

# 1 Introduction

Beside [3] alternative reference for the Ito process are [1], [2] and Varadhan's lecture notes [4].

# 2 Martingales and Doob-Meyer decomposition

This section follows [2]. In the last lecture we identified a  $d$ -dimensional Ito process as the stochastic differential

$$d\xi_t = \mathbf{b}_t dt + \boldsymbol{\sigma}_t[d\mathbf{w}_t]$$

$$\sigma_t^i[d\mathbf{w}_t] := \sum_{j=1}^n \sigma_t^{ij} dw_t^j \quad \begin{array}{l} i = 1, \dots, d \\ j = 1, \dots, n \end{array}$$

for  $\mathbf{b}_t, \boldsymbol{\sigma}$  *non-anticipative* with respect to the Wiener-process. Alternatively we can write

$$\xi_t = \xi_o + \int_0^t ds \mathbf{b}_s + \int_0^t \boldsymbol{\sigma}_s[d\mathbf{w}_s]$$

The right hand side comprises three terms.

- An initial value  $\xi_o$ , eventually deterministic, for the process.
- A *bounded variation* component

$$\mathbf{B}_t := \int_0^t ds \mathbf{b}_s$$

This means that for any vector component  $i$  of  $\mathbf{B}$  we have

$$\lim_{|p| \downarrow 0} \sum_k (t_{k+1} - t_k) |B_{t_{k+1}}^i - B_{t_k}^i| = \int_0^t dt |B^i| < \infty$$

which implies the vanishing of the quadratic variation.

- A *martingale* component

$$\mathbf{M}_t = \int_0^t \boldsymbol{\sigma}_t[d\mathbf{w}_t]$$

From the properties of the Ito-integral we have in fact:

- Conservation of the average: for any  $t$

$$\langle \int_0^t \boldsymbol{\sigma}_s[d\mathbf{w}_s] \rangle = 0$$

since  $\boldsymbol{\sigma}$  is non-anticipating.

- finite quadratic variation

$$\langle \int_0^t \sigma_s^i[d\mathbf{w}_s] \int_0^t \sigma_{s'}^j[d\mathbf{w}_{s'}] \rangle = \int_0^t ds \sigma_s^{ik} \sigma_s^{jk} < 0$$

As usual Einstein convention is implied: repeated indices stand for index contraction.

The representation of an Ito process in the form

$$\xi_t = \xi_o + B_t + M_t$$

is often referred to as the *Doob-Meyer representation*

**Example 2.1** (*Exponential martingale*). Let us consider the process

$$\xi_t = e^{\lambda w_t - \frac{\lambda^2 t}{2}} \xi_o \tag{2.1}$$

by Ito lemma we have

$$d\xi_t = \lambda dw_t e^{\lambda w_t - \frac{\lambda^2 t}{2}} \xi_o = \lambda \xi_t dw_t$$

If we recast the Ito differential into Doob-Meyer form we find

$$\xi_t = \xi_o + \lambda \int_0^t dw_s \xi_s$$

The exponential process does not have bounded variation component. It is therefore a *martingale*

$$\langle \xi_t \rangle = \langle \xi_o \rangle$$

if

$$\langle |\xi_o| \rangle < \infty$$

### 3 Stochastic calculus with Hermite polynomials

This section expands example D.3 of chapter 4 of [3].

**Proposition 3.1** (*Expansion of the transition probability of the Wiener process*). We have

$$\frac{e^{-\frac{(x-y)^2}{2t}}}{\sqrt{2\pi t}} = \sum_{n=0}^{\infty} \frac{e^{-\frac{x^2}{2t}}}{\sqrt{2\pi t}} \left(\frac{y}{t}\right)^n h_n(x, t)$$

where

$$h_n(x, t) = \frac{(-t)^n}{\Gamma(n+1)} e^{\frac{x^2}{2t}} \frac{d^n}{dx^n} e^{-\frac{x^2}{2t}} \tag{3.1}$$

*Proof.* The  $n$ -th order of the Taylor expansion can be couched into the form

$$\frac{y^n}{\Gamma(n+1)} \frac{d^n}{dz^n} \Big|_{z=0} \frac{e^{-\frac{(x-z)^2}{2t}}}{\sqrt{2\pi t}} := \frac{e^{-\frac{x^2}{2t}}}{\sqrt{2\pi t}} \left(\frac{y}{t}\right)^n h_n(x, t)$$

whence we can calculate the explicit form of the polynomial  $h_n$ . Namely

$$\begin{aligned} \frac{e^{-\frac{x^2}{2t}}}{\sqrt{2\pi t}} \left(\frac{y}{t}\right)^n h_n(x, t) &= \\ \frac{y^n}{\Gamma(n+1)} \frac{d^n}{dz^n} \Big|_{z=0} \sum_{k=0}^{\infty} \int_{\mathbb{R}} e^{ipx - \frac{tp^2}{2}} \frac{(-ipz)^k}{\Gamma(k+1)} &= \frac{y^n}{\Gamma(n+1)} \int_{\mathbb{R}} e^{ipx - \frac{tp^2}{2}} (-ip)^n \end{aligned}$$

Observing that powers of  $p$  are generated by taking derivatives with respect to  $x$  we get into

$$\frac{e^{-\frac{x^2}{2t}}}{\sqrt{2\pi t}} \left(\frac{y}{t}\right)^n h_n(x, t) = \frac{(-y)^n}{\Gamma(n+1)} \frac{d^n}{dx^n} \int_{\mathbb{R}} e^{ipx - \frac{tp^2}{2}}$$

Performing the integral and contrasting the left to the right hand side yields the claim.  $\square$

The polynomials  $h_n$  defined by (3.1) are called the **Hermite polynomials**. It is readily checked that they enjoy the scaling property

$$h_n(\lambda x, \lambda^2 t) = \lambda^n h_n(x, t) \quad \Rightarrow \quad (x \partial_x + 2t \partial_t) h_n(x, t) = n h_n(x, t)$$

Furthermore

**Proposition 3.2** (*Expected value of Hermite polynomials*).

$$\langle h_n(w_t + x, t) \rangle = \frac{x^n}{\Gamma(n+1)} = h_n(x, 0) \quad (3.2)$$

*Proof.*

$$\langle h_n(w_t + x, t) \rangle := \int_{\mathbb{R}} dy h_n(y, t) \frac{e^{-\frac{(y-x)^2}{2t}}}{\sqrt{2\pi t}} = \frac{(-t)^n}{\Gamma(n+1)} \int_{\mathbb{R}} dy \frac{e^{-\frac{y^2}{2t} + \frac{xy}{t}}}{\sqrt{2\pi t}} \frac{d^n}{dx^n} e^{-\frac{x^2}{2t}}$$

integrating by parts yields the claim.  $\square$

The reason why the expectation value is preserved is that the differential of Hermite along realizations of the Wiener process takes the form.

**Proposition 3.3** (*Stochastic differential of Hermite polynomials*).

$$dh_n(w_t, t) = dw_t \partial_{w_t} h_n(w_t, t)$$

*Proof.* By Ito lemma we have

$$dh_n(w_t, t) = dt \left( \partial_t + \frac{1}{2} \partial_{w_t}^2 \right) h_n(w_t, t) + dw_t \partial_{w_t} h_n(w_t, t)$$

In order to prove the claim we need to show that

$$\left( \partial_t + \frac{1}{2} \partial_{w_t}^2 \right) h_n(w_t, t) = 0$$

Such result can be achieved by direct calculation. It is instructive to proceed in a slightly indirect way. For any  $t > 0$

$$0 = \left( \partial_t - \frac{1}{2} \partial_x^2 \right) \frac{e^{-\frac{(x-y)^2}{2t}}}{\sqrt{2\pi t}} \\ \sum_{n=0}^{\infty} \frac{e^{-\frac{x^2}{2t}}}{\sqrt{2\pi t}} \left(\frac{y}{t}\right)^n \frac{1}{t} \left( -n + t \partial_t - \frac{t}{2} \partial_x^2 + x \partial_x \right) h_n = - \sum_{n=0}^{\infty} \frac{e^{-\frac{x^2}{2t}}}{\sqrt{2\pi t}} \left(\frac{y}{t}\right)^n \left( \partial_t + \frac{1}{2} \partial_x^2 \right) h_n$$

which implies

$$\left( \partial_t + \frac{1}{2} \partial_x^2 \right) h_n = 0$$

as each of these multiply positive definite terms of different order in  $y$ .  $\square$

We have therefore a probabilistic interpretation of the statistical conservation law

$$h_n(w_t + x, t) = h_n(x, 0) + \int_0^t dw_t \partial_{w_t} h_n(w_t, t)$$

From the property of the Ito integral

$$\prec h_n(w_t + x, t) \succ = h_n(x, 0) = \frac{x^n}{\Gamma(n+1)}$$

### 3.1 Recursion relation and multiple integrals over the Wiener process

**Proposition 3.4** (*Integrals over Hermite polynomials*).

$$\int_0^t dw_s h_n(w_s, s) = h_{n+1}(w_t, t)$$

*Proof.* Consider the exponential martingale process (2.1). It satisfies

$$e^{\lambda w_t - \frac{\lambda^2}{2} t} = 1 + \lambda \int_0^t dw_s e^{\lambda w_s - \frac{\lambda^2}{2} s}$$

whence

$$\frac{d^n}{d\lambda^n} \Big|_{\lambda=0} e^{\lambda w_t - \frac{\lambda^2}{2} t} = \frac{d^n}{d\lambda^n} \Big|_{\lambda=0} \lambda \int_0^t dw_s e^{\lambda w_s - \frac{\lambda^2}{2} s}$$

Contrasting the left-hand side with the definition of Hermite polynomials we conclude

$$\frac{d^n}{d\lambda^n} \Big|_{\lambda=0} e^{\lambda w_t - \frac{\lambda^2}{2} t} = t^n \frac{d^n}{dz^n} \Big|_{z=0} e^{\frac{z}{t} w_t - \frac{z^2}{2t}} = \Gamma(n+1) h_n(w_t, t)$$

The right hand side is

$$\frac{d^n}{d\lambda^n} \Big|_{\lambda=0} \lambda \int_0^t dw_s e^{\lambda w_s - \frac{\lambda^2}{2} s} = n \frac{d^{n-1}}{d\lambda^{n-1}} \Big|_{\lambda=0} \int_0^t dw_s e^{\lambda w_s - \frac{\lambda^2}{2} s} = \Gamma(n+1) \int_0^t dw_s h_{n-1}(w_s, s)$$

We have therefore proved that

$$h_n(w_t, t) = \int_0^t dw_s h_{n-1}(w_s, s)$$

□

An important consequence is the following. Since

$$h_0(w_t, t) = 1$$

we have that

$$\int_0^t dw_s = \int_0^t dw_s h_0(w_s, s) = h_1(w_t, t)$$

and

$$\int_0^t dw_{s_1} \int_0^{s_1} dw_{s_0} = h_2(w_t, t)$$

or in full generality

$$\int_0^t dw_{s_1} \prod_{i=1}^{n-1} \int_0^{s_{i-1}} dw_{s_{i-1}} = h_n(w_t, t)$$

## 4 The Stratonovich integral

This section follows section E of chapter 6 of [3]. We have seen that for

$$\theta_k = s t_k + (1 - s) t_{k-1} \quad \forall s \in [0, 1] \quad (4.1)$$

the sum

$$\sum_{k=1}^n w_{\theta_k} (w_{t_k} - w_{t_{k-1}}) = \frac{w_{t_n}^2}{2} - \sum_k \left[ \frac{(w_{t_k} - w_{\theta_k})^2}{2} - \frac{(w_{\theta_k} - w_{t_{k-1}})^2}{2} \right]$$

in  $\mathbb{L}^2(\Omega)$  converges to

$$\int_0^t w_s^{(\theta)} dw_s = \frac{w_t^2}{2} - \frac{t(1-2s)}{2}$$

Choosing  $s = 1/2$  the second term on the right hand side disappears and we recover the result from ordinary calculus. The example suggests to define

**Fisk-Stratonovich integral**  $\int_0^t dw_s \diamond \xi_s := \lim_{|\mathfrak{p}(n)| \downarrow 0} \sum_{t_k \in \mathfrak{p}_n} \frac{\xi_{t_{k-1}+t_k}}{2} (w_{t_k} - w_{t_{k-1}})$  (4.2)

where  $\{\mathfrak{p}_n\}_{n=0}^\infty$  is a sequence of refining partitions of  $[0, t]$ . Note that

$$\begin{aligned} \frac{\xi_{\frac{t_{k+1}+t_k}{2}} - \frac{\xi_{t_{k+1}} + \xi_{t_k}}{2}}{2} &= \\ \frac{\xi_{t_{k+1} - \frac{t_{k+1}-t_k}{2}} - \xi_{t_{k+1}}}{2} + \frac{\xi_{t_k + \frac{t_{k+1}-t_k}{2}} - \xi_{t_k}}{2} &= O(\xi_{t_{k+1}} - \xi_{t_k})^2 \end{aligned}$$

Thus we can equivalently write

$$\int_0^t dw_s \diamond \xi_s := \lim_{|\mathfrak{p}(n)| \downarrow 0} \sum_{t_k \in \mathfrak{p}_n} \frac{\xi_{t_{k+1}} + \xi_{t_k}}{2} (w_{t_k} - w_{t_{k-1}})$$

As in the Ito case the limit converges in **mean square sense**. At variance with the Ito case, the integrand in the definition (4.2) is **anticipating**:

$$\langle \xi_t \diamond dw_t \rangle \neq \langle \xi_t \rangle \langle dw_t \rangle = 0$$

Thus the **martingale property** of the Ito integral is **lost**. To appreciate the advantage of the definition consider

$$\int_0^t dw_s \diamond w_s = \lim_{|\mathfrak{p}(n)| \downarrow 0} \sum_{t_k \in \mathfrak{p}_n} \frac{(w_{t_k} + w_{t_{k-1}})(w_{t_k} - w_{t_{k-1}})}{2} = \frac{w_t^2}{2} \quad (4.3)$$

in agreement with the rules of **ordinary differential calculus**. The example illustrates the general situation.

## 4.1 Relation with the Ito differential

Let us consider

$$\xi_t = g(\chi_t, t) \quad (4.4)$$

with

$$d\chi_t = b(\chi_t, t) dt + \sigma(\chi_t, t) dw_t \quad (4.5)$$

then by Ito lemma we can write

$$d\xi_t = dg(\chi_t, t) = dt \left\{ \partial_t + b_t \partial_{\chi_t} + \frac{\sigma^2}{2} \partial_{\chi_t}^2 \right\} g + dw_t \sigma_t \partial_{\chi_t} g \quad (4.6)$$

and use the this result to establish the relation between the Fisk-Stratonovich and the Ito integral. Namely given a non-anticipating process  $\eta_t$  we can couch the definition of the Fisk-Stratonovich integral into the form

$$\int_0^t dw_s \diamond \eta_s = \lim_{|\mathfrak{p}^{(n)}| \downarrow 0} \sum_{t_k \in \mathfrak{p}^{(n)}} \left\{ \eta_{t_{k-1}} (w_{t_k} - w_{t_{k-1}}) + \frac{(\eta_{t_{k-1}} - \eta_{t_k})(w_{t_k} - w_{t_{k-1}})}{2} \right\}$$

In the literature the latter equality is sometimes written in the continuum limit as

$$\int_0^t dw_s \diamond \eta_s = \int_0^t dw_s \eta_s + \langle \eta, w \rangle_t$$

where  $\langle \xi, w \rangle_t$  is *quadratic co-variation* of the processes  $\xi_t$  and  $w_t$ . The essential point is that in the limit (which converges in the mean square sense under our hypotheses) the quadratic co-variation receives finite contributions only from the term proportional to the increment of the Wiener process

$$dw_t \sim O(\sqrt{dt}) \quad \Rightarrow \quad dw_t^2 \sim O(dt)$$

In such a case if

$$\eta_t = f(\chi_t, t)$$

we find

|   |   |
|---|---|
| <p style="margin: 0;"><b>Stratonovich to Ito conversion</b></p> | $\int_0^t dw_s \diamond f(\chi_s, s) = \int_0^t dw_s f(\chi_s, s) + \frac{1}{2} \int_0^t ds \sigma(\chi_s, s) \partial_{\chi_s} f(\chi_s, s) \quad (4.7)$ |
|---|---|

In particular for

$$f(\chi_t, t) = \sigma(\chi_t, t) \partial_{\chi_t} g(\chi_t, t)$$

we obtain

$$\begin{aligned} dw_t \diamond [\sigma(\chi_t, t) \partial_{\chi_t} g(\chi_t, t)] &= \\ dw_t \sigma(\chi_t, t) \partial_{\chi_t} g(\chi_t, t) &+ \frac{dt}{2} \sigma(\chi_t, t) \partial_{\chi_t} [\sigma(\chi_t, t) \partial_{\chi_t} g(\chi_t, t)] \end{aligned}$$

which allows us to write

$$d\xi_t = dg(\chi_t, t) = dt \left\{ \partial_t + \left[ b - \frac{\sigma}{2} (\partial_{\chi_t} \sigma) \right] \partial_{\chi_t} \right\} g + dw_t \diamond [\sigma \partial_{\chi_t} g] \quad (4.8)$$

As expected, the right hand side does not include any-longer a second derivative of  $g$ , the hallmark of Ito lemma. The function  $g$  is, however, transported by the Stratonovich stochastic differential

$$d\xi_t = dt \left[ b_t - \frac{1}{2} (\sigma \partial_{\chi_t} \sigma)_t \right] + dw_t \sigma_t$$

## 4.2 Examples

- Consider the process

$$\xi_t = w_t^2$$

In such a case the role of the process  $\chi_t$  of the previous section is played by the Brownian motion itself

$$\chi_t = w_t \quad \Rightarrow \quad d\chi_t = dw_t$$

i.e.  $b = 0$  and  $\sigma = 1$  in (4.5). Ito lemma yields

$$d\xi_t = dg(w_t) = 2w_t dw_t + dt$$

The differential admits the **equivalent** Stratonovich representation

$$d\xi_t = 2w_t \diamond dw_t$$

with again  $\chi_t = w_t$ .

- Consider now

$$\xi_t = \chi_t^2$$

with

$$d\chi_t = \chi_t dt + \chi_t dw_t \tag{4.9}$$

This case corresponds to

$$b = \sigma = \chi_t$$

in (4.5). It follows by Ito lemma

$$d\xi_t = 3\chi_t^2 dt + 2\chi_t^2 dw_t$$

On the other hand (4.9) admits the Stratonovich representation

$$d\chi_t = \frac{\chi_t}{2} dt + \chi_t \diamond dw_t$$

whence

$$d\xi_t = \chi_t^2 dt + 2\chi_t^2 \diamond dw_t$$

## References

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