

1 Introduction

The content of these notes is also covered by chapter 3 section C of [1].

2 Properties of the Brownian motion

- Covariance

$$\langle w_t w_s \rangle = \sigma^2 (t \wedge s) := R(t, s) \quad (2.1)$$

- Stationary increments: for any $t, s \in \mathbb{R}_+$

$$\langle (w_t - w_s)^2 \rangle = \sigma^2 |t - s|$$

- Characteristic function: for $t_1 \leq t_2 \leq t_3 \leq \dots \leq t_m$

$$\langle e^{i \sum_{k=1}^{m-1} \lambda_k (w_{t_{k+1}} - w_{t_k})} \rangle = \prod_{k=1}^{m-1} \langle e^{i \lambda_k (w_{t_{k+1}} - w_{t_k})} \rangle = \prod_{k=1}^{m-1} e^{-\frac{\lambda_k^2 \sigma^2}{2} (t_{k+1} - t_k)}$$

3 Karhunen-Loève representation of the Wiener process

The Karhunen-Loève representation provides us with a convenient *deterministic* functional orthonormal basis $\{\psi_n\}_{n=1}^\infty$, which will allow us to write the Brownian motion

$$w_t : \Omega \times [0, T] \rightarrow \mathbb{R}$$

as a series

$$w_t = \sum_{n=0}^{\infty} c_n \psi_n(t)$$

Randomness is encoded in the coefficients

$$c_n = \int_0^T \frac{dt}{T} \psi_n(t) w_t$$

which need to form a sequence $\{c_n\}_{n=0}^\infty$ of independent Gaussian random variable with zero average and variance determined by the Karhunen-Loève representation.

Let us consider the space of the Lebesgue square integrable $\mathbb{L}^2([0, T])$ real functions on $[0, T]$. This space is an Hilbert space with respect to the scalar product:

$$\langle f, g \rangle := \int_0^T \frac{dt}{T} f(t)g(t) \quad f, g \in \mathbb{L}^2([0, T])$$

Proposition 3.1. *The covariance R of the Brownian motion (2.1) defines the kernel of an operator mapping $\mathbb{L}^2([0, T])$ into itself:*

$$g(t) = \int_0^T \frac{dt}{T} R(t, s) f(s) := R[f](t)$$

i.e.

$$f \in \mathbb{L}^2([0, T]) \rightarrow g \equiv R[f] \in \mathbb{L}^2([0, T])$$

Proof. The following chain of equalities holds by definition

$$\|g\|_2^2 := \langle R[f], R[f] \rangle = \int_0^T \frac{dt}{T} \left[\int_0^T \frac{ds}{T} R(t, s) f(s) \right] \left[\int_0^T \frac{ds'}{T} R(t, s') f(s') \right]$$

where

$$\int_0^T \frac{ds}{T} R(t, s) f(s) = \sigma^2 \int_0^t \frac{ds}{T} s f(s) + \sigma^2 t \int_t^T \frac{ds}{T} f(s)$$

so that by Lyapunov inequality

$$|R[f](t)| \leq \sigma^2 T \int_0^T \frac{ds}{T} |f(s)| \leq \sigma^2 T \left(\int_0^T \frac{ds}{T} |f(s)|^2 \right)^{1/2}$$

We conclude

$$\|g\|_2^2 \leq \sigma^4 T^2 \|f\|_2^2$$

□

The as an operator kernel $R(t, s)$ enjoys two further properties

- $R(t, s)$ is positive definite: i.e. for any collection of $\{c_i\}_{i=1}^n \in \mathbb{C}$ and sampling of the arguments $\{t_i\}_{i=1}^n$ we have

$$\sum_{ij=1}^n c_i R(t_i, s_i) c_j^* \geq 0$$

Namely

$$\sum_{ij=1}^n c_i R(t_i, t_j) c_j^* = \sum_{i,j=1}^n \langle w_{t_i} w_{t_j} \rangle c_i c_j^* = \langle \left| \sum_{i=1}^n c_i w_{t_i} \right|^2 \rangle \geq 0$$

- The kernel is symmetric:

$$R(t, s) = R(s, t)$$

These properties of the kernel allows us to invoke

Theorem 3.1 (Mercer). *Let R be a continuous symmetric non-negative definite kernel. Then there is an orthonormal basis $\{\psi_n\}_{n=0}^\infty$ of $\mathbb{L}^2([0, T])$ consisting of eigenfunctions of the operator*

$$R[f] = \int_0^T \frac{ds}{T} R(t, s) f(s)$$

such that the corresponding sequence of eigenvalues $\{r_n\}_{n=0}^\infty$ is non-negative. The eigenfunctions corresponding to non-zero eigenvalues are continuous on $[0, T]$ and R admit the representation

$$R(t, s) = \sum_{n=0}^{\infty} r_n \psi_n(t) \psi_n(s)$$

Proof. The theorem is proved at pag. 138 of [2]. It can also be regarded as a consequence of the Hilbert-Schmidt theorem [3] \square

In other words the operator specified by R is diagonalizable with discrete spectrum. The idea underlying the Karhunen-Loève representation is to use the orthonormal basis of eigenvectors to write the Brownian motion as the \mathbb{L}^2 -convergent series:

$$w_t = \sum_{n=0}^{\infty} c_n \psi_n(t) \quad (3.1)$$

Namely, if the $\{c_n\}_{n=0}^{\infty}$ are Gaussian independent random variables satisfying

$$\langle c_n \rangle = 0 \quad \& \quad \langle c_n^2 \rangle = r_n$$

it follows immediately that (3.1) satisfies all the requirements to be a Brownian motion:

i Covariance:

$$\langle w_t w_s \rangle = \sum_{n=0}^{\infty} r_n \psi_n(t) \psi_n(s) = R(t, s)$$

ii Independent increments: for $t_1 \leq t_2 \leq t_3 \leq t_4$

$$\langle (w_{t_4} - w_{t_3})(w_{t_2} - w_{t_1}) \rangle = \sigma^2 \sum_n r_n [t_4 \wedge t_2 - t_4 \wedge t_1 - t_3 \wedge t_2 + t_3 \wedge t_1] = 0$$

iii Gaussian structure of the characteristic function :

$$\langle e^{i \sum_{k=1}^m \sum_{n=0}^{\infty} c_n \lambda_k [\psi_n(t_k) - \psi_n(t_{k-1})]} \rangle = e^{-\frac{1}{2} \sum_n r_n \{ \sum_{k=1}^m \lambda_k [\psi_n(t_k) - \psi_n(t_{k-1})] \}^2}$$

We use the identity

$$\begin{aligned} \sum_{k=1}^m \lambda_k [\psi_n(t_k) - \psi_n(t_{k-1})] &= \\ \sum_{k=1}^{m-1} (\lambda_k - \lambda_{k+1}) \psi_n(t_k) + \lambda_m \psi_n(t_m) - \lambda_1 \psi_n(t_0) &:= \sum_{k=0}^m \Delta_k \psi_n(t_k) \end{aligned}$$

Thus

$$\langle e^{i \sum_{k=1}^m \sum_{n=0}^{\infty} c_n \lambda_k [\psi_n(t_k) - \psi_n(t_{k-1})]} \rangle = e^{-\frac{1}{2} \sum_n r_n [\sum_{k=1}^m \Delta_k \psi_n(t_k)]^2}$$

Now we observe that

$$\begin{aligned} \sum_n r_n \left[\sum_{k=1}^m \Delta_k \psi_n(t_k) \right]^2 &= \\ \sum_{k,l=1}^m \sum_n r_n \Delta_k \Delta_l \psi_n(t_k) \psi_n(t_l) &= \sum_{k,l=1}^m \Delta_k \Delta_l R(t_k, t_l) = \langle \left(\sum_{l=0}^m \Delta_l w_{t_l} \right)^2 \rangle \end{aligned}$$

Finally using the definition of the Δ_l coefficients

$$\sum_n r_n \left[\sum_{k=1}^m \Delta_k \psi_n(t_k) \right]^2 = \langle \left[\sum_{k=1}^m \lambda_k (w_{t_k} - w_{t_{k-1}}) \right]^2 \rangle = \sum_{k=1}^m \langle \lambda_k^2 (w_{t_k} - w_{t_{k-1}})^2 \rangle$$

3.1 Explicit construction of the Karhunen-Loève basis

We need to solve:

$$\int_0^T \frac{ds}{T} R(t, s) \psi_n(s) = r_n \psi_n(t)$$

or equivalently

$$\int_0^t \frac{ds}{T} ds \sigma^2 s \psi_n(s) + \sigma^2 t \int_t^T \frac{ds}{T} \psi_n(s) = r_n \psi_n(t)$$

The first derivative is

$$\sigma^2 \int_t^T \frac{ds}{T} \psi_n(s) = r_n \dot{\psi}_n(t) \quad \Rightarrow \quad \dot{\psi}_n(T) = 0$$

and the second derivative

$$-\frac{\sigma^2}{T} \psi_n(t) = r_n \ddot{\psi}_n(t) \quad \Rightarrow \quad \ddot{\psi}_n(t) + \frac{\sigma^2}{T r_n} \psi_n(t) = 0$$

Thus the problem is equivalent to solving the differential equation

$$\ddot{\psi}_n(t) + \frac{\sigma^2}{T r_n} \psi_n(t) = 0$$

with boundary conditions

$$\psi_n(0) = \dot{\psi}_n(T) = 0$$

The boundary condition in zero yields (the factor $\sqrt{2}$ comes from unit normalization)

$$\psi_n(t) = \sqrt{2} \sin \left(\sqrt{\frac{\sigma^2}{T r_n}} t \right)$$

Imposing the condition for $t = T$ gives the “*quantization*” condition for the eigenvalues:

$$\sqrt{\frac{\sigma^2 T}{r_n}} = (2n + 1) \frac{\pi}{2} \quad n = 0, 1, \dots$$

The conclusion is that the explicit Karhunen-Loève representation of the covariance of the Wiener process is

$$R(t, s) = \sum_{n=0}^{\infty} \frac{\sigma^2 T \psi_n(t) \psi_n(s)}{(n + \frac{1}{2})^2 \pi^2}$$

4 Final remark

An example of concrete application of the Karhunen-Loève representation can be found in [4].

References

- [1] L.C. Evans, *An Introduction to Stochastic Differential Equations*, lecture notes, <http://math.berkeley.edu/~evans/>
- [2] R. Courant, and D. Hilbert, *Methods of Mathematical Physics, Vol I* Wiley, (1989).
- [3] A. N. Kolmogorov, S. V. Fomin *Elements of the Theory of Functions and Functional Analysis* Dover Publications, (1999).
- [4] T. Sapsis, and P. Lermusiaux, *Dynamically Orthogonal field equations for continuous stochastic dynamical systems*, *Physica D*, **238** 2347 (2009) and http://web.mit.edu/sapsis/www/papers/DO_fld_eqs.pdf