

1 Introduction

The lecture notes cover sections F and G of chapter 2 and appendix Appendix A of [1]. The notes also discuss some extra examples.

2 Borel -Cantelli lemma

Let $\{F_k\}_{k=1}^{\infty}$ a sequence of events in a probability space.

Definition 2.1 (F_n *infinitely often*). The event specified by the simultaneous occurrence an infinite number of the events in the sequence $\{F_k\}_{k=1}^{\infty}$ is called “ F_n infinitely often” and denoted F_n i.o.. In formulae

$$F_n \text{ i.o.} := \bigcap_{n=1}^{\infty} \bigcup_{k=n}^{\infty} F_k = \{\omega \in \Omega \mid \omega \text{ belongs to infinitely many of the } F_n\}$$

An alternative notation may help. Observe that the set-theoretic union (\cup) operation has the probabilistic meaning of “at least”. The event $\tilde{F}_n = \bigcup_{k=n}^{\infty} F_k$ defines a sampling of the tail (starting from n) of the sequence. It occurs if at least one event in the tail occurs. The intersection $\bigcap_{n=1}^{\infty} \tilde{F}_n$ differs from the empty set if disregarding how long is the tail (how large is n) we can observe the occurrence of an $F_{n'}$ for $n' \geq n$. This is possible only if an infinite number of the events in the sequence indeed occur:

$$F_n \text{ i.o.} = \lim_{n \uparrow \infty} \bigcup_{k \geq n} F_k = \lim_{n \uparrow \infty} \sup_{k \geq n} F_k$$

Recall also that

$$P(F_k) = \int dP \chi_{F_k}(\omega)$$

for χ_{F_k} the characteristic function of the event χ and that

$$\lim_{n \uparrow \infty} \sup_{k \geq n} \chi_{F_k}(\omega) = \chi_{\lim_{n \uparrow \infty} \sup_{k \geq n} F_k}(\omega)$$

Lemma 2.1 (*Borel-Cantelli*). The following claims hold:

- if $\sum_n P(F_n) < \infty$ then $P(F_n \text{ i.o.}) = 0$
- if $\sum_n P(F_n) = \infty$ and $\{F_n\}_{n=1}^{\infty}$ consists of independent events $P(F_n \text{ i.o.}) = 1$

Proof. :

- By definition

$$P(F_n \text{ i.o.}) = P(\bigcap_{n=0}^{\infty} \bigcup_{k=n}^{\infty} F_k) = P(\lim_{n \uparrow \infty} \bigcup_{k=n}^{\infty} F_k) \tag{2.1}$$

Lebesgue’s dominated convergence theorem (see e.g. [2] pag. 187), allows us to carry out the limit from the integral

$$P(F_n \text{ i.o.}) = \lim_{n \uparrow \infty} P(\bigcup_{k=n}^{\infty} F_k) \leq \lim_{n \uparrow \infty} \sum_{k=n}^{\infty} P(F_k) = 0$$

whilst the definition of probability measure enforces the inequality

$$P(F_n \text{ i.o.}) \leq \lim_{n \uparrow \infty} \sum_{k=n}^{\infty} P(F_k)$$

The proof of the first statement follows from the hypothesized convergence of the series.

- We can turn to the complementary event:

$$(\cup_{k \geq n} F_k)^c = \cap_{k \geq n} F_k^c$$

and use independence

$$P(\cap_{k \geq n} F_k^c) = \prod_{k=n}^{\infty} P(F_k^c) = \prod_{k=n}^{\infty} [1 - P(F_k)]$$

The inequality

$$1 - x \leq e^{-x}$$

then provides us with an upper bound for each factor in the product

$$P(\cap_{k \geq n} F_k^c) \leq \prod_{k=n}^{\infty} e^{-P(F_k)} = e^{-\sum_{k=n}^{\infty} P(F_k)}$$

whence the claim follows if the series diverges.

□

The Borel-Cantelli lemma provides an extremely useful tool to prove asymptotic results about random sequences holding *almost surely* (acronym: *a.s.*). This means that such results hold true but for events of zero probability. An obvious synonym for *a.s.* is then *with probability one*.

3 Characteristic function of a random variable

Definition 3.1 (*Characteristic function*). Let

$$\xi : \Omega \rightarrow \mathbb{R}^d$$

the expectation value

$$\check{p}_{\xi}(\mathbf{q}) := \langle e^{i\boldsymbol{\xi} \cdot \mathbf{q}} \rangle$$

is referred to as the characteristic function of the random variable

Example 3.1 (*Characteristic function of a Gaussian random*). Let

$$\xi : \Omega \rightarrow \mathbb{R}$$

distributed with Gaussian PDF. The characteristic function is

$$\check{g}_{\bar{x}, \sigma}(q) = \int dx e^{iqx} g_{\bar{x}, \sigma}(x) = e^{iq\bar{x} - \frac{\sigma^2 q^2}{2}}$$

For a Gaussian variable it is also true

$$\check{G}_{\bar{x}, \sigma}(q) = \frac{i^n q^n}{\Gamma(n+1)} \langle \xi^n \rangle$$

having used the Γ -function representation of the factorial (see appendix B). The remaining expectation value is

$$\langle \xi^n \rangle = \frac{\Gamma(2n+1)}{2^n \Gamma(n+1)} = (2n-1)!! \quad (3.1)$$

Formally for a random variable ξ one can write

$$\langle \xi^n \rangle = \frac{1}{i^n} \frac{d^n}{dq^n} \check{p}_\xi(q) \Big|_{q=0} \quad (3.2)$$

The relation is formal because it may be a relation between infinities.

Example 3.2 (Lorentz distribution). Let

$$p : \mathbb{R} \rightarrow \mathbb{R}_+$$

be

$$p_{y\sigma}(x) = \frac{\sigma}{\pi \{(x-y)^2 + \sigma^2\}}$$

the Lorentz probability density so that $(\mathbb{R}, \mathcal{B}, P_{y\sigma}(\mathcal{B}))$ a probability space. Note that

$$p_{0\sigma}(x) = p_{0\sigma}(-x)$$

Using a change of variable and it is straightforward to verify that

$$\int_{\mathbb{R}} dx x p_{y\sigma}(x) = y$$

however

$$\int_{\mathbb{R}} dx x^2 p_{y\sigma}(x) = \infty$$

The characteristic function can be computed using Cauchy theorem

$$\begin{aligned} \check{p}_{y\sigma}(q) &= e^{iqy} \int dx e^{iqx} p_{0\sigma}(x) \\ &= \frac{e^{iqy}}{2i\pi} \int_{\mathbb{R}} dx e^{iqx} \left\{ \frac{1}{x-i\sigma} - \frac{1}{x+i\sigma} \right\} = e^{iqy} \begin{cases} e^{-q\sigma} & \text{if } q > 0 \\ e^{q\sigma} & \text{if } q < 0 \end{cases} \end{aligned}$$

The characteristic function develops a cusp for $q = 0$

$$\check{p}_{y\sigma}(q) = e^{iqy - \sigma|q|}$$

Finally note that

$$\delta(x-y) \stackrel{w}{=} \lim_{\sigma \downarrow 0} p_{y\sigma}(x)$$

Remark 3.1 (The Fourier representation of the δ -Dirac). Let $f_i : \mathbb{R}^d \rightarrow \mathbb{R}$, $i = 1, 2$ some smooth integrable functions and

$$\check{f}_i(\mathbf{q}) = \int_{\mathbb{R}^d} d^d x e^{i\mathbf{q}\cdot\mathbf{x}} f_i(\mathbf{x}) \quad i = 1, 2$$

their Fourier transform. The convolution identity

$$(f_1 \star f_2)(\mathbf{x}) := \int d^d y f_1(\mathbf{x} - \mathbf{y}) f_2(\mathbf{x}) = \int \frac{d^d q}{(2\pi)^d} e^{i\mathbf{q}\cdot\mathbf{x}} \check{f}_1(\mathbf{q}) \check{f}_2(\mathbf{q})$$

maybe thought as a consequence of

$$\delta^{(d)}(\mathbf{x} - \mathbf{y}) \stackrel{w}{=} \int \frac{d^d q}{(2\pi)^d} e^{i\mathbf{q}\cdot(\mathbf{x}-\mathbf{y})}$$

The δ -Dirac in such a case could be very formally thought as the characteristic function of a "random variable" uniformly distributed over \mathbb{R}^d .

4 Probability density and Dirac- δ

Let

$$\boldsymbol{\xi} : \Omega \rightarrow \mathbb{R}^d$$

with PDF $p_{\boldsymbol{\xi}}(\cdot)$. From the properties of the δ function we have

$$\int_{\mathbb{R}^d} \delta^{(d)}(\boldsymbol{\xi} - \mathbf{x}) \boldsymbol{\gamma} = \int_{\mathbb{R}^d} d^d \mathbf{y} p_{\boldsymbol{\xi}}(\mathbf{y}) \delta^{(d)}(\mathbf{y} - \mathbf{x}) = p_{\boldsymbol{\xi}}(\mathbf{x})$$

This relation allows us to derive the relation between the PDF's of functionally dependent random variables

4.1 Derivation in one dimension

Suppose the random variable has the

$$P_{\xi}(x < \xi < x + dx) = p_{\xi}(x) dx$$

Functional relation between random variables

$$\phi = f(\xi) \tag{4.1}$$

again

$$P_{\phi}(y < \phi < y + dy) = p_{\phi}(y) dy$$

From

$$P_{\xi}(x < \xi < x + dx) = P_{\phi}(y < \phi < y + dy)$$

one gets into

$$p_{\xi}(x) = p_{\phi}(f(x)) \frac{df}{dx}$$

4.2 Multi-dimensional case using the δ -function

Suppose f is **one-to-one** and write

$$p_{\phi}(\mathbf{y}) = \int_{\mathbb{R}^d} \delta^{(d)}(\phi - \mathbf{y}) \boldsymbol{\gamma} = \int_{\mathbb{R}^d} \delta^{(d)}(f(\boldsymbol{\xi}) - \mathbf{y}) \boldsymbol{\gamma}$$

It follows

$$p_{\phi}(\mathbf{y}) = \int_{\mathbb{R}^d} d^d \mathbf{x} \delta^{(d)}(f(\mathbf{x}) - \mathbf{y}) p_{\boldsymbol{\xi}}(\mathbf{x}) \equiv \lim_{\sigma \downarrow 0} \int_{\mathbb{R}^d} d^d \mathbf{x} \frac{e^{-\frac{|f(\mathbf{x}) - \mathbf{y}|^2}{2\sigma^2}}}{(2\pi\sigma^2)^{\frac{d}{2}}} p_{\boldsymbol{\xi}}(\mathbf{x})$$

Since f is one-to-one

$$f^{-1}(\mathbf{y}) = \mathbf{x}_{\star} \tag{4.2}$$

is globally well defined. Taylor-expanding the argument of the exponential we get for the i -th component of \mathbf{y}

$$y^i = f^i(\mathbf{x}_{\star}) + (x^j - x_{\star}^j) \frac{\partial f^i}{\partial x^j}(\mathbf{x}_{\star}) + O((x^j - x_{\star}^j)^2)$$

Call

$$A_{ij} := \frac{\partial f^i}{\partial x^j}(\mathbf{x}_\star)$$

$$z^i = \frac{x^i - x_\star^i}{\sigma}$$

the integral becomes

$$p_\phi(\mathbf{y}) = \lim_{\sigma \downarrow 0} \int_{\mathbb{R}^d} d^d z \frac{e^{-\frac{z^i A_{li} A_{lj} z^j + O(\sigma)}{2}}}{(2\pi)^{\frac{d}{2}}} p_\xi(\mathbf{x}_\star + O(\sigma)) = \frac{p_\xi(\mathbf{x}_\star)}{|\det A|} \quad (4.3)$$

5 Limit theorems for Bernoulli variables

Definition 5.1 (*i.i.d. random variables*). A sequence of random variables $\{\xi_i\}_{i=1}^\infty$ is said identically distributed if

$$P_{\xi_1}(x) = P_{\xi_2}(x) = \dots = P_{\xi_n}(x) = \dots \quad (5.1)$$

Furthermore if they are mutually independent they are usually referred to with the acronym *i.i.d.*

Let $\{\xi_i\}_{i=1}^\infty$ a sequence of i.i.d. Bernoulli variables i.e. for all i , $\xi_i \stackrel{d}{=} \xi$ (*equality in distribution*) and

$$\xi : \Omega \rightarrow \{-x, x\} \quad \& \quad P_\xi(x) = p$$

From the characteristic function

$$\langle e^{it\xi} \rangle = \cos(xt) + i(2p-1) \sin(xt)$$

we find

$$\langle \xi^n \rangle = \begin{cases} x^n & n = 2k \\ x^n(2p-1) & n = 2k+1 \end{cases}$$

whence

$$\langle \xi \rangle = (2p-1)x \quad \& \quad \langle (\xi - \langle \xi \rangle)^2 \rangle = 4p(1-p)x^2$$

and

$$\langle (\xi - \langle \xi \rangle)^4 \rangle = 16p(1-p)[1 - 3(1-p)p]$$

If we introduce the random variable

$$S_n := \frac{\sum_{i=1}^n \xi_i}{n}$$

then

$$\langle S_n \rangle = \langle \xi \rangle \quad \& \quad \langle (S_n - \langle \xi \rangle)^2 \rangle = \frac{4p(1-p)}{n}$$

- We can apply Chebyshev lemma to show

$$P(|S_n - \langle \xi \rangle| \geq \varepsilon) \leq \frac{\langle (S_n - \langle \xi \rangle)^2 \rangle}{\varepsilon^2} = \frac{4p(1-p)}{n\varepsilon^2} \xrightarrow{n \uparrow \infty} 0 \quad (5.2)$$

This is *law of large numbers* for i.i.d. Bernoulli variables.

- Upon setting

$$\tilde{\xi}_i = \xi_i - \langle \xi \rangle \quad (5.3)$$

$$\begin{aligned} \langle (S_n - \langle \xi \rangle)^4 \rangle &= \sum_{i=1}^n \frac{\langle \tilde{\xi}_i^4 \rangle}{n^4} + 3 \sum_{ijkl} \delta_{ij} \delta_{kl} (1 - \delta_{jk}) \frac{\langle \tilde{\xi}_i^2 \rangle^2}{n^4} \\ &= \frac{16p(1-p)}{n^3} [1 - 3(1-p)p] + 3n(n-1) \frac{16p^2(1-p)^2}{n^4} \leq \frac{C}{n^2} \end{aligned}$$

for $n \gg 1$, we also find

$$P(|S_n - \langle \xi \rangle| \geq \varepsilon) \leq \frac{\langle (S_n - \langle \xi \rangle)^4 \rangle}{\varepsilon^4} \leq \frac{C}{n^2 \varepsilon^4} \xrightarrow{n \uparrow \infty} 0$$

The new bound tends to zero sufficiently fast at infinity to be useful in order to apply the Borel-Cantelli lemma which entitles us to conclude:

$$S_n \xrightarrow{n \uparrow \infty} \langle \xi \rangle \quad a.s.$$

i.e. we have proved the *strong law of large numbers* for Bernoulli schemes.

- Let us convene that S_0 is zero. Then we can write (see appendix A)

$$P\left(S_n = \frac{mx}{n}\right) = \frac{\Gamma(n+1)}{\Gamma\left(\frac{n+m}{2}+1\right)\Gamma\left(\frac{n-m}{2}+1\right)} p^{\frac{n+m}{2}} (1-p)^{\frac{n-m}{2}}$$

For $n, m \gg 1$ we can extricate the asymptotics of the probability using Stirling's formula (see formula B.1 in appendix (B)):

$$P\left(S_n = \frac{mx}{n}\right) \simeq \frac{1}{\sqrt{2\pi \frac{n^2-m^2}{n}}} e^{\frac{n+m}{2} \ln p + \frac{n-m}{2} \ln(1-p) + n(\ln n - 1) - \frac{n+m}{2}(\ln \frac{n+m}{2} - 1) - \frac{n-m}{2}(\ln \frac{n-m}{2} - 1)} \quad (5.4)$$

The strong law of large numbers allows us to relate asymptotically m to the *empirical probability* (i.e. the observed frequency) of a displacement to the right

$$m = (2\tilde{p} - 1)n$$

As n tends to infinity the observation allows us to recast (5.4) into the form

$$P\left(S_n = \frac{mx}{n}\right) \simeq \frac{e^{-nK(\tilde{p}|p) + o(n)}}{\sqrt{2\pi \frac{4\tilde{p}(1-\tilde{p})}{n}}}$$

with

$$K(\tilde{p}|p) = - \left\{ \tilde{p} \ln \frac{\tilde{p}}{p} + (1 - \tilde{p}) \ln \frac{1 - \tilde{p}}{1 - p} \right\} \quad (5.5)$$

the *Kullback-Leibler divergence (entropy)* between the empirical and the Bernoulli distribution. This quantity measure the (rate of) discrepancy between two probability measures.

- Let z be an indicator of the discrepancy between p and \tilde{p} i.e.

$$\tilde{p} = p + \frac{z}{2}$$

Taylor expanding (5.5) around z equal zero yields

$$K(\tilde{p}|p) = \frac{z^2}{8p(1-p)} + O(z^3) \quad (5.6)$$

We can use this result to estimate the probability with which S_n asymptotically deviates from its expected value

$$P\left(\frac{S_n - \xi}{\sqrt{(S_n - \xi)^2}} = z\right) \stackrel{n \uparrow \infty}{\simeq} \frac{e^{-\frac{z^2}{2}}}{\sqrt{2\pi \frac{(S_n - \xi)^2}{x^2}}}$$

This is the content of the *central limit theorem* for Bernoulli variables. It is tempting to interpret the vanishing of the variance in the denominator as n tends to infinity, as a weight in a Riemann sum so to infer

$$P\left(\frac{S_n - \xi}{\sqrt{(S_n - \xi)^2}} \leq z\right) \stackrel{n \uparrow \infty}{\simeq} \int_{-\infty}^z du \frac{e^{-\frac{u^2}{2}}}{\sqrt{2\pi}}$$

- Consider now the characteristic function

$$\langle e^{iS_n t} \rangle = \langle e^{i\xi t} \rangle^n = \left[\cos\left(\frac{x t}{n}\right) + i(2p-1) \sin\left(\frac{x t}{n}\right) \right]^n$$

The limit

$$\lim_{n \uparrow \infty} \ln \langle e^{iS_n t} \rangle = i(2p-1) x t - \frac{x^2 t^2}{2} \frac{4p(p-1)}{n} + o\left(\frac{1}{n}\right)$$

contains the same type of information of the central limit theorem. Namely we can couch the result into the form

$$\begin{aligned} \lim_{n \uparrow \infty} \ln \langle e^{iS_n t} \rangle &= i \langle \xi \rangle t - \frac{t^2}{2} \frac{\langle (\xi - \langle \xi \rangle)^2 \rangle}{n} + o\left(\frac{1}{n}\right) \\ &= i \langle S_n \rangle t - \frac{t^2}{2} \frac{\langle (S_n - \langle S_n \rangle)^2 \rangle}{n} + o\left(\frac{1}{n}\right) \end{aligned}$$

The rescaling

$$t \rightarrow \frac{t}{\sqrt{\langle (S_n - \langle S_n \rangle)^2 \rangle^{1/2}}}$$

suggests that the characteristic function

$$\langle e^{i \frac{S_n - \langle S_n \rangle}{\sqrt{\langle (S_n - \langle S_n \rangle)^2 \rangle^{1/2}}} t} \rangle \stackrel{n \uparrow \infty}{\rightarrow} e^{-\frac{t^2}{2}}$$

tends indeed to the characteristic function of the Gaussian distribution.

Appendices

A Random walk

Let

$$R_n = n S_n$$

and suppose that of the n steps n_r were to the right and n_l to the left:

$$n = n_l + n_r \tag{A.1}$$

The displacement from the origin in units of x is then given by

$$m = n_r - n_l \tag{A.2}$$

We can solve for n_l , n_r and obtain

$$n_r = \frac{n + m}{2} \quad \& \quad n_l = \frac{n - m}{2}$$

The probability of an *individual sequence* of samples of Bernoulli variables such that $R_n(\omega) = m x$ is

$$P(\omega) = p^{\frac{n+m}{2}} (1-p)^{\frac{n-m}{2}}$$

In order to evaluate the total probability $P(R_n = m x)$ we must *count* all possible sequences of samples such that (A.1), (A.2) are verified. This number is equal to the number of ways we can extract n_r out of n *indistinguishable object* (this means that the extraction order does not matter):

$$C_n^{n_r} = \frac{n!}{n_r!(n-n_r)!} \equiv \frac{n!}{n_r!n_l!}$$

The conclusion is

$$P(R_n = m x) = \frac{n!}{\left(\frac{n+m}{2}\right)! \left(\frac{n-m}{2}\right)!} p^{\frac{n+m}{2}} (1-p)^{\frac{n-m}{2}}$$

Using binomial formula it is straightforward to check the normalization condition for $m = -n, -(n-1), \dots, n-1, n$.

B Gamma function

The Γ function for any $x \in \mathbb{R}_+$ is specified by the integral

$$\Gamma(x) = \int_0^\infty \frac{dy}{y} y^x e^{-y}$$

For $x \in \mathbb{N}$ the integral can be performed explicitly and it is equal to the factorial:

$$\Gamma(x) = (x-1)! \quad x \in \mathbb{N}$$

For $x \in \mathbb{R}_+$, integration by parts yields the identity

$$\Gamma(x+1) = \int_0^\infty \frac{dy}{y} y^{x+1} e^{-y} = x \Gamma(x)$$

which is trivially satisfied by factorials. For $x \gg 1$ the value of the integral is approximated by *Laplace's stationary point* method

$$\Gamma(x+1) \simeq e^{x(\ln x - 1)} \int_{\mathbb{R}} dy e^{-\frac{y^2}{2x}} = \sqrt{2\pi x} e^{x(\ln x - 1)} \quad x \gg 1 \tag{B.1}$$

Such asymptotic estimation is usually referred to as *Stirling formula*.

References

- [1] L.C. Evans, *An Introduction to Stochastic Differential Equations*, lecture notes, <http://math.berkeley.edu/~evans/>
- [2] A. N. Shiryaev, *Probability*, 2nd Ed. Springer (1996), <http://books.google.com/>