S}(\lambda, k) = \lim_{\Lambda \uparrow \mathbb{Z}^d} (f_{\Lambda, k}, S(\lambda) f_{\Lambda, k}) - g^2$

$$f_{\Lambda, k} = \frac{1}{\sqrt{|\Lambda|}} \sum_{y \in \Lambda} e^{-iky} \varphi(y) \quad (\text{for } k \neq 0)$$

$$|f_{\Lambda, k}| \leq f_{\Lambda, 0} \Rightarrow \hat{S}(\lambda, k) \leq \hat{S}(\lambda, 0)$$

Also  $\partial_k \hat{S}(\lambda, k) \Big|_{k=0} = 0$  so  $k=0$  is a

maximum  $\Rightarrow$  claim  $\square$