

Time Frequency Analysis - Winter 2012

EXERCISE SET 5

5.1. Investigate the commutation relations between S_ξ and T_y , M_η , D_λ^2 : Find a ξ' (possibly different in the different identities below) so that

$$S_\xi T_y = T_y S_{\xi'}, \quad S_\xi M_\eta = M_\eta S_{\xi'}, \quad S_\xi D_\lambda^2 = D_\lambda^2 S_{\xi'}.$$

Finally, find a value of $\eta = \eta(\xi, \lambda)$ such that $S_\xi D_\lambda^2 M_\eta = D_\lambda^2 M_\eta S_\xi$. We begin by recalling the definition of the operator S_ξ :

$$S_\xi f(x) = \int_{-\infty}^{\xi} \hat{f}(\eta) e^{2\pi i x \eta} d\eta.$$

We now have

$$\begin{aligned} S_\xi T_y f(x) &= \int_{-\infty}^{\xi} \widehat{T_y f}(\eta) e^{2\pi i \eta x} d\eta = \int_{-\infty}^{\xi} M_{-y} \hat{f}(\eta) e^{2\pi i \eta x} d\eta \\ &= \int_{-\infty}^{\xi} \hat{f}(\eta) e^{2\pi i \eta (x-y)} d\eta = S_\xi f(x-y) = T_y S_\xi f(x). \\ S_\xi M_\eta f(x) &= \int_{-\infty}^{\xi} \widehat{M_\eta f}(s) e^{2\pi i s x} ds = \int_{-\infty}^{\xi} T_\eta \hat{f}(s) e^{2\pi i s x} ds \\ &= \int_{-\infty}^{\xi} \hat{f}(s-\eta) e^{2\pi i s x} ds = \int_{-\infty}^{\xi-\eta} \hat{f}(s) e^{2\pi i (\eta+s)x} ds \\ &= M_\eta S_{\xi-\eta} f(x). \\ S_\xi D_\lambda^2 f &= \int_{-\infty}^{\xi} \widehat{D_\lambda^2 f}(\eta) e^{2\pi i \eta x} d\eta = \int_{-\infty}^{\xi} D_{\lambda^{-1}}^2 \hat{f}(\eta) e^{2\pi i \eta x} d\eta \\ &= \int_{-\infty}^{\xi} \lambda^{\frac{1}{2}} \hat{f}(\lambda \eta) e^{2\pi i \eta x} d\eta = \lambda^{\frac{1}{2}} \int_{-\infty}^{\lambda \xi} \hat{f}(\eta) e^{2\pi i x \eta / \lambda} \frac{d\eta}{\lambda} \\ &= \lambda^{-\frac{1}{2}} \int_{-\infty}^{\lambda \xi} \hat{f}(\eta) e^{2\pi i \eta \frac{x}{\lambda}} d\eta = D_\lambda^2 S_{\lambda \xi} f(x). \end{aligned}$$

Using the previous identities we can write

$$S_\xi D_\lambda^2 M_\eta f(x) = D_\lambda^2 S_{\lambda \xi} M_\eta f(x) = D_\lambda^2 M_\eta S_{\lambda \xi - \eta} f(x),$$

so we get the desired identity with $\eta = \lambda \xi - \xi = \eta(\xi, \lambda)$.

5.2. We have used several times the identity $\langle f, g \rangle = \langle \hat{f}, \hat{g} \rangle$, where $\langle f, g \rangle = \int_{\mathbb{R}} f \bar{g}$ is the L^2 inner product. Prove this identity in the following two ways:

- (a) Write the identity $\|h\|_{L^2} = \|\hat{h}\|_{L^2}$ for $h = f + ug$, where $u \in \{1, -1, i, -i\}$.
- (b) In the identity $\int_{\mathbb{R}} f \hat{h} = \int_{\mathbb{R}} \hat{f} h$ substitute $f = \hat{f}$, and manipulate the right

hand side. (a) We have

$$\begin{aligned} \|f + ug\|_2^2 &= \int_{\mathbb{R}} (f + ug)\overline{(f + ug)} = \int_{\mathbb{R}} (|f|^2 + |u|^2|g|^2 + f\bar{u}\bar{g} + ug\bar{f}) \\ &= \int_{\mathbb{R}} |f|^2 + \int_{\mathbb{R}} |g|^2 + \bar{u} \int_{\mathbb{R}} f\bar{g} + u \int_{\mathbb{R}} g\bar{f}. \end{aligned}$$

Using the same expansion for $\widehat{f + ug}$ and Plancherel's theorem $\|f + ug\|_2^2 = \|\widehat{f + ug}\|_2^2$ we get

$$\bar{u} \int_{\mathbb{R}} f\bar{g} + u \int_{\mathbb{R}} \bar{f}g = \bar{u} \int_{\mathbb{R}} \hat{f}\hat{\bar{g}} + u \int_{\mathbb{R}} \hat{\bar{f}}\hat{g}$$

For $u = 1$ we get

$$\int_{\mathbb{R}} f\bar{g} + \int_{\mathbb{R}} \bar{f}g = \int_{\mathbb{R}} \hat{f}\hat{\bar{g}} + \int_{\mathbb{R}} \hat{\bar{f}}\hat{g},$$

while for $u = i$

$$-i \int_{\mathbb{R}} f\bar{g} + i \int_{\mathbb{R}} \bar{f}g = -i \int_{\mathbb{R}} \hat{f}\hat{\bar{g}} + i \int_{\mathbb{R}} \hat{\bar{f}}\hat{g}.$$

Multiplying the second identity by i and adding them together gives the claim.

(b) We have already proved that

$$\int_{\mathbb{R}} f\hat{h} = \int_{\mathbb{R}} \hat{f}h.$$

Applying this identity to $\hat{\bar{f}}$ in place of f we get

$$\langle \hat{h}, \hat{f} \rangle = \int_{\mathbb{R}} \hat{h}\hat{\bar{f}} = \int_{\mathbb{R}} \hat{\bar{f}}h = \int_{\mathbb{R}} \hat{\hat{f}}h = \int_{\mathbb{R}} h\bar{f} = \langle h, f \rangle.$$

5.3. Let \mathbb{T} be an up-tree of tiles. Show that it can be divided into 20 subcollections \mathbb{T}_i so that if $P, P' \in \mathbb{T}_i$ for the same i , then $\phi_P, \phi_{P'}$ are orthogonal to each other. Let $P = I_P \times \omega_P$ be a tile. Remember that

$$\phi_P(x) \stackrel{\text{def}}{=} M_{c(\omega_P)} T_{c(I)} D_{|I|}^2 \phi,$$

and that $\text{supp } \hat{\phi} \subset [-\frac{1}{20}, \frac{1}{20}]$. For fixed $k \in \mathbb{Z}$ consider all the tiles in \mathbb{T} such that $|\omega_P| = 2^k$ so that $|I_P| = 2^{-k}$. The centers of the time intervals of any two tiles $P, P' \in \mathbb{T}$ satisfy $c(I_P) - c(I_{P'})/|I| \in \mathbb{Z}$. Call two tiles in \mathbb{T} equivalent if $|\omega_P| = |\omega_{P'}| = 2^k$ and $c(I_P) - c(I_{P'}) \equiv 0 \pmod{20}$. Call $\mathbb{T}_k^{(1)}, \mathbb{T}_k^{(2)}, \dots, \mathbb{T}_k^{(20)}$, the different equivalent classes. Finally consider the subcollections $\mathbb{T}_j \stackrel{\text{def}}{=} \cup_{k \in \mathbb{Z}} \mathbb{T}_k^{(j)}$ for $j = 1, 2, \dots, 20$. The enumeration of the equivalence classes is irrelevant so it can be arbitrary. Now let $P \in \mathbb{T}_j, P' \in \mathbb{T}_{j'}$ for some j, j' . If P, P' correspond to the same k then they satisfy $|\omega_P| = |\omega_{P'}| = 2^k$. In this case the centers of the time intervals satisfy $c(I_P) - c(I_{P'}) = 20n2^{-k}$. By the calculation in the proof of Proposition 5.2 in the notes we get $\langle \phi_P, \phi_{P'} \rangle = 0$. In the complementary case we have $|\omega_P| \neq |\omega_{P'}|$ so suppose $|\omega_P| < |\omega_{P'}|$. Since both tiles were part of an up-tree, the upper frequency intervals

intersect so we must have $\omega_{P_u} \subsetneq \omega_{P'_u}$ so in fact the whole frequency interval satisfies $\omega_P \subset \omega_{P'_u}$. However the frequency support of each ϕ_P is contained in ω_{P_d} :

$$\text{supp}(\hat{\phi}_P) \subset c(\omega_d) + |\omega_P|[-\frac{1}{20}, \frac{1}{20}] \subset \omega_{P_d}.$$

Thus in this case as well we have $\langle \phi_P, \phi_{P'} \rangle = \langle \hat{\phi}_P, \hat{\phi}_{P'} \rangle = 0$.

5.4. Prove the following fact that was implicitly used in transforming the probabilistic expectation \mathbf{E} into the Lebesgue integral over $[0, 1]$: For numbers $0 \leq a \leq b \leq 1$, we have

$$\mathbf{P}\left(\sum_{j=1}^{\infty} 2^{-j} \beta_j \in [a, b)\right) = b - a,$$

where β_j are independent random variables with $\mathbf{P}(\beta_j = 0) = \mathbf{P}(\beta_j = 1) = 1/2$. First of all observe that the probability that $\sum_{j=1}^{\infty} 2^{-j} \beta_j$ attains any single value is zero. Indeed, for any given $a \in [0, 1)$ we either have that the value a is never attained, thus the probability is zero, or there is a deterministic sequence $\{\epsilon_j\}_{j=1}^{\infty} \subset \{0, 1\}^{\mathbb{Z}}$, depending on a , such that

$$\sum_{j=1}^{\infty} 2^{-j} \beta_j = a \Leftrightarrow \beta_j = \epsilon_j \quad \text{for all } j.$$

Thus

$$\left\{\sum_{j=1}^{\infty} 2^{-j} \beta_j = a\right\} \subseteq \bigcap_{j=1}^N \{\beta_j = \epsilon_j\}$$

for all positive integers N . We conclude that in any case

$$\mathbf{P}\left(\sum_{j=1}^{\infty} 2^{-j} \beta_j = a\right) \leq \prod_{j=1}^N \mathbf{P}(\beta_j = \epsilon_j) = \frac{1}{2^N},$$

by the independence of the β_j 's. Since this holds for arbitrary any we get $\mathbf{P}\left(\sum_{j=1}^{\infty} 2^{-j} \beta_j = a\right) = 0$.

Consider the case where $[a, b) = [0, 1/2)$. Then we have

$$0 \leq \sum_{j=1}^{\infty} 2^{-j} \beta_j \leq 1/2 \Leftrightarrow \beta_1 = 0,$$

so in this case the claim is immediate. Now let $[a, b)$ be any dyadic interval of the form $2^{-k}[j, j+1)$. We have

$$\sum_{j=k+1}^{\infty} 2^{-j} \beta_j \leq \sum_{j=k+1}^{\infty} 2^{-j} \leq 2^{-k},$$

and $\sum_{j=1}^k 2^{-j} \beta_j$ is some number of the form $\ell 2^{-k}$ with $\ell \leq 2^k - 1$ where the exact value of ℓ depends only on the values of β_1, \dots, β_k . Independently of what this exact value is, there exists some sequence of numbers $\epsilon_1, \dots, \epsilon_k \in \{0, 1\}$ such that

$$\mathbf{P}\left(\sum_{j=1}^{\infty} 2^{-j} \beta_j \in 2^{-k}[j, j+1)\right) = \mathbf{P}\left(\bigcap_{j=1}^k \{\beta_j = \epsilon_j\}\right) = 2^{-k} = |[a, b)|,$$

by the independence of the β_j 's. Finally any interval $[a, b)$ can be written as a (possibly infinite) disjoint union of dyadic intervals except maybe the endpoints, that is

$$(a, b) = \bigcup_{m=1}^{\infty} \Delta_m$$

where the Δ_j 's are dyadic and disjoint. Indeed one considers the maximal dyadic intervals in (a, b) so they are disjoint by construction. Now since (a, b) is open any point $x \in (a, b)$ is contained in some dyadic interval and thus in some maximal dyadic interval. The claim follows easily since

$$\begin{aligned} \mathbf{P}\left(\sum_{j=1}^{\infty} 2^{-j} \beta_j \in [a, b)\right) &= \mathbf{P}\left(\sum_{j=1}^{\infty} 2^{-j} \beta_j \in (a, b)\right) = \mathbf{P}\left(\sum_{j=1}^{\infty} 2^{-j} \beta_j \in \bigcup_m \Delta_m\right) \\ &= \sum_m \mathbf{P}\left(\sum_{j=1}^{\infty} 2^{-j} \beta_j \in \Delta_m\right) = \sum_m |\Delta_m| \\ &= |[a, b)| = b - a. \end{aligned}$$

5.5. Given a function ψ and an interval J , denote $\psi_J \stackrel{\text{def}}{=} T_{c(J)} D_{|J|}^2 \psi$. Let I be a fixed standard dyadic interval. Prove that

$$\mathbf{E} \psi_{I+\beta}(x) = (\psi * \mathbf{1}_{[0,1)}) \mathbf{1}_I(x),$$

where \mathbf{E} is the expectation over the random choice of the shift parameter $\beta \in \{0, 1\}^{\mathbb{Z}}$. First of all let us look at $\psi_{I+\beta}$ for a fixed β . We have

$$\psi_{I+\beta}(x) = T_{c(I+\beta)} D_{|I|}^2 \psi = T_{c(I+\beta)} \frac{1}{|I|^{\frac{1}{2}}} \psi\left(\frac{x}{|I|}\right)$$

As in the proof of Theorem 6.1 in the notes we write for $|I| = 2^{-k}$:

$$\begin{aligned} c(I+\beta) - c(I) &= \sum_{2^{-j} < 2^{-k}} \beta_j 2^{-j} = \sum_{j \geq k+1} \beta_j 2^{-j} = \sum_{s=1}^{\infty} \beta_{s+k} 2^{-(s+k)} \\ &= |I| \sum_{j=1}^{\infty} \beta_{j+k} 2^{-s}, \end{aligned}$$

and the binary series above is uniformly distributed in the interval $[0, 1)$ as β_j 's are chosen randomly in $\{0, 1\}^{\mathbb{Z}}$ by exercise 5.4. Thus

$$\begin{aligned} \mathbf{E}\psi_{I+\beta} &= \frac{1}{|I|^{\frac{1}{2}}} \int_0^1 \psi\left(\frac{x - c(I) - u|I|}{|I|}\right) du = \frac{1}{|I|^{\frac{1}{2}}} \psi * \mathbf{1}_{[0,1)}\left(\frac{x - c(I)}{|I|}\right) \\ &= D_{|I|}^2 \psi * \mathbf{1}_{[0,1)}(x - c(I)) = T_{c(I)} D_{|I|}^2 \psi * \mathbf{1}_{[0,1)}(x) \\ &= (\psi * \mathbf{1}_{[0,1)})_I(x). \end{aligned}$$